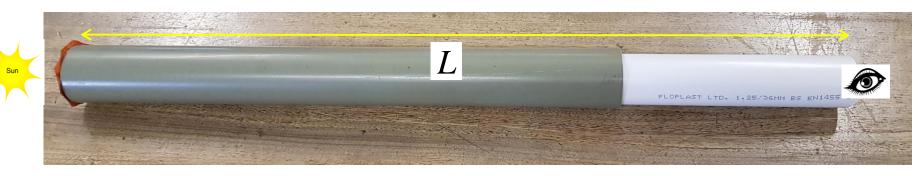


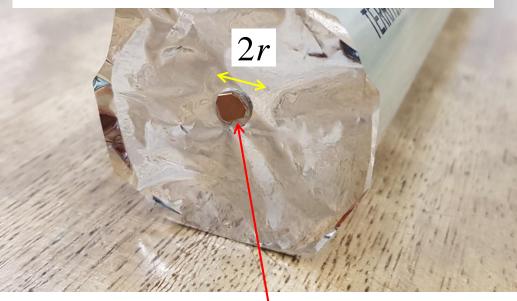
i.e. measuring the angular size of the Sun and hence the surface temperature of the Sun

Tony Ayres & Andrew French. June 2021



WARNING: <u>NEVER</u> STARE DIRECTLY AT THE SUN WITH A NAKED EYE

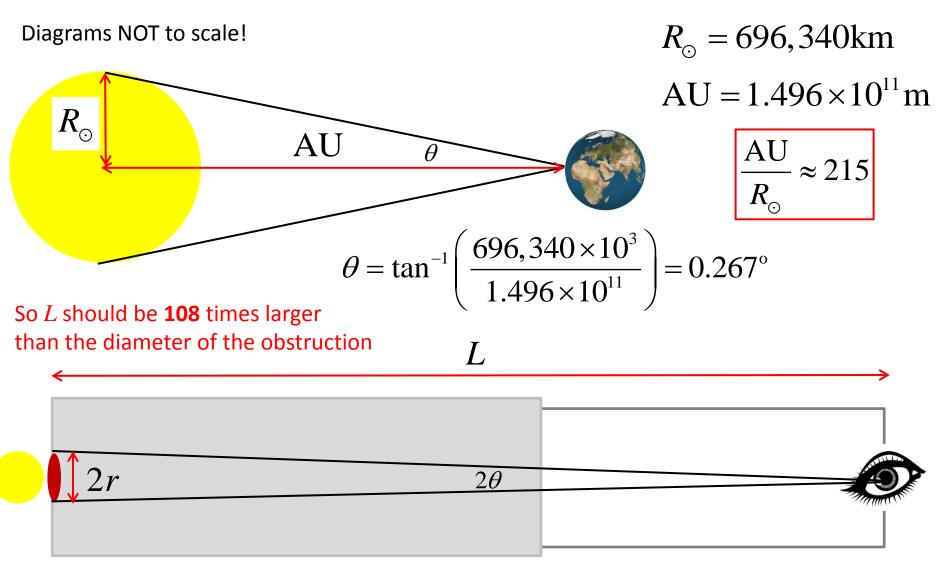
Thin aluminium 'space blanket' foil. Attenuates sunlight to allow for safe viewing, but not opaque.



Look at the Sun from this end. Adjust the white tube length until the opaque obstruction eclipses the Solar disc.



Opaque obstruction of radius r



Without the opaque obstruction, the apparent size of the solar disc would look like:

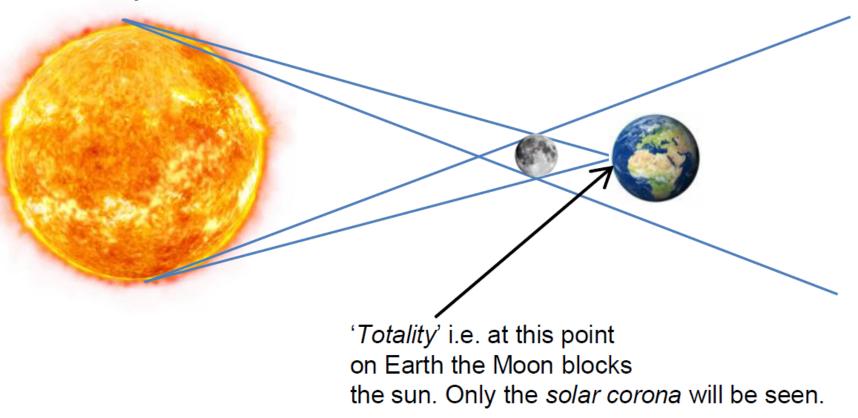
So if we extend L until the disc is *just eclipsed* by the obstruction, we can find angle θ , if we measure r and L.

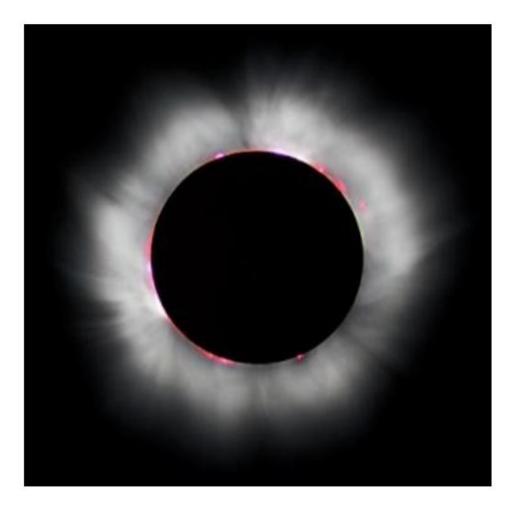
 $\theta = \tan^{-1} \left(\frac{r}{I} \right)$

Solar and lunar eclipses

A total lunar eclipse is when the moon is positioned in the umbra of the Sun's light, obscured by the Earth. i.e. Light cannot reach the surface of the moon from the Sun as the Earth is in the way. This means no light is reflected back to Earth from the Moon's surface, which is why it appears black during a lunar eclipse.

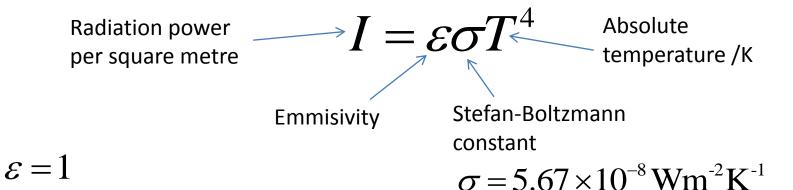
Solar Eclipse

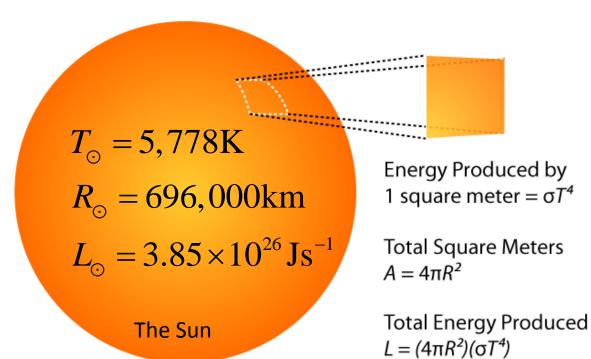




Solar flares and corona observed from Earth during a total solar eclipse

Looking at the Sun through our Solar Tube should reveal something similar. (Although we won't see a corona!)



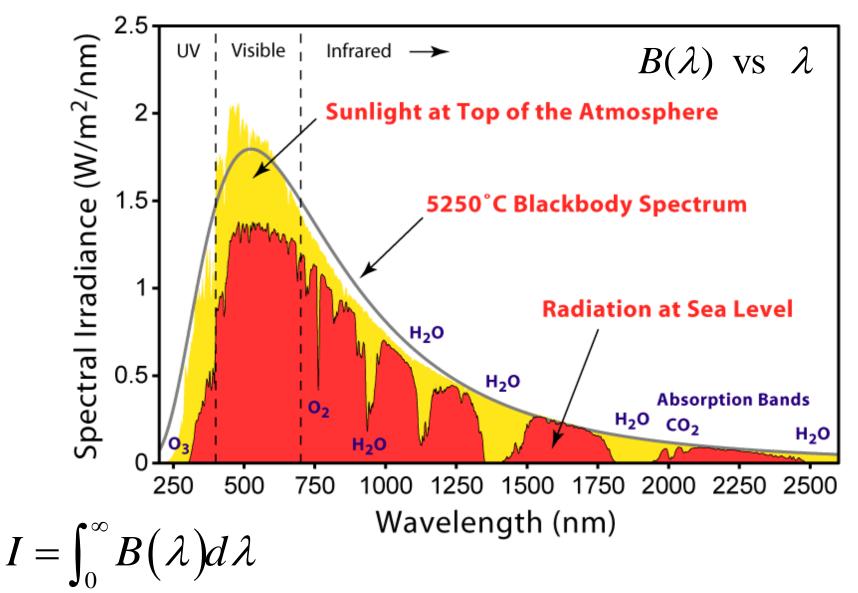


For a 'Black Body' at 20°C = 293K

$$I = 418 \text{Wm}^{-2}$$

It is interesting to compare this to the maximum solar energy incident upon the Earth, which is on average about 1,361 Wm⁻² The measured solar irradiance (i.e. power received on Earth per square metre within a wavelength interval

i.e. $\lambda \rightarrow \lambda + d\lambda$



Red	620-750nm
Yellow	570-590nm
Green	495-570nm
Blue	450-495nm

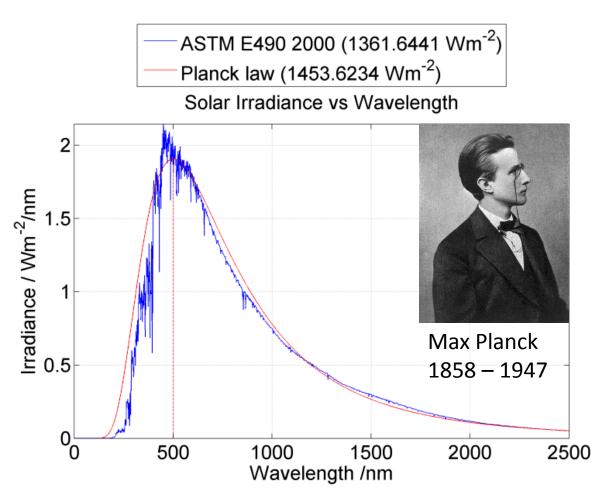
Boltzmann's constant

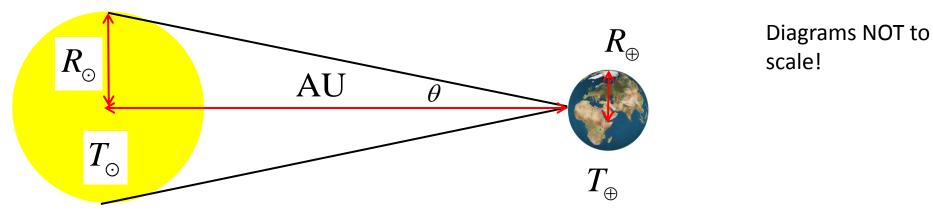
Planck's constant

 $k_{\rm B} = 1.381 \times 10^{-23} \,{\rm m}^2 {\rm kg s}^{-2} {\rm K}^{-1}$ $h = 6.626 \times 10^{-34} \,\mathrm{m^2 kg s^{-1}}$ Speed of light $c = 2.998 \times 10^8 \,\mathrm{ms}^{-1}$

$$I = \int_0^\infty B(\lambda, T) d\lambda = \sigma T^4$$
$$\sigma = \frac{2\pi^5 k_B^4}{15c^2 h^3}$$
$$B(\lambda, T) = \frac{2hc^2}{\lambda^5} \frac{1}{e^{\frac{hc}{\lambda k_B T}} - 1}$$







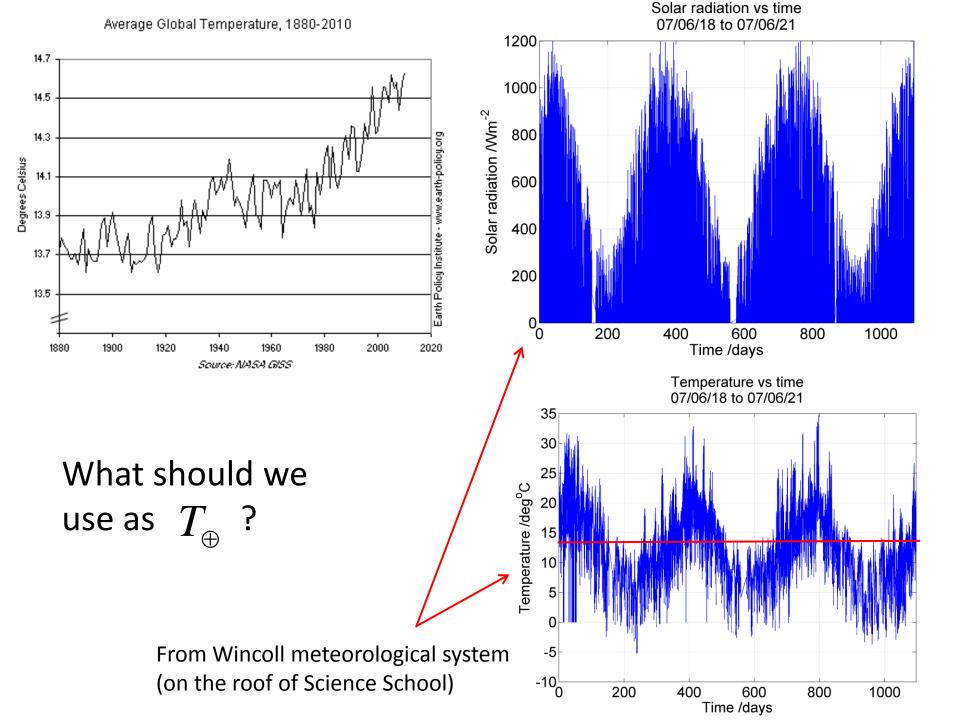
Assume the Sun radiates isotropically. The power per unit area at distance AU illuminating the Earth is:

$$\Phi = \frac{4\pi R_{\oplus}^2 \sigma T_{\odot}^4}{4\pi A U^2}$$

Assume the Earth receives this power over area πR_{\oplus}^2 and then, on average, *re-radiates this energy as infra-red radiation*, isotropically, based upon a mean average Earth temperature. By equating power received = power re-radiated:

$$\frac{4\pi R_{\odot}^2 \sigma T_{\odot}^4}{4\pi A U^2} \times \pi R_{\oplus}^2 = 4\pi R_{\oplus}^2 \sigma T_{\oplus}^4$$
$$\therefore T_{\odot} = \sqrt{\frac{2AU}{R_{\odot}}} T_{\oplus} = \sqrt{\frac{2}{\tan \theta}} T_{\oplus}$$

So if we can calculate angle θ , we can calculate the temperature of the Sun from an average Earth temperature.



To get some idea, let's work backwards from the actual angle θ and the actual temperature of the surface of the Sun.

$$T_{\odot} = \sqrt{\frac{2}{\tan \theta}} T_{\odot}$$
$$T_{\oplus} = \sqrt{\frac{1}{2} \tan \theta} \times T_{\odot}$$
$$T_{\oplus} = \sqrt{\frac{1}{2} \tan \left(0.267^{\circ}\right)} \times 5,778 \text{K}$$
$$T_{\oplus} = 279 \text{K} = 5.9^{\circ} \text{C}$$

So perhaps a bit too cold to be believable, but remember the Earth will not re-radiate *all* the solar energy. Plenty of energy is locked up in biomass and indeed in the kinetic and potential energies of moving molecules within the troposphere. Including us!

