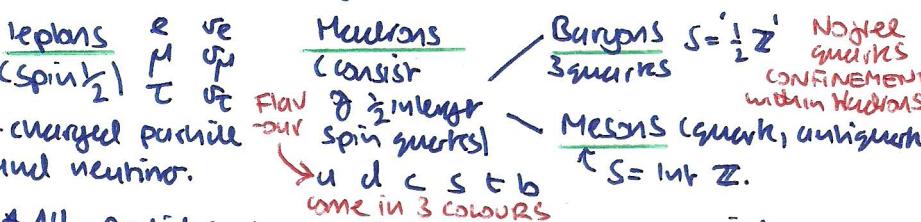


PARTICLE PHYSICS

Study of MATTER (the most fundamental elements that constitute the observable world) and FORCES (the most basic interactions between elements of matter) current understanding embodied in the 'STANDARD MODEL'.

MATTER - consists of leptons and quarks (which form hadrons).



FORCES are due to exchange of elementary particles GAUGE BOSONS

EM photon & WEAK $w^{\pm} Z^0$ + Higgs boson
STRONG gluon g

gravit
yet to be discovered.

All particles have ANTI-PARTICLES. Identified apart from sign interaction. i.e. charge - positron e^+ , electron e^- .

Forces have 'Yukawa' form $F_i = \alpha_i \hbar c \left(\frac{1}{r^2} + \frac{m_i}{r} \right) e^{-m_i r/\hbar c}$ for type of interaction i

at a distance r via exchange of a gauge boson of mass m_i .

for EM, STRONG $m_i=0$ $\alpha_{EM} \sim \frac{1}{137}$ $\alpha_S \sim 1$. $m_W \sim 80.4 \text{ GeV}/c^2$ $m_Z \sim 91.2 \text{ GeV}/c^2$
 $d\omega \sim \frac{1}{r^2}$

Study particle physics using i) static properties {mass, spin, magnetic moments, parity}

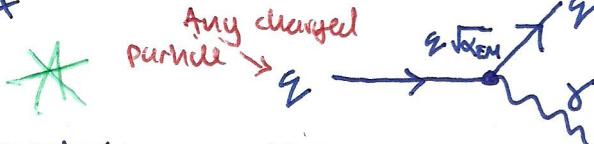
iii) Particle scattering cross section $\sigma = \frac{\pi}{4} \frac{\Gamma_{int}}{f_{inel}}$ $\Gamma_{int} = \frac{2\pi}{\hbar} |M_{int}|^2 g(E)$
and $d\sigma/d\Omega$. (iii) Decays $N(t) = N_0 e^{-\Gamma t/E}$ F.O.R.
Total decay rate = 'Total width' $\Gamma = \sum \Gamma_i$

use theory to predict angular dependence of $|M_{int}|^2$

Partial width $\Gamma_i \Rightarrow \Gamma = \sum \Gamma_i$ Branching ratio $B_i = \Gamma_i / \Gamma$ frame independent

\rightarrow decay mode i . * Feynmann diagrams allow us to compute matrix elements for a particular interaction between particles. Represents a sum over all tree orderings. Add matrix elements for all possible Feynmann diagrams then Γ^2 to get $\Gamma_{int} = \frac{\Gamma^2}{2\pi g(E)}$.
Add matrix elements for all possible Feynmann diagrams from BASIC EM

VERTEX



$$1) \text{e}^- \rightarrow \text{e}^+ \gamma$$

$$\Gamma \propto \frac{1}{137}$$

* Energy, momentum, (linear, angular) charge must be conserved at each vertex.

3) Gauge Theories describe the interaction of particle wavefunctions ensuring that the physics remains invariant under certain SYMMETRY transformations which reflect CONSERVATION LAWS observed in nature.

SYMMETRY operator \Rightarrow existence of UNITARY OPERATOR \hat{U} If $\psi(r,t)$ is an eigenstate of Hamiltonian \hat{H} then so is $\psi'(r,t) = \hat{U}\psi(r,t)$ and $[\hat{U}, \hat{H}] = 0$

E.g. translational invariance $\psi'(r,t) = \psi(r+E, t) = \psi(r, t) + E \frac{d}{dr} \psi(r, t) + \dots$

Let $\hat{U} = 1 + E \frac{d}{dr} = 1 + i E \hat{P}_x$ ($P = \hbar \vec{v}$) Now since $[\hat{U}, \hat{H}] = 0 \Rightarrow [\hat{P}_x, \hat{H}] = 0$

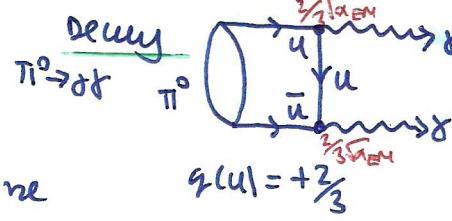
\Rightarrow MOMENTUM CONSERVED. consider unitary operator $\hat{U} = e^{i\vec{q} \cdot \vec{r}}$ if

$\vec{q}(r, t)$ varies in spacetime \rightarrow acts like a field. Local transformation $\psi' = e^{i\vec{q}(r)} \psi$ changes phase (not $1/q^2$) of all particles depending on their position \rightarrow require a physical field to carry these changes. i.e. photon field for EM.

$$\Gamma \propto \left| \frac{2}{3} \frac{1}{\sqrt{137}} \frac{3}{3} \frac{1}{\sqrt{137}} \right|^2$$

and h.c. in diagram Yes.

M for this diagram $\propto \sqrt{\Gamma_{EM}}$ Vertex factors.
if propagator effects are ignored
matrix elements are \prod vertex factors.
 $q = \text{charge}/e.$



$$q(u) = +\frac{2}{3}$$

Note # leptons - # antileptons is conserved [LEPTON # CONSERVATION] as well as # Quarks - # antiquarks [BARYON # CONSERVATION]

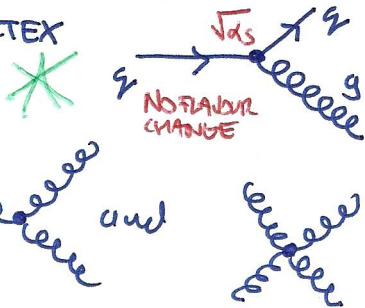
$$e^- \rightarrow \bar{e} \gamma \quad p \rightarrow e^+ \gamma$$

$$\text{e.g. pr E.M. S.E. is } \left[\frac{1}{2m} (-i\nabla - q\vec{A})^2 + q\phi \right] \psi = \frac{\partial A}{\partial t} \quad \text{vector potential } \vec{A} = q\vec{A}$$

Now $\vec{A} \rightarrow \vec{A} + \nabla \lambda$
 $\phi \rightarrow \phi - \frac{\partial \lambda}{\partial t}$ does not alter $E = -\nabla\phi - \frac{\partial A}{\partial t}$
 $\underline{B} = \nabla \times \underline{A}$

1) QCD - Quantum chromodynamics is the theory of the strong interaction.

BASIC VERTEX



plus:



gluons assumed to be massless Spin 1 vector gauge bosons carrying colour quantum number

3 colours + 3 anti-colours.

rgb Twisted by antiquarks

→ could make a colourless state (Glueball) out of gluons. Not observed yet.

* Quarks and gluons

are permanently confined inside hadrons.

Self interaction of gluons when squeezed into a line into tube-like regions. ⇒ Energy / unit length stored in field ~ constant ⇒ separation energy → ∞. If one tries to separate quarks in a $q\bar{q}$ pair (e.g.) this energy creates new pairs - HADRONISATION.



Real gluons are orthogonal linear combinations of colour states $r\bar{b}, r\bar{g}, b\bar{r}, b\bar{g}, g\bar{r}, g\bar{b}, r\bar{r}, b\bar{b}, g\bar{g}$.

Note $r\bar{r} + b\bar{b} + g\bar{g}$ has no 'net colour'; plays no part in strong interaction. Hence 8 coloured gluons.

* Asymptotic freedom

In QED the charges are screened by the charges.

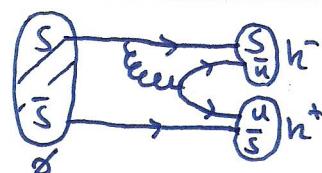
At small distances (\Rightarrow high energy probing particle) - smaller screening effect
- larger effective charge - larger effective value of α_{EM}

Analogous effect in QCD - "color screening":

except due to gluon self interaction effect
is opposite. Gluons appear 'free' at high energies i.e. α_S decreases

* Scattering - → in QED with Feynmann rules for calculating matrix elements.

e.g. $\phi \rightarrow h^+ h^-$



NOTE: PARITY CONSERVED IN EM, STRONG INTERACTIONS

(can be violated in WEAK interactions) - useful diagnostic.

Note whether bosons or fermions are produced ⇒ even if parity conserved the wrong sign can mean interaction not possible e.g. $W^+ \rightarrow \pi^+ \pi^0$ if produced same parity $P=-1 \Rightarrow$ Fermion $P=1 \Rightarrow$ Bosons.

* Evidence for colour

insulin ratio $R = \sigma(e^+ e^- \rightarrow \text{Hadron}) / \sigma(e^+ e^- \rightarrow \mu^+ \mu^-)$

i.e. QED interaction.

i.e. $\text{e}^+ \text{e}^- \rightarrow \text{gg}$ Exptl $R = 3 \sum g_i^2$ for quark colour charge g_i . (Note each event is independent so add $|n|^2$). ↑ 3 due to 3 colours. i sum depends on energy of $e^+ e^-$ collision. As raise energy expect to see steps as more quarks can be created.

$$E > 2M_S \quad R = 3(\frac{4}{3}g_u + \frac{1}{3}g_d + \frac{1}{3}g_s) = 2$$

$$E > 2M_F \quad R = 3(\dots + \frac{4}{3}g_s) = 5$$

$$E > 2M_b \quad R = 3(\dots + \frac{4}{3}g_s) = 3\frac{1}{3}$$

$$R = 3(\dots + \frac{4}{3}g_s) = 3\frac{2}{3}$$

e.g. $\text{e}^+ \text{e}^- \rightarrow \text{gggg}$ - Extra ds in cross section. Experimentally observe hadronic jets - angle depends on gluon spin ($=1$)

Triple gluon vertex is lowest energy gluon radiation

→ modifies angular distribution. $\text{gg} \rightarrow \text{gggg}$

*Hadron masses: Mass differences amongst hadrons are partly due to differences in quark masses and partly due to the spin-spin interaction of confined quarks.

c.g. Hyperfine splitting in H_2 ($\ell=1$) in QED: $\Delta E \propto \frac{\Sigma_p \cdot \Sigma_e}{m_p m_e}$

MESON MASSES $m_m = m_1 + m_2 + A_m \frac{\Sigma_1 \cdot \Sigma_2}{m_1 m_2}$ Note all quarks are spin $\frac{1}{2}$
 Σ_1, Σ_2 are constant. Now $\Sigma_1 \cdot \Sigma_2 = \frac{1}{2} (\Sigma^2 - \Sigma_1^2 - \Sigma_2^2)$ where $\Sigma = \Sigma_1 + \Sigma_2$ so: $m_m = m_1 + m_2 - \frac{3A_m}{4m_1 m_2}$

BARYON MASSES $m_m = m_1 + m_2 + m_3 + A_B \left[\frac{\Sigma_1 \cdot \Sigma_2}{m_1 m_2} + \frac{\Sigma_1 \cdot \Sigma_3}{m_1 m_3} + \frac{\Sigma_2 \cdot \Sigma_3}{m_2 m_3} \right]$ use $\Sigma_1 \cdot \Sigma_2$ above again... 1-MESONS $m_m = m_1 + m_2 + \frac{A}{4m_1 m_2}$
 $\Sigma_1, \Sigma_2, \Sigma_3$ are constant

*Baryon magnetic moments: If all the quarks have zero orbital angular momentum

$M_{\text{Baryon}} = M_1 + M_2 + M_3$ Now for pointlike spin $\frac{1}{2}$ particle of charge g $M = \frac{g}{m} \Sigma$
 $\therefore \Delta M = \frac{g}{m} \frac{\Sigma}{2}$. $\therefore \langle \uparrow | M_{\text{Baryon}} | \uparrow \rangle = \frac{e\hbar}{2} \left(\frac{g_1}{m_1} + \frac{g_2}{m_2} + \frac{g_3}{m_3} \right)$
 actually state $| \uparrow \uparrow \uparrow \rangle_{123}$

*Hadron Resonances: Most hadronic states decay via the strong interaction \Rightarrow short lifetime. Identify state by peak in cross section vs. energy. i.e., Breit-Wigner.

$\sigma(\pi + Y \rightarrow Z \rightarrow A + B) = \frac{\pi \tilde{x}^2 g \Gamma_{\pi Y} \Gamma_{AB}}{(E_0 - E)^2 + \Gamma^2/4}$ For Hadron Z of energy E_0
 $g = \frac{2S_z + 1}{(2S_x + 1)(2S_y + 1)}$ $\Gamma_{\pi Y} = \Gamma(Z \rightarrow \pi + Y)$ $\Gamma_{AB} = \Gamma(Z \rightarrow A + B)$
 $\tilde{x} = \frac{p_x}{p_Z}$ momentum in c.o.m system.

5) Relativistic Wave Equations: Generate (non-relativistic) S.E. via replacing E, P in energy expression by Q.M. operators $E \rightarrow i\hbar \frac{\partial}{\partial t}, P \rightarrow \hbar \vec{k} \nabla$

i.e. $E = \frac{P^2}{2m} + V \rightarrow i\hbar \frac{\partial}{\partial t} \psi = -\frac{\hbar^2}{2m} \nabla^2 \psi + V \psi$ S.E. Now using special relativity $E^2 = p^2 c^2 + m^2 c^4 \Rightarrow -\hbar^2 \frac{\partial^2 \psi}{\partial t^2} = -c^2 \nabla^2 \psi + m^2 c^4 \Rightarrow \nabla^2 \psi - \frac{m^2 c^2}{\hbar^2} \psi = \frac{1}{c^2} \frac{\partial^2 \psi}{\partial t^2}$
 This is the KLEIN GORDON equation. Solutions are of the form $\psi = e^{i(Et - \vec{p} \cdot \vec{r})/\hbar}$ i.e., -ve / +ve Energy solutions
 \hookrightarrow Interpret as $\begin{cases} \text{Particle} - \text{ve } E \\ \text{Antiparticle} + \text{ve } E \end{cases}$ k.b. equation not much use since probability of finding a particle is not conserved..... though we can use it as a foundation to generating Weyl and Dirac equations.

* The Dirac Equation \Rightarrow Aim to generalize k.b. equation $(\nabla^2 - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \frac{m^2 c^2}{\hbar^2}) \psi = 0$

$$\Rightarrow (i \gamma^1 \frac{\partial}{\partial x} + i \gamma^2 \frac{\partial}{\partial y} + i \gamma^3 \frac{\partial}{\partial z} + \frac{i \gamma^0}{c} \frac{\partial}{\partial t} + \frac{mc}{\hbar})(i \gamma^1 \frac{\partial}{\partial x} + i \gamma^2 \frac{\partial}{\partial y} + i \gamma^3 \frac{\partial}{\partial z} + \frac{i \gamma^0}{c} \frac{\partial}{\partial t} - \frac{mc}{\hbar}) \psi = 0$$

γ 's must anticommute to remove cross terms. 4 γ 's \Rightarrow 4x4 matrices.

DIRAC eq is $i \left(\frac{1}{c} \frac{\partial}{\partial t} + \frac{\gamma^1}{c} \frac{\partial}{\partial x} + \frac{\gamma^2}{c} \frac{\partial}{\partial y} + \frac{\gamma^3}{c} \frac{\partial}{\partial z} \right) \psi - \frac{mc}{\hbar} \psi = 0$ from 4 are 4 unpaired SPINORS

$$\gamma^0 = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \quad \gamma^1 = \begin{pmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & -1 & 0 & 0 \\ -1 & 0 & 0 & 0 \end{pmatrix} \quad \gamma^2 = \begin{pmatrix} 0 & 0 & 0 & -i \\ 0 & 0 & i & 0 \\ 0 & i & 0 & 0 \\ -i & 0 & 0 & 0 \end{pmatrix} \quad \gamma^3 = \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \\ -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{pmatrix}$$

$$= \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \end{pmatrix} \quad = \begin{pmatrix} 0 & \gamma_x & 0 & 0 \\ -\gamma_x & 0 & 0 & 0 \end{pmatrix} \quad = \begin{pmatrix} 0 & 0 & \gamma_y & 0 \\ 0 & 0 & 0 & -\gamma_y \end{pmatrix} \quad = \begin{pmatrix} 0 & 0 & \gamma_z & 0 \\ 0 & 0 & 0 & -\gamma_z \end{pmatrix}$$

\hookrightarrow get correct spin eigenvalues when solving Dirac equation for $P=0$

i.e. $i \frac{\gamma^0 \partial_t}{\hbar} - \frac{mc^2}{\hbar} \psi = 0 \Rightarrow \psi = \begin{pmatrix} \exp(-mc^2 t/\hbar) \\ " \\ \exp(mc^2 t/\hbar) \end{pmatrix}$ ^{fermions} _{antifermions} ^{Dirac Spin} $\Sigma = \begin{pmatrix} \Sigma_x & 0 & 0 & 0 \\ 0 & \Sigma_y & 0 & 0 \\ 0 & 0 & \Sigma_z & 0 \\ 0 & 0 & 0 & \Sigma_z \end{pmatrix} \frac{\hbar}{2}$ pp(3)

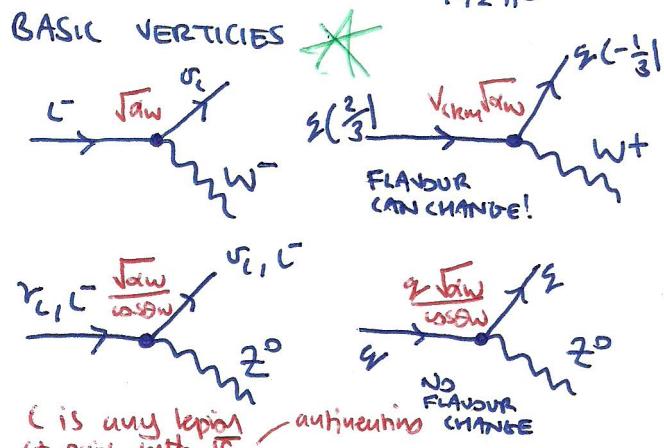
$$\Rightarrow \underline{S}_z \psi = \frac{i}{2} \begin{pmatrix} \exp(-mc^2 t/\hbar) \\ -\exp(-mc^2 t/\hbar) \\ \vdots \\ \vdots \end{pmatrix} \quad \begin{array}{l} \text{Particle, Spin up} \\ \text{" Spin down} \\ \text{Antiparticle spin up} \\ \text{" Spin down} \end{array}$$

For solutions $p \neq 0$
 Try $\psi = U e^{i(Et - p \cdot r)/\hbar}$
 where $E = \sqrt{p^2 + m^2 c^2}$
 pr $U_1, U_2 \sim \sqrt{p^2 + m^2} e^{iE t}$
 pr U_3, U_4 .

- Equipped with Dirac equation
 one can calculate matrix elements properly
 from Feynmann Diagrams. (using Feynmann rules not
 discussed here).

6) Weak interactions. Fermi theory of β decay fails at high energies (Sargent rule
 for leptonic decay $\Gamma_X = \frac{G_F^2 M_X^5}{192\pi^3} \quad G_F = 1.166 \times 10^{-5} \text{ GeV}^{-2}$).

BASIC VERTICES



(is any lepton
 & pairs with τ_i) - antineutrino change

Weak interactions allow parity violation (unlike EM and strong). reason? vertex
 factor for weak interactions includes a term $\propto \frac{1}{2}(1-\delta^5)$ which acts as the
 'HELICITY' operator i.e. selects left handed particles (and right handed antiparticles)
 ONLY at high energy. Left handedness for particles $\Rightarrow \underline{S}$ is uniparallel to \underline{p} .

* Z^0 decays - why there are only 3 types of neutrinos

- measure total width for $Z^0 \rightarrow \dots$ decay, Γ_{Z^0} . Now since all Z^0
 \rightarrow leptons and $Z^0 \rightarrow$ quarks (well nucleons) have the same decay rates
 can write $\Gamma_{Z^0} = 3\Gamma_{e\bar{e}} + 3\Gamma_{u\bar{u}} + N_\nu \Gamma_{\nu\bar{\nu}} = 3\Gamma_{e\bar{e}} + 3\Gamma_{d\bar{d}} + N_\nu \Gamma_{\nu\bar{\nu}}$
 $(N_\nu \Gamma_{\nu\bar{\nu}}$ is an 'invisible' part)
 width with it found by experimentally
 measuring $\Gamma_{Z^0}, \Gamma_{e\bar{e}}, \Gamma_{d\bar{d}}$ etc. Using
 theory for $\Gamma_{\nu\bar{\nu}}$ one can work out N_ν)

$$\begin{aligned} \Delta W &\sim \frac{1}{2\theta} \\ V_{CKM} &\approx \begin{pmatrix} 1-\frac{\Delta^2}{2} & \Delta & \Delta^2 \\ -\Delta & 1-\frac{\Delta^2}{2} & \Delta^2 \\ \Delta^2 & -\Delta^2 & 1 \end{pmatrix} \\ &= \begin{pmatrix} \sqrt{u_d} & \sqrt{u_s} & \sqrt{u_b} \\ \sqrt{c_d} & \sqrt{c_s} & \sqrt{c_b} \\ \sqrt{t_d} & \sqrt{t_s} & \sqrt{t_b} \end{pmatrix} \end{aligned}$$

parametrizes coupling constants between W^+ decaying quark states

W^\pm, Z^0 are massive

so can only exist above threshold energy.

$$\begin{aligned} \Delta &= \sin \theta_C \\ \theta_C &\approx 13^\circ \\ \theta_W &\approx 29^\circ \end{aligned}$$

exists due to standard model.

* Beyond the Standard Model

Successes: Describes all existing particle physics data to extreme accuracy
Problems: lots of free parameters (~ 20) i.e. masses of leptons, quarks, couplings $\alpha_m, \alpha_s, \Delta_W$,
 θ_W, θ_C , quark and lepton charges, mixing (V_{CKM}); gravity not included.

* Higgs Boson H^0 (neutral spin 0 boson) thought to give particles mass via their interaction with the Higgs field (consisting of scalar called Higgs Boson!).

New possible vertices depending on Higgs mass. i.e. $H^0 \rightarrow Z^0 \rightarrow 2\mu^+ \mu^-$ or $H^0 \rightarrow 2f$.

* Neutrino oscillations ν_e, ν_μ, ν_τ can oscillate i.e. ν_e can become ν_τ spontaneously over time! Explains odd results observed at SUPERKAMIOKANDE which detects annual rate of atmospheric/Solar muon and electron neutrinos. (uses CHERENKOV radiation - light emitted when charged particle exceeds $\frac{c}{n}$ in a medium). use $\begin{pmatrix} \nu_\mu \\ \nu_e \end{pmatrix} = \begin{pmatrix} \cos \theta_C & \sin \theta_C \\ -\sin \theta_C & \cos \theta_C \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix}$ & mixing angle θ_C & lepton # & neutrino mass eigenstates

$$\frac{\partial \nu_\mu}{\partial t} = \frac{1}{n} \theta_C \quad \text{for when } \frac{c}{n} < 1 \quad \Rightarrow \nu_\mu = \frac{1}{n} \theta_C \nu_\mu$$

$$\text{yield} \quad \left| \frac{\nu_\mu(t)}{\nu_\mu(0)} \right|^2 = 1 - \sin^2 \theta_C \sin^2 \left[\frac{(E_1 - E_2)t}{2\pi} \right] \quad \text{as a model for N. oscillations}$$

$$1 \text{ b} = 10^{-28} \text{ m}^2$$

↑ cross section units are often in BARNS (b)

V.Gibson
Lent 2000

PARTICLE PROPERTIES

From the *Review of Particle Physics*, European Physical Journal C3 (1998)

| Quarks (spin 1/2) | | | |
|-------------------|---------|----------------------------|------------|
| Name | Flavour | Mass (GeV/c ²) | Charge (e) |
| up | u | ≈ 0.35 | +2/3 |
| down | d | $m_d \approx m_u$ | -1/3 |
| charm | c | 1.5 | +2/3 |
| strange | s | 0.5 | -1/3 |
| top | t | $174(\pm 5)$ | +2/3 |
| bottom | b | 4.5 | -1/3 |

Parity :

Quark and charged leptons $P = +1$

Antiquarks and antileptons $P = -1$

Parity of 2 particles with A.M C is

$$P = P_1 P_2 (-1)^L$$

| Leptons (spin 1/2) | | | | | |
|--------------------|---------|----------------------------|-------------------------------------|--|--|
| Lepton | Charge | Mass (MeV/c ²) | Mean life (s) | Lepton Decay Mode | Branching Ratio (%) |
| ν_e | 0 | < 15 eV/c ² | stable | | |
| ν_μ | 0 | < 0.17 | stable | | |
| ν_τ | 0 | < 18.2 | stable | | |
| e | ± 1 | 0.511 ^a | stable | | |
| μ | ± 1 | 105.658 ^b | 2.197×10^{-6} ^c | $e^- \bar{\nu}_e \nu_\mu$ | ≈ 100 |
| τ | ± 1 | 1777.0(± 3) | $290.0(\pm 12) \times 10^{-15}$ | $\mu^- \bar{\nu}_\mu \nu_\tau$ $e^- \bar{\nu}_e \nu_\tau$ hadrons $+ \nu_\tau$ | 17.37(± 9) 17.81(± 7) ≈ 65 |

^a The error on the e mass is 1.5×10^{-7} MeV/c².

^b The error on the μ mass is 3.4×10^{-5} MeV/c².

^c The error on the μ lifetime is 4×10^{-11} s.

N.B. Numbers given in brackets correspond to the error in the last digit.

For example, $m_\tau = 1777.0(\pm 3)\text{MeV}/c^2 \equiv (1777.0 \pm 0.3)\text{MeV}/c^2$.

Gauge Bosons ($J^P = 1^-$)

| Force | Gauge Boson | Charge (e) | Mass (GeV/c 2) | Full Width (GeV) | Decay Mode | Branching Ratio (%) |
|----------------|-------------|-----------------------|--------------------------------------|------------------|---|--|
| E-M | γ | $< 5 \times 10^{-30}$ | $< 2 \times 10^{-16} \text{ eV}/c^2$ | stable | | |
| Weak (Charged) | W^\pm | ± 1 | $80.41(\pm 10)$ | $2.06(\pm 6)$ | $e\nu_e$ $\mu\nu_\mu$ $\tau\nu_\tau$ hadrons | $10.9(\pm 4)$ $10.2(\pm 5)$ $11.3(\pm 8)$ $67.8(\pm 10)$ |
| Weak (Neutral) | Z^0 | 0 | $91.187(\pm 7)$ | $2.490(\pm 7)$ | ee $\mu\mu$ $\tau\tau$ $\nu\nu$ hadrons | $3.366(\pm 8)$ $3.367(\pm 13)$ $3.360(\pm 15)$ $20.01(\pm 16)$ $69.90(\pm 15)$ |
| Strong | g | 0 | 0 | stable | | |

Pseudoscalar Mesons ($J^P = 0^-$)

| Particle | Quark Content | Mass (MeV/c ²) | Mean Life (s) or Width (keV) | Decay Mode | Branching Ratio (%) |
|----------------------|--|----------------------------|---|--|---|
| π^\pm | u \bar{d} , d \bar{u} | 139.5700(± 4) | $2.6033(\pm 5) \times 10^{-8}$ | $\mu^- \bar{\nu}_\mu$ | ≈ 100 |
| π^0 | (u \bar{u} - d \bar{d})/ $\sqrt{2}$ | 134.9764(± 6) | $8.4(\pm 6) \times 10^{-17}$ | $\gamma\gamma$ | 98.80(± 3) |
| η | see note a | 547.3(± 1) | 1.2(± 1) | $\gamma\gamma$ $\pi^0 \pi^0 \pi^0$ $\pi^+ \pi^- \pi^0$ $\pi^+ \pi^- \gamma$ $\pi^+ \pi^- \eta$ $\rho^0 \gamma$ $\pi^0 \pi^0 \eta$ | 39.2(± 3) 32.2(± 4) 23.1(± 5) 4.8(± 1) 44(± 2) 30(± 1) 21(± 1) |
| η' | see note a | 957.8(± 1) | 0.20(± 2) | | |
| K^\pm | u \bar{s} , s \bar{u} | 493.677(± 16) | $1.239(\pm 2) \times 10^{-8}$ | $\mu^- \bar{\nu}_\mu$ $\pi^- \pi^0$ $\pi^+ \pi^- \pi^-$ $\pi^0 \mu^- \bar{\nu}_\mu$ $\pi^0 e^- \bar{\nu}_e$ $\pi^+ \pi^-$ $\pi^0 \pi^0$ | 63.5(± 2) 21.2(± 1) 5.59(± 5) 3.18(± 8) 4.82(± 6) 68.6(± 3) 31.4(± 3) |
| K^0, \bar{K}^0 | d \bar{s} , s \bar{d} | 497.67(± 3) | $K_S^0 0.8934(\pm 8) \times 10^{-10}$ $K_L^0 5.17(\pm 4) \times 10^{-8}$ | $\pi^0 \pi^0 \pi^0$ $\pi^+ \pi^- \pi^0$ $\pi^\pm \mu^\mp \nu_\mu$ $\pi^\pm e^\mp \nu_e$ | 21.1(± 3) 12.6(± 2) 27.2(± 3) 38.8(± 3) |
| D^\pm | c \bar{d} , d \bar{c} | 1869.3(± 5) | $1.06(\pm 2) \times 10^{-12}$ | $e^- + \text{any}^b$ $K^- + \text{any}$ $K^+ + \text{any}$ $K^0 + \text{any}$ plus $\bar{K}^0 + \text{any}$ $K^- + \text{any}^c$ $K^+ + \text{any}$ $e^+ + \text{any}$ $\mu^+ + \text{any}$ $\bar{K}^0 + \text{any}$ plus $K^0 + \text{any}$ | 17(± 2) 24(± 3) 6(± 1) 59(± 7) 53(± 4) 3.4(± 5) 6.8(± 3) 6.6(± 8) 42(± 5) |
| D^0, \bar{D}^0 | u \bar{c} , c \bar{u} | 1864.6(± 5) | $0.415(\pm 4) \times 10^{-12}$ | | |
| D_s^\pm | c \bar{s} , s \bar{c} | 1968.5(± 6) | $0.47(\pm 2) \times 10^{-12}$ | seen | |
| B^\pm | u \bar{b} , b \bar{u} | 5279(± 2) | $1.65(\pm 4) \times 10^{-12}$ | seen | |
| B^0, \bar{B}^0 | d \bar{b} , b \bar{d} | 5279(± 2) | $1.56(\pm 4) \times 10^{-12}$ | seen | |
| B_s^0, \bar{B}_s^0 | s \bar{b} , b \bar{s} | 5369(± 2) | $1.54(\pm 7) \times 10^{-12}$ | seen | |
| η_c | c \bar{c} | 2980(± 4) | 13(± 4) MeV | hadrons | |

^a η and η' are linear combinations of the quark state (u \bar{u} + d \bar{d})/ $\sqrt{2}$ and s \bar{s} .

^b D⁻ decay modes; ^c D⁰ decay modes.

| Vector Mesons ($J^P = 1^-$) | | | | | |
|-------------------------------|--|----------------------------|-------------------|---|--|
| Particle | Quark Content | Mass (MeV/c ²) | Full Width (MeV) | Decay Mode | Branching Ratio (%) |
| ρ^\pm | $u\bar{d}, d\bar{u}$ | 770.0(± 8) | 151(± 1) | $\pi\pi$ | 100 |
| ρ^0 | $(u\bar{u} - d\bar{d})/\sqrt{2}$ | | | | |
| ω | $(u\bar{u} + d\bar{d})/\sqrt{2}$ | 781.9(± 1) | 8.41(± 9) | $\pi^+\pi^-\pi^0$ $\pi^0\gamma$ $\pi^+\pi^-$ | 88.8(± 7) 8.5(± 5) 2.2(± 3) |
| ϕ | $s\bar{s}$ | 1019.413(± 8) | 4.43(± 5) | K^+K^- $K_L^0 K_S^0$ | 49.1(± 8) 34.1(± 6) |
| $K^{*\pm}$ | $u\bar{s}, s\bar{u}$ | 891.7(± 3) | 50.8(± 9) | $K\pi$ | ≈ 100 |
| K^{*0}, \bar{K}^{*0} | $d\bar{s}, s\bar{d}$ | 896.1(± 3) | 50.5(± 6) | $K\pi$ | ≈ 100 |
| $D^{*\pm}$ | $c\bar{d}, d\bar{c}$ | 2010.0(± 5) | < 0.13 | $D^0\pi^{-a}$ $D^-\pi^0$ | 68(± 1) 31(± 3) |
| D^{*0}, \bar{D}^{*0} | $u\bar{c}, c\bar{u}$ | 2006.7(± 5) | < 2.1 | $D^0\pi^{0b}$ $D^0\gamma$ | 62(± 3) 38(± 3) |
| $D_s^{*\pm}$ | $c\bar{s}, s\bar{c}$ | 2112.4(± 7) | < 1.9 | seen | |
| B^* | $u\bar{b}, b\bar{u}, d\bar{b}, b\bar{d}, s\bar{b}, b\bar{s}$ | 5325(± 2) | | $B\gamma$ seen | |
| J/ψ | $c\bar{c}$ | 3096.88(± 4) | 87(± 5) keV | hadrons e^+e^- $\mu^+\mu^-$ $\tau^+\tau^-$ e^+e^- $\mu^+\mu^-$ | 87.7(± 5) 6.0(± 2) 6.0(± 2) 2.7(± 2) 2.5(± 2) 2.48(± 7) |
| $\Upsilon(1s)$ | $b\bar{b}$ | 9460.4(± 2) | 53(± 2) keV | | |

^a D^{*-} decay modes; ^b D^{*0} decay modes.

Baryons ($J^P = 1/2^+$)

| Particle | Quark Content | Mass (MeV/c ²) | Mean Life (s) or Full Width (MeV) | Decay Mode | Branching Ratio (%) |
|---------------|---------------|----------------------------|-----------------------------------|----------------------|------------------------------------|
| p | uud | 938.2723(± 3) | $> 1.6 \times 10^{25}$ years | | |
| n | udd | 939.5656(± 3) | 887(± 2) | $pe^- \bar{\nu}_e$ | 100 |
| Λ^0 | uds | 1115.683(± 6) | $2.63(\pm 2) \times 10^{-10}$ | $p\pi^-$ $n\pi^0$ | 63.9(± 5) 35.8(± 5) |
| Σ^+ | uus | 1189.37(± 7) | $0.799(\pm 4) \times 10^{-10}$ | $p\pi^0$ $n\pi^+$ | 51.6(± 3) 48.3(± 3) |
| Σ^0 | uds | 1192.64(± 2) | $7.4(\pm 7) \times 10^{-20}$ | $\Lambda^0\gamma$ | 100 |
| Σ^- | dds | 1197.45(± 3) | $1.48(\pm 1) \times 10^{-10}$ | $n\pi^-$ | 99.848(± 5) |
| Ξ^0 | uss | 1314.9(± 6) | $2.90(\pm 9) \times 10^{-10}$ | $\Lambda^0\pi^0$ | 99.54(± 5) |
| Ξ^- | dss | 1321.3(± 1) | $1.64(\pm 2) \times 10^{-10}$ | $\Lambda^0\pi^-$ | 99.89(± 4) |
| Λ_c^+ | udc | 2284.9(± 6) | $2.1(\pm 1) \times 10^{-13}$ | seen | |
| Λ_b | udb | 5624(± 9) | $1.14(\pm 8) \times 10^{-12}$ | seen | |

Baryons ($J^P = 3/2^+$)

| | | | | | |
|------------|----------------------|-------------------|-------------------------------|--|--|
| Δ | uuu, uud udd, ddd | ≈ 1232 | ≈ 120 | $N\pi$ | > 99 |
| Σ^* | uus, uds, dds | ≈ 1385 | ≈ 36 | $\Lambda^0\pi$ $\Sigma\pi$ | 88(± 2) 12(± 2) |
| Ξ^* | uss, dss | ≈ 1530 | ≈ 9 | $\Xi\pi$ | 100 |
| Ω^- | sss | 1672.5(± 3) | $0.82(\pm 1) \times 10^{-10}$ | Λ^0K^- $\Xi^0\pi^-$ $\Xi^-\pi^0$ | 67.8(± 7) 23.6(± 7) 8.6(± 4) |

MATTER

Leptons

Spin $\frac{1}{2}$

charged
neutral

neutrinos (massless?)

Nucleons (consist of quarks)

Baryons (3 quarks) $\frac{1}{2}$ integer spin
(anti....)

Mesons (quark + antiquark pair)
1 integer spin