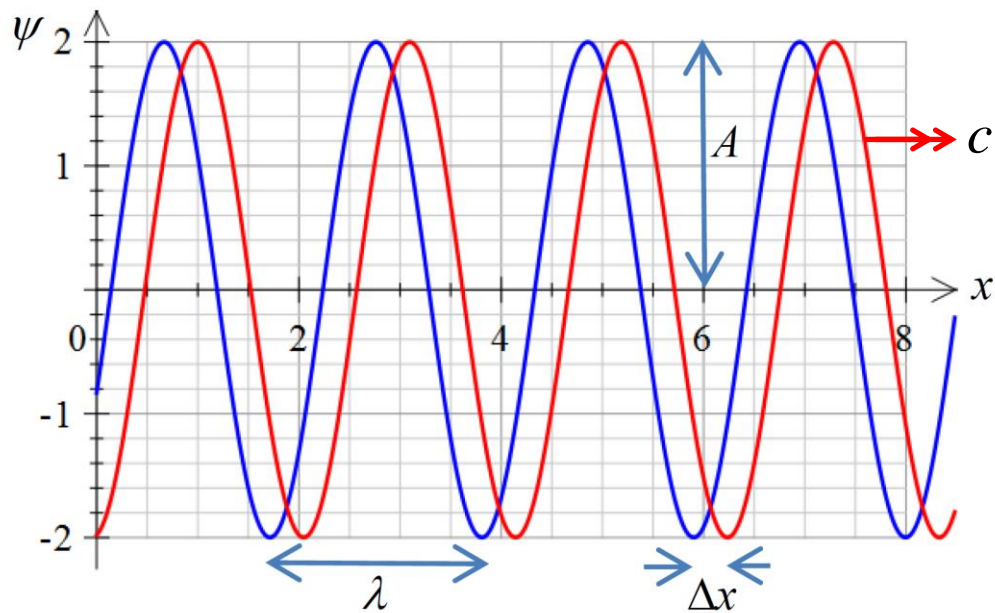


# BPhO

## Computational Challenge

# Waves and optics

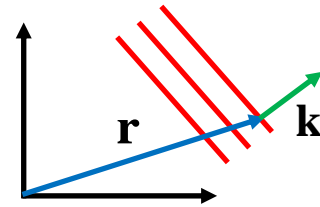
Dr Andrew French.  
December 2023.



—  $\psi(x, t = 0)$   
—  $\psi\left(x, t = \frac{\Delta x}{c}\right)$

$$\psi(x, t) = A e^{i(kx - \omega t - \phi)}$$

$$\psi(\mathbf{r}, t) = A e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t - \phi)}$$



Wave speed  $c = f \lambda$

Frequency = 1/period  $f = \frac{1}{T}$

$$\omega = 2\pi f$$

Wavenumber  $k = \frac{2\pi}{\lambda}$

$$\omega = ck$$

$$\psi(x, t) = A \cos\left(2\pi \frac{x - ct}{\lambda} - \phi\right)$$

$$\psi(x, t) = A \cos(kx - \omega t - \phi)$$

Wave Fourier components are **translated sinusoids**

```

A = 1;           %Amplitude
lambda = 0.5;   %Wavelength /m
c = 340;        %Wavespeed /ms^-1
f_unit = 'kHz'; %Frequency unit
t_unit = 'ms';  %Time unit

%Spatial sinusoid separation (fraction of
wavelength)
dx = lambda/8;

fsize = 18;     %FontSize
N = 1000;      %Number of data points
Nwaves = 3;    %Number of waves to plot
linewidth = 2; %Line width

```

```
%
```

```

%Determine period /s and frequency /Hz
f = c/lambda;
T = 1/f;

```

```

%Determine time delay associated with dx
dt = dx/c;

```

```

%Generate x,y and t,y data vectors
x = linspace( 0, Nwaves*lambda, N);
t = linspace( 0, Nwaves*T, N);

```

```

yx1 = A*sin( 2*pi*x/lambda );
yx2 = A*sin( 2*pi*( x- dx)/lambda );
yt = A*sin( 2*pi*t/T );

```

```

%Plot y vs x graph
figure('color',[1 1 1],'name','waves anatomy')
plot(x,yx1,'b-',x,yx2,'r-', 'linewidth',linewidth);
legend( {'t = 0','t = \Deltat'}, 'fontsize',fsize )
xlabel( 'x /m', 'fontsize',fsize )
title([' A = ',num2str(A), ', \lambda = ',num2str(lambda), ...
      ',m, c = ',num2str(c), 'm/s'], 'fontsize',fsize )
grid on;
set( gca, 'fontsize',fsize )
axis tight
print((gcf, 'wave y vs x.png', '-dpng', '-r300' )
      clf;

```

```

%Plot y vs t graph
plot(t/tuf( t_unit ),yt, 'linewidth',linewidth);
xlabel( ['t /',t_unit], 'fontsize',fsize )
title([' A = ',num2str(A), ', \lambda = ',num2str(lambda), ...
      ',m, c = ',num2str(c), 'm/s, f = ',...
      num2str( f/fuf( f_unit),5),f_unit, ', T = ',...
      num2str(T/tuf( t_unit ),5),t_unit], 'fontsize',fsize )
grid on;
set( gca, 'fontsize',fsize )
axis tight
print( gcf, 'wave y vs t.png', '-dpng', '-r300' )
close(gcf);

```

```
%%
```

```

%Frequency unit factor
function m = fuf( f_unit )
if strcmp( f_unit, 'THz' )==1
    m = 1e12;
elseif strcmp( f_unit, 'GHz' )==1
    m = 1e9;
elseif strcmp( f_unit, 'MHz' )==1
    m = 1e6;
elseif strcmp( f_unit, 'kHz' )==1
    m = 1e3;
else; m = 1; end

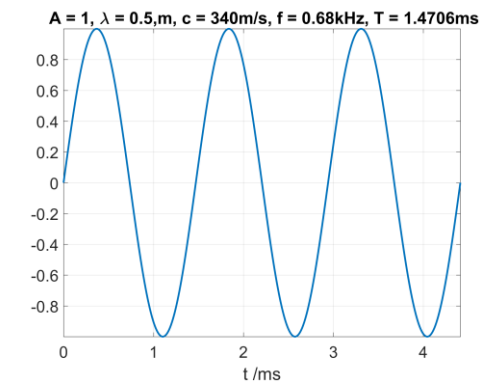
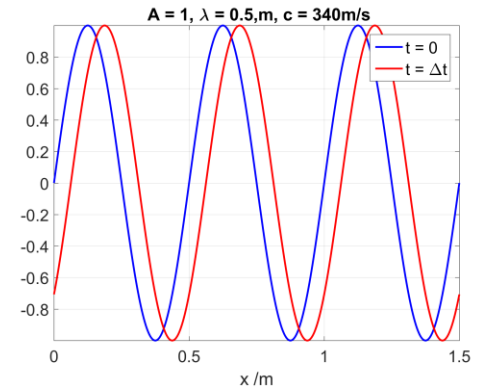
```

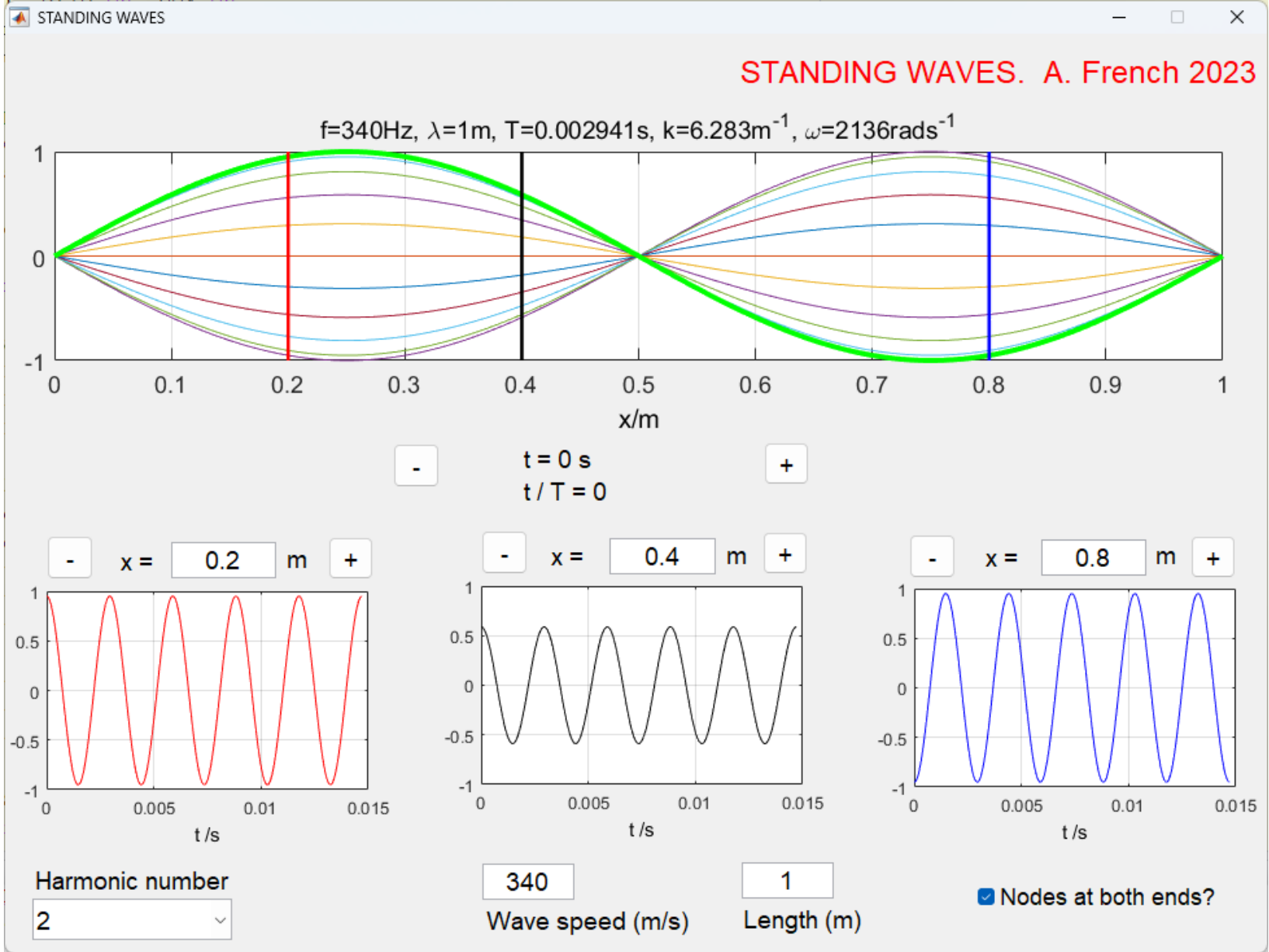
```
%
```

```

%Time unit factor
function m = tuf( t_unit )
if strcmp( t_unit, 'ps' )==1
    m = 1e-12;
elseif strcmp( t_unit, 'ns' )==1
    m = 1e-9;
elseif strcmp( t_unit, '\mus' )==1
    m = 1e-6;
elseif strcmp( t_unit, 'ms' )==1
    m = 1e-3;
else; m = 1; end

```

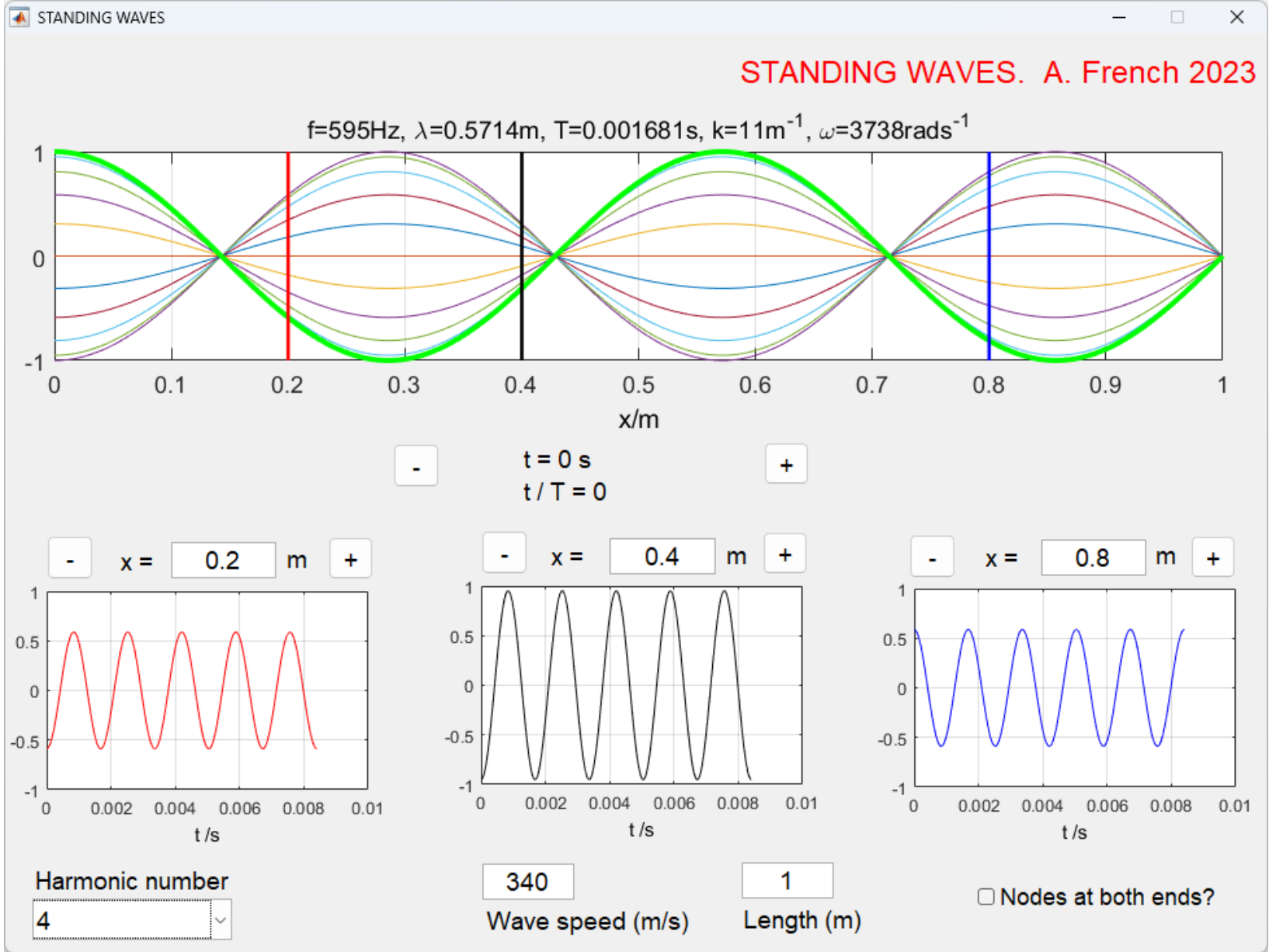




Length is a whole number of half wavelengths

$$L = n \frac{1}{2} \lambda$$

$$n = 1, 2, 3, 4, \dots$$

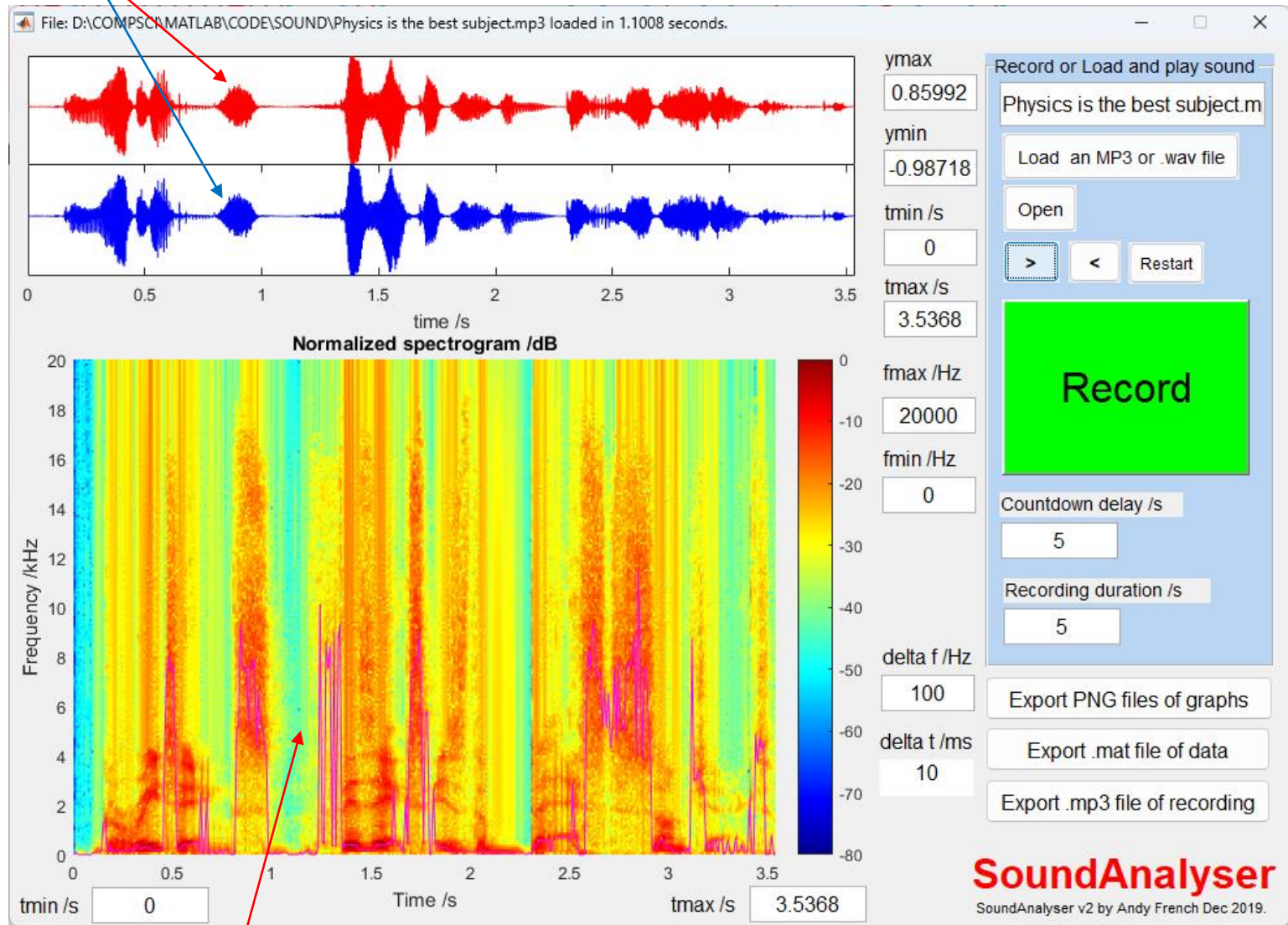


Length (plus *end correction*) is an odd number of quarter wavelengths

$$L + \frac{2}{3}r = (2n - 1)\frac{1}{4}\lambda$$

$$n = 1, 2, 3, 4, \dots$$

# Left and right stereo channels of a sound wave recording



Plot of **frequency content** of waves vs time.  
Colour scale is **decibels** (dB) with max power set at 0dB.

$$\text{dB} = 10 \log_{10} (\text{signal power})$$

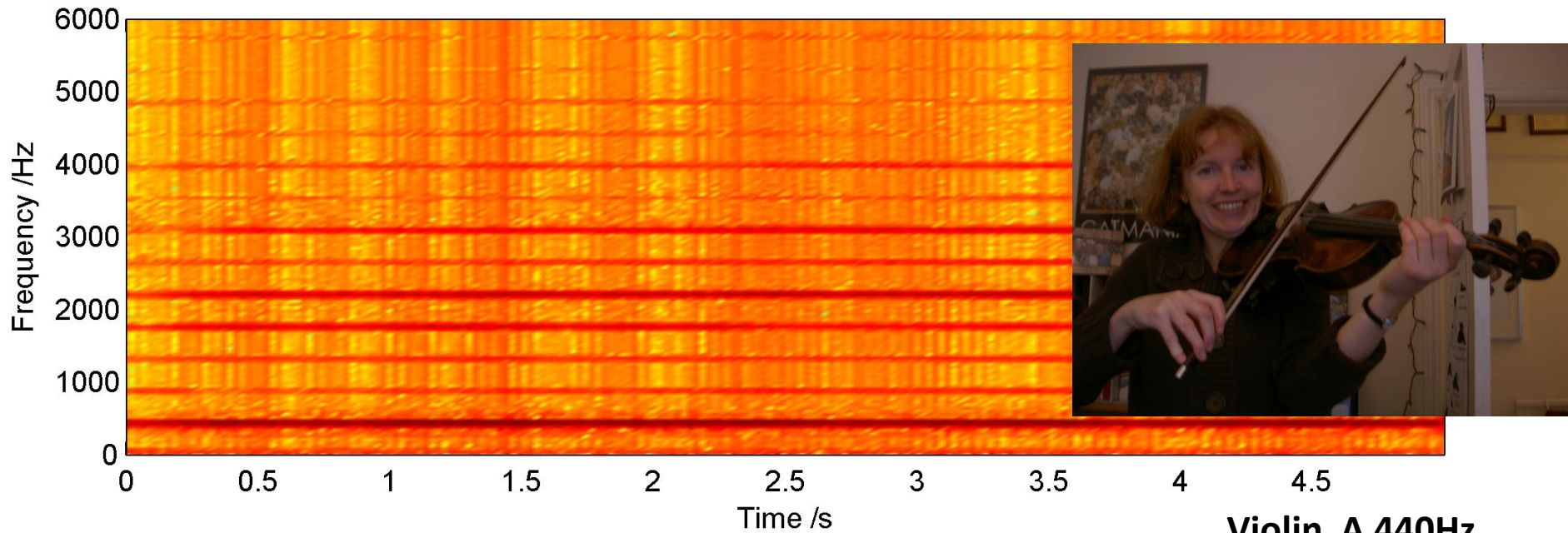
# Musical harmony



- The mathematics of music has been known since the time of Pythagoras, 2500 years ago
- Frequency intervals of simple fractions e.g. 3:2 (a fifth) yield 'harmonious' music
- An **octave** means a **frequency ratio of 2**. An octave above concert A (440Hz) is therefore 880Hz. An octave below is 220Hz.
- The modern 'equal-tempered scale' divides an octave (the frequency ratio 2) into twelve parts such that

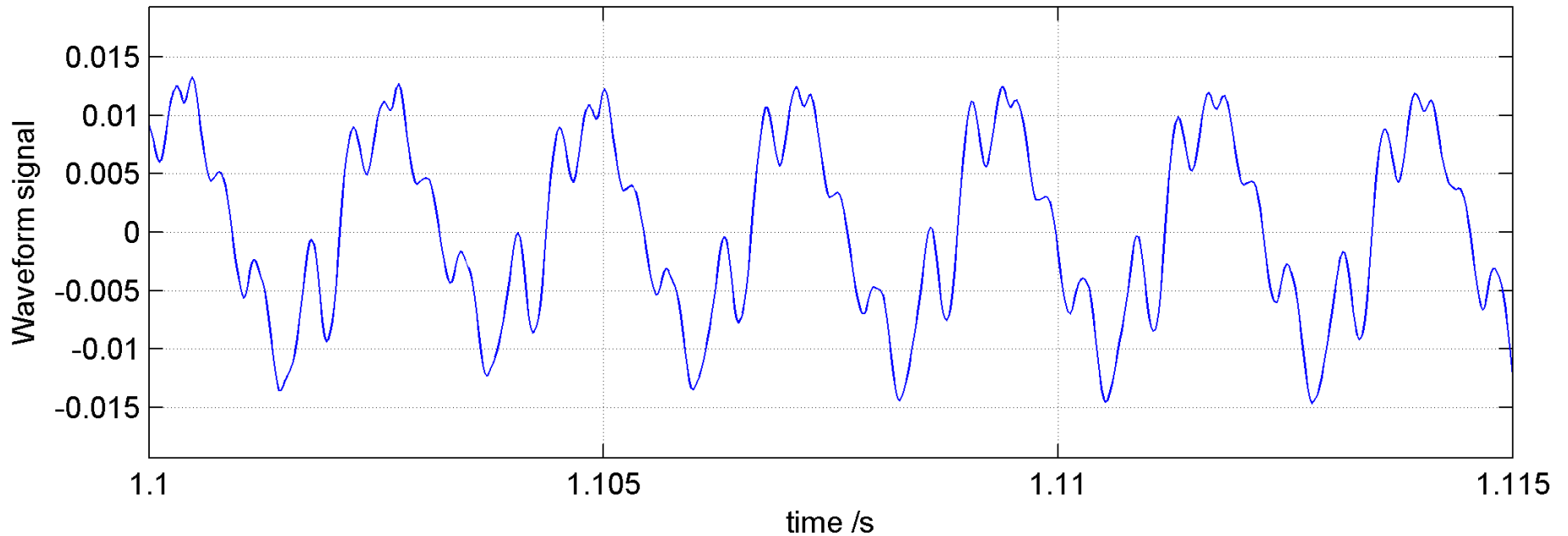
$$f_n = 2^{n/12} = \sqrt[12]{2^n}$$

Normalized spectrogram /dB: Frequency spectrum variation with time



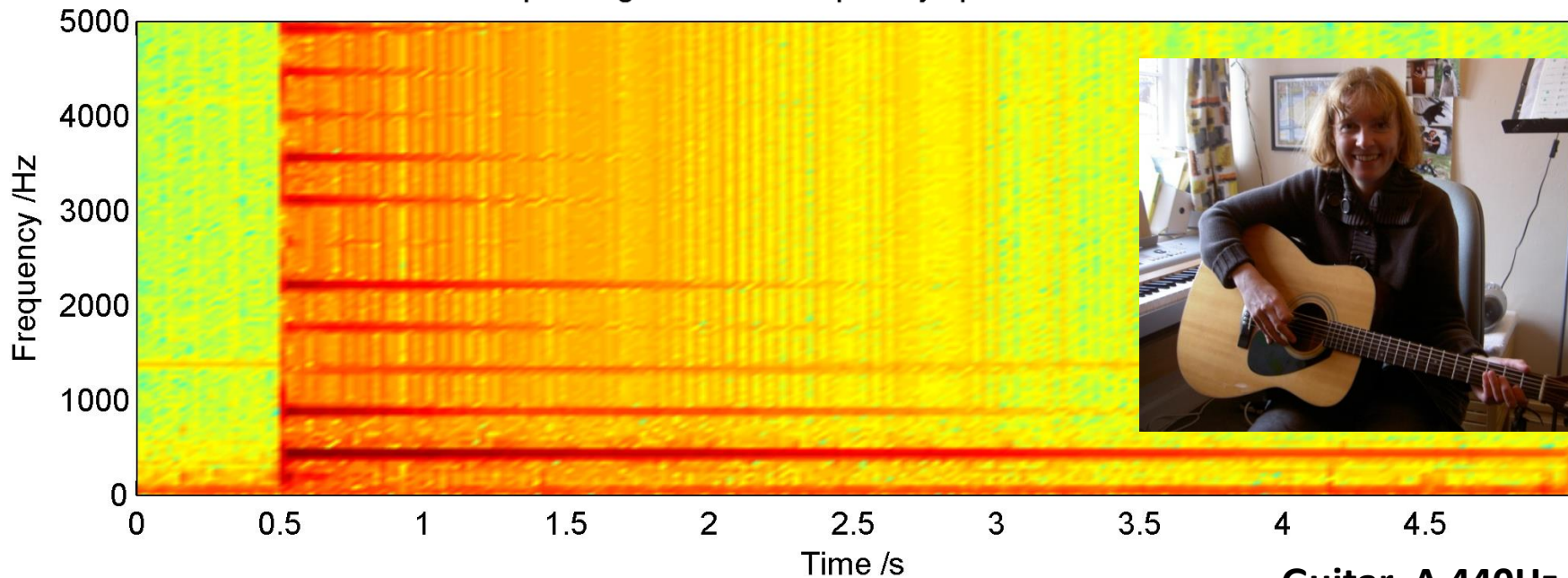
**Violin A 440Hz**

Waveform signal vs time: Right channel



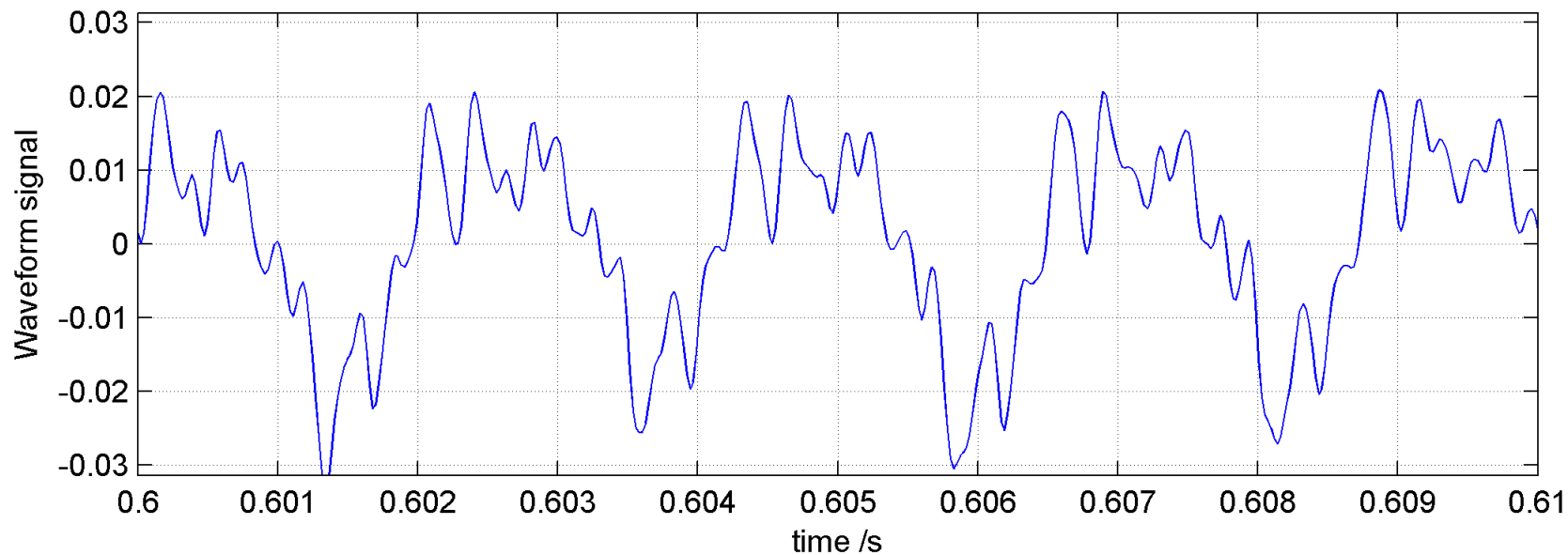


Normalized spectrogram /dB: Frequency spectrum variation with time

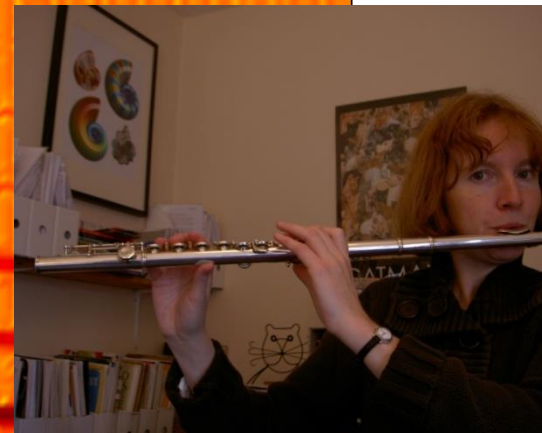
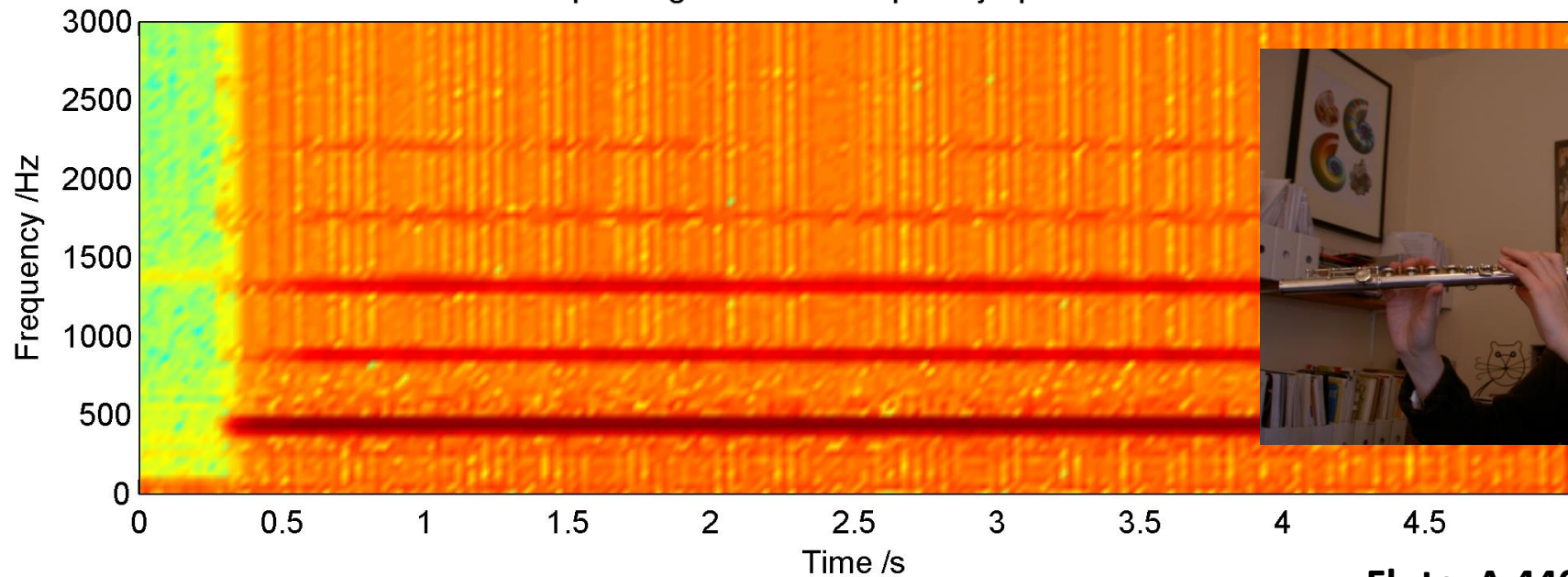


**Guitar A 440Hz**

Waveform signal vs time: Right channel

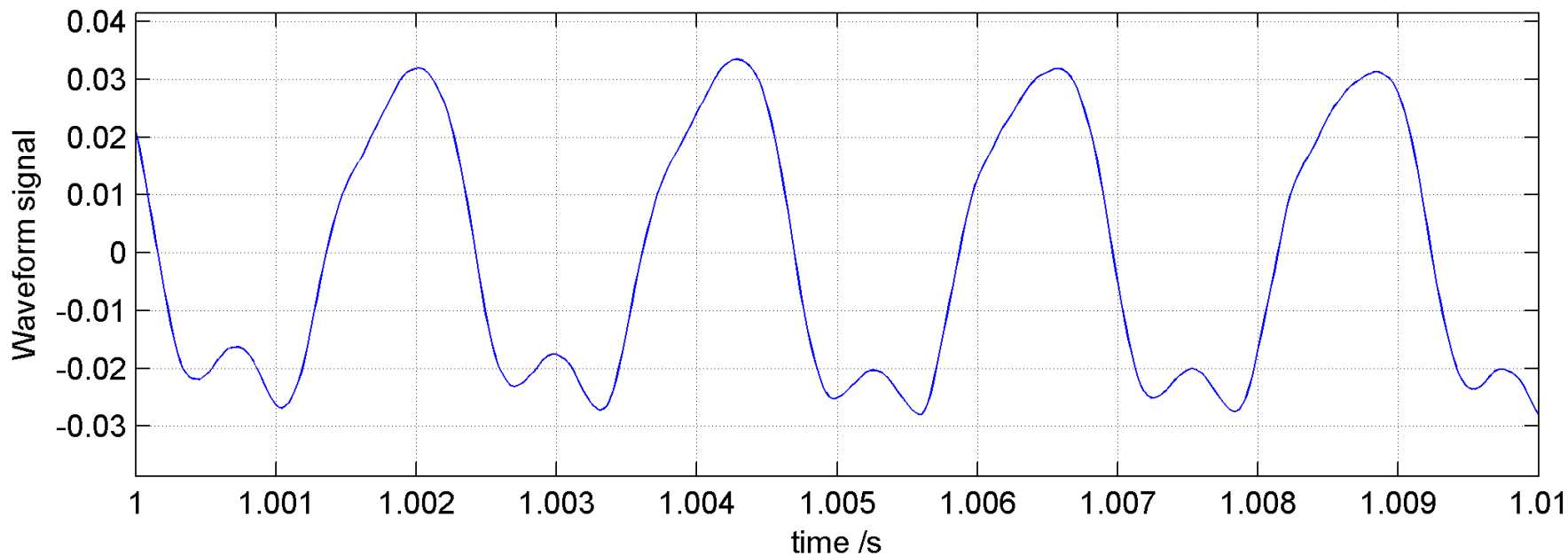


Normalized spectrogram /dB: Frequency spectrum variation with time

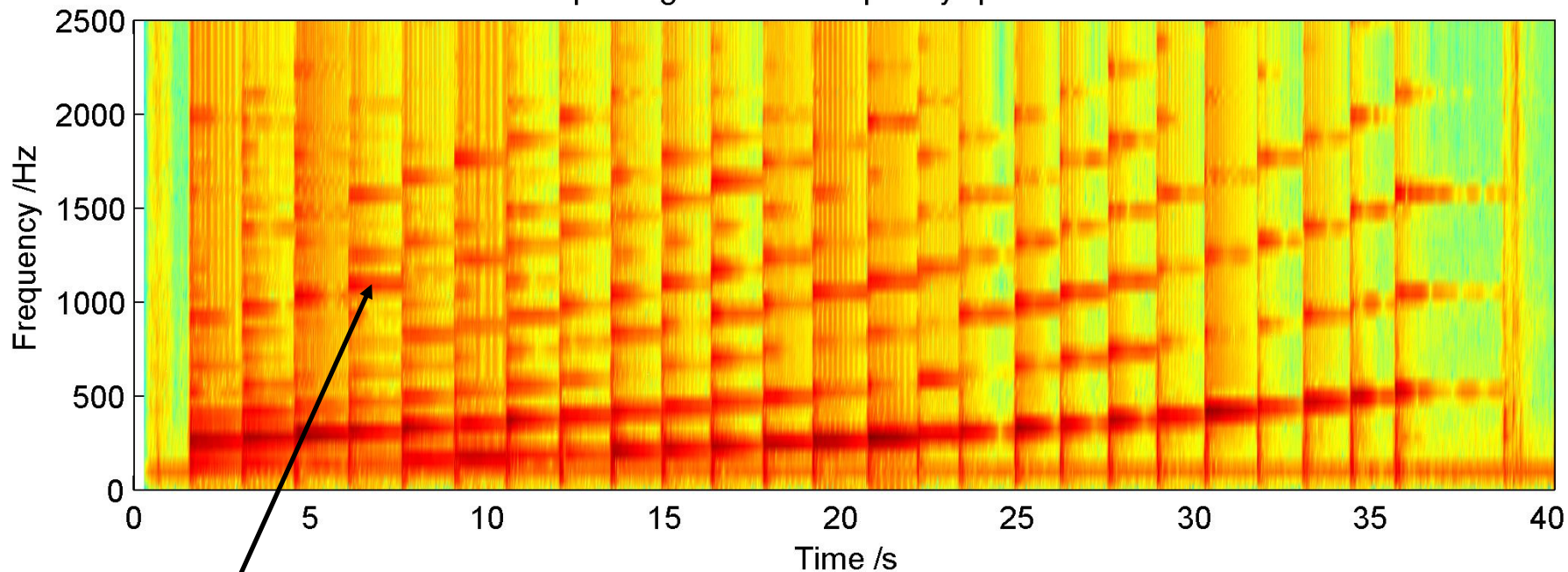


**Flute A 440Hz**

Waveform signal vs time: Right channel

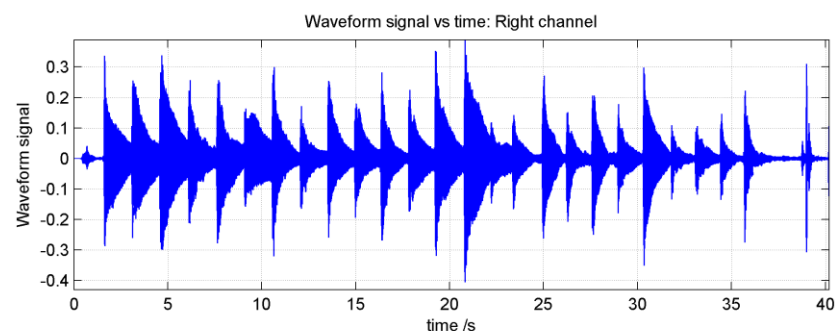
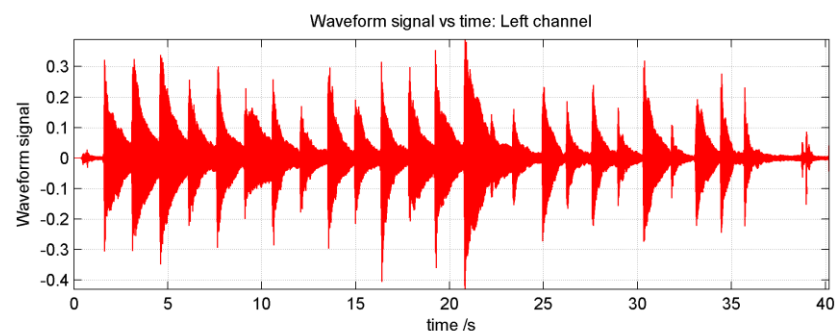


# Normalized spectrogram /dB: Frequency spectrum variation with time



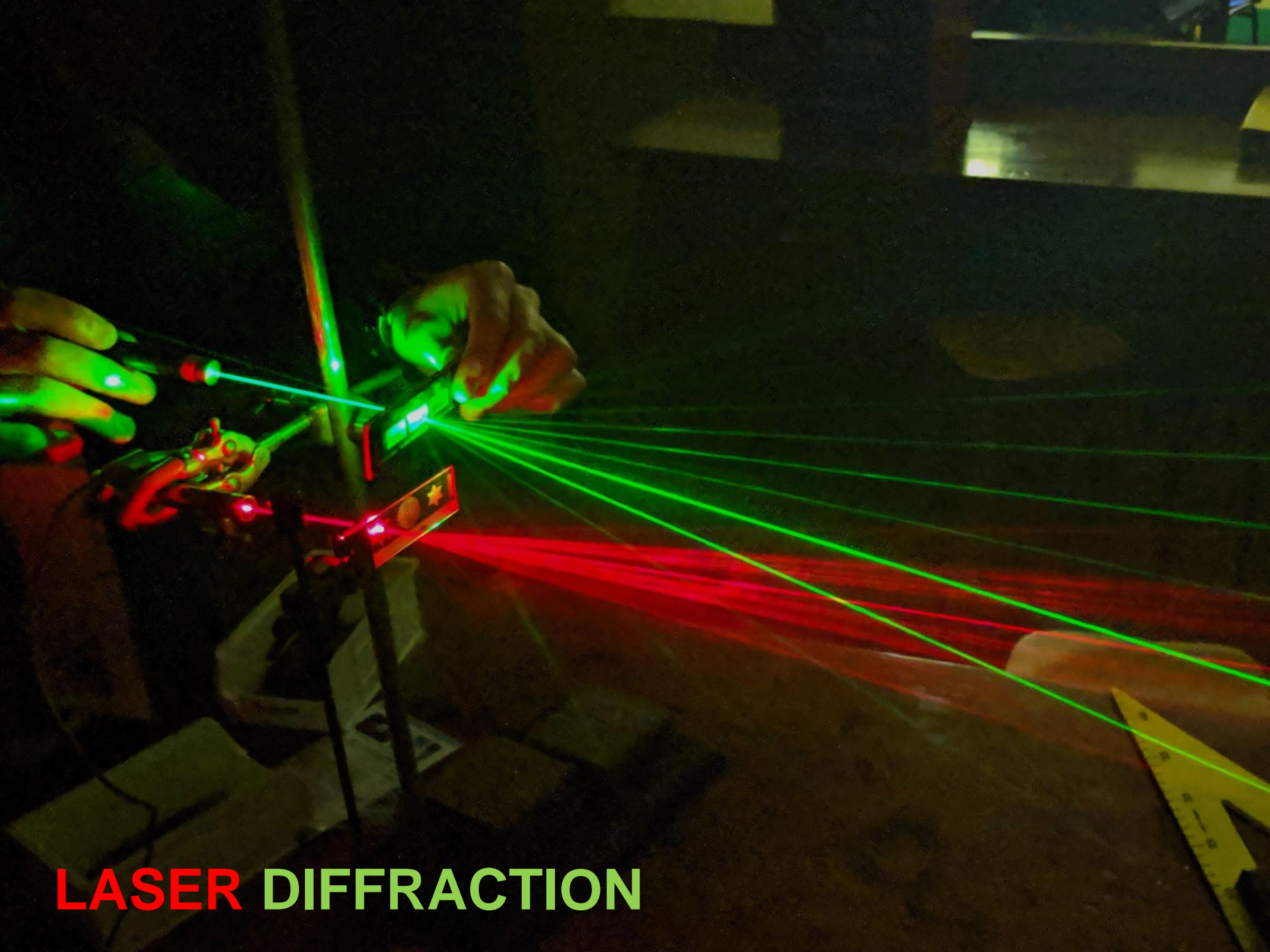
$$f_n = 2^{n/12}$$

Notice  
*power law* of  
harmonics



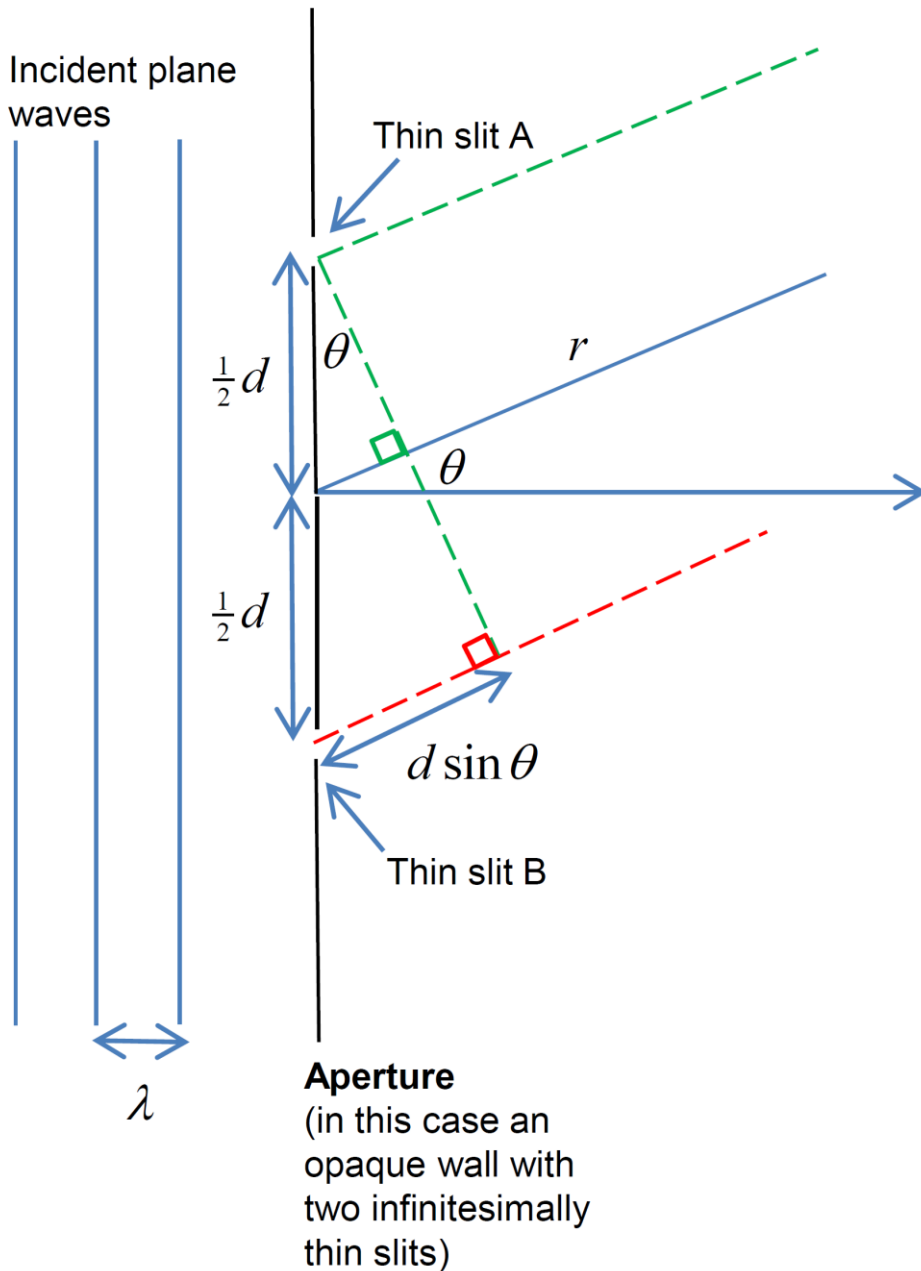
*Chromatic scale* (played on a piano) over two octaves

C,C#...Bb,B,  
C,C#...Bb,B,C



**LASER** DIFFRACTION

# Key geometrical idea from two infinitesimally thin slits ('Young's Slits')



Spherical waves will emanate from the slits, and interfere with each other.

For distances such that:

$$r \gg \frac{d^2}{\lambda}$$

(we call this the **Far Field**) we can assume waves from each slit are **plane waves**, for any given observational angle  $\theta$ .

**Constructive interference** occurs when the *phase difference* between the waves from slits A and B is an integer multiple of  $2\pi$  radians.

$$\frac{2\pi}{\lambda} d \sin \theta = 2\pi n$$

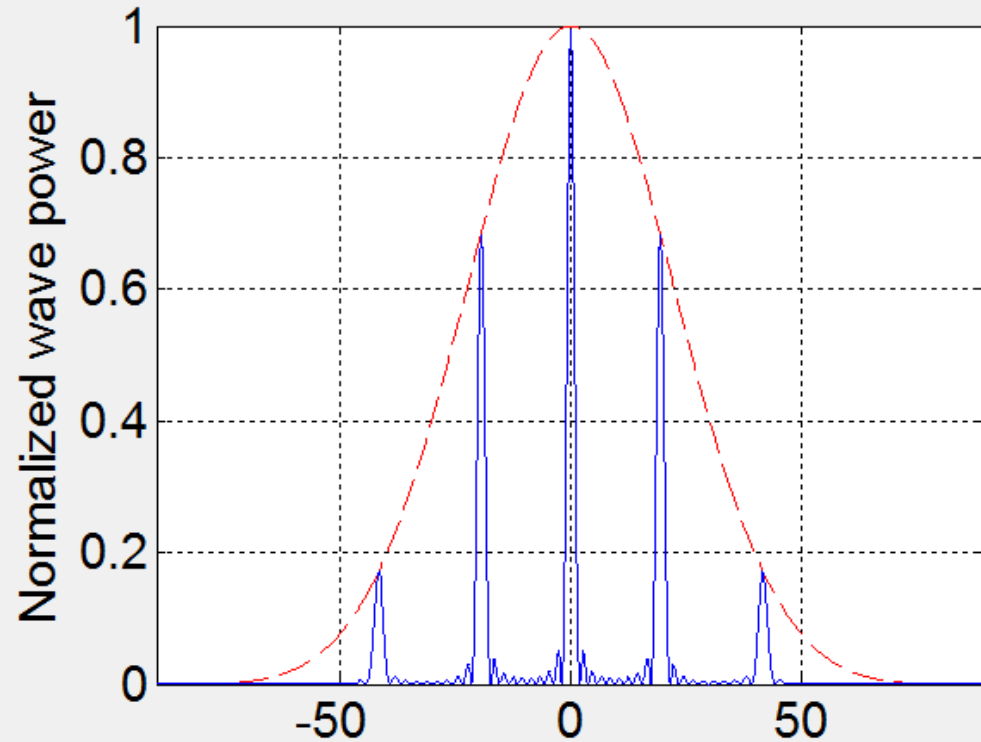
Wavenumber  $k$       Path difference between waves from A and B      Integer  $n$

Hence expect **maxima** in the resulting **Far Field Diffraction pattern** (e.g. spots of a laser on a wall) at angles

$$\sin \theta = \frac{n\lambda}{d}$$

# Grating Fraunhofer far field diffraction

$\lambda = 650\text{nm}$ ,  $s = 3\lambda$ ,  $w = 1\lambda$ ,  $N = 10$



Wavelength /nm  $\theta^\circ$

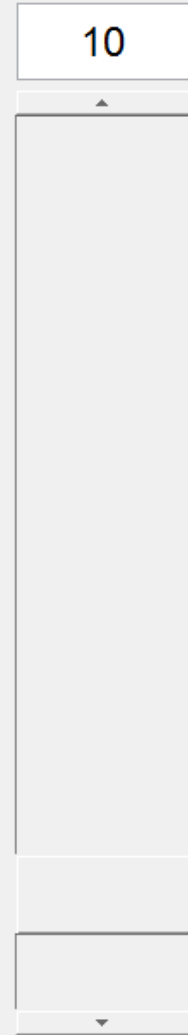
Speed of light /ms<sup>-1</sup>

Frequency /THz

Number of slits

Slit spacing /wavelength

Slit width /wavelength

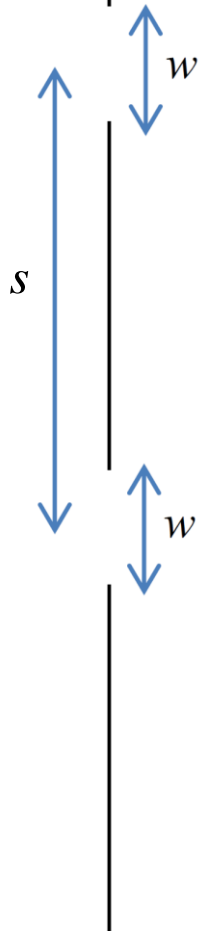


Save .PNG

dB

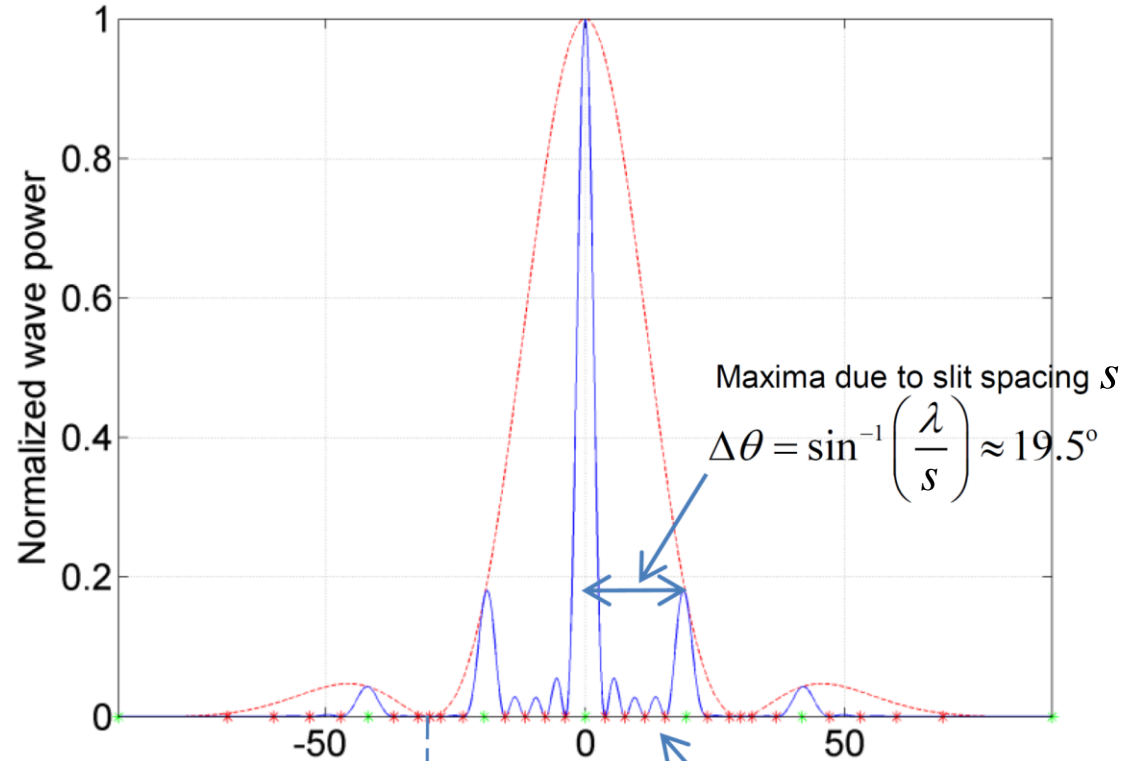
$$|\psi|^2 = \frac{A^2}{N^2 r^2} \left( \frac{\sin\left(\frac{\pi}{\lambda} w \sin \theta\right)}{\frac{\pi}{\lambda} w \sin \theta} \times \frac{\sin\left(\frac{\pi}{\lambda} N s \sin \theta\right)}{\sin\left(\frac{\pi}{\lambda} s \sin \theta\right)} \right)^2$$

Slits of width  
and spacing  $s$



Grating Fraunhofer far field diffraction

$\lambda = 650\text{nm}$ ,  $d = 3\lambda$ ,  $w = 2\lambda$ ,  $N = 5$



**Envelope due to finite slit width**

Zeros at:  $\theta = \sin^{-1}\left(\frac{n\lambda}{w}\right)$ ;  $n \neq 0$

Maxima due to slit spacing  $S$

$\Delta\theta = \sin^{-1}\left(\frac{\lambda}{s}\right) \approx 19.5^\circ$

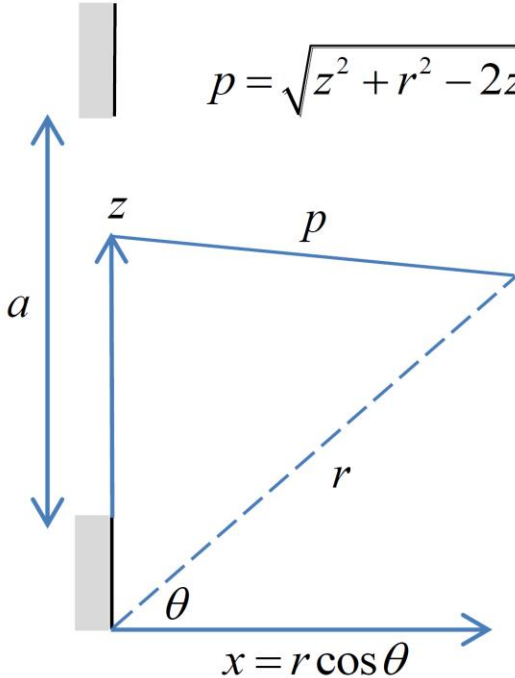
**Fine structure due to number of slits**  
(i.e. overall size of aperture)

Zeros at:  $\theta = \sin^{-1}\left(\frac{p\lambda}{Ns}\right)$

But maxima when  $\frac{p}{N}$  integer  $m$

# Modelling general diffraction effects from a finite width slit

We can use a computer to evaluate the wavefield in the vicinity of a finite width slit which is uniformly illuminated. We are therefore not restricted to the limitations of the Fraunhofer and Fresnel regimes

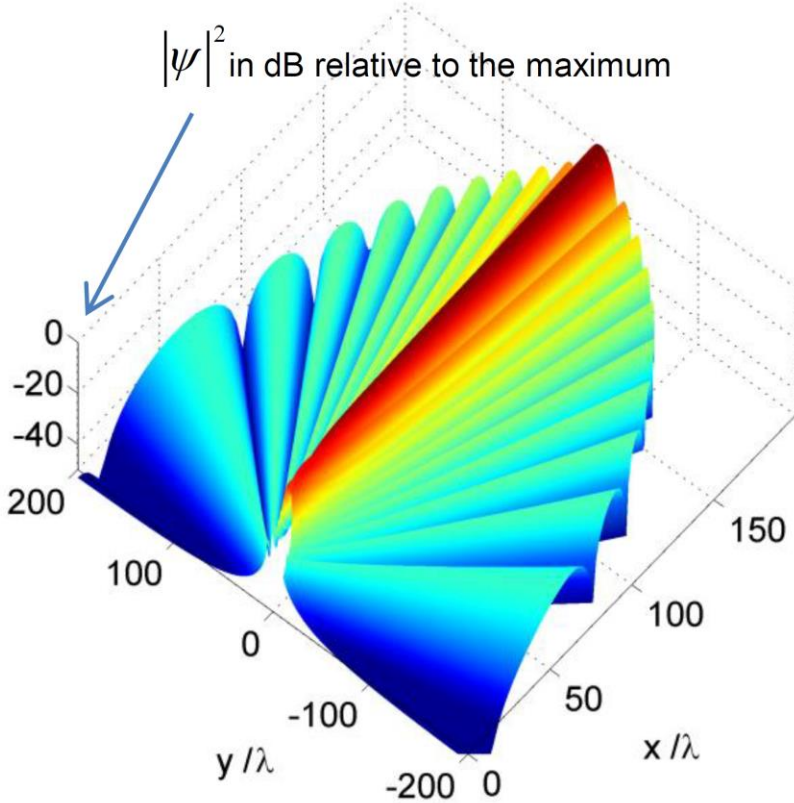


$$p = \sqrt{z^2 + r^2 - 2zr \cos\left(\frac{1}{2}\pi - \theta\right)} = \sqrt{z^2 + r^2 - 2zr \sin \theta}$$

$$\psi(r, t) = \frac{A}{a} e^{-i\omega t} \int \frac{e^{ikp}}{p} dz$$

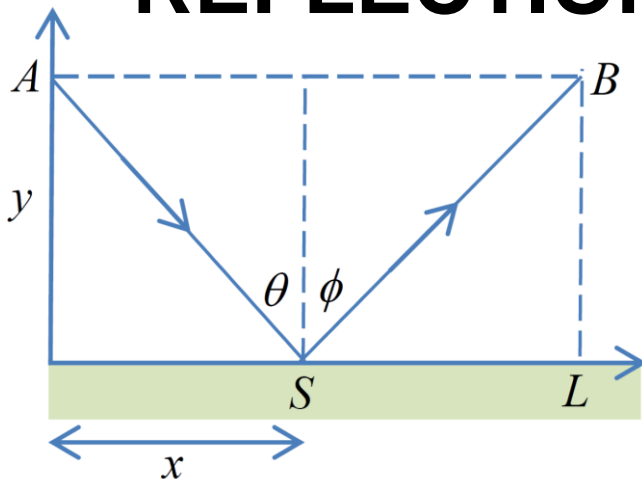
The idea is to evaluate the wavefield power at various  $r$  and  $\theta$  values. The integral is approximated by a sum based on a large number of finite  $dz$  values

In the examples below :  
 $dz = \frac{1}{500} a$





# REFLECTION



$$x = y \tan \theta$$

$$L - x = Y \tan \phi$$

$$t = \frac{\sqrt{x^2 + y^2}}{c/n} + \frac{\sqrt{(L-x)^2 + y^2}}{c/n}$$

$$\frac{\partial t}{\partial x} = \frac{n}{c} \left( \frac{\frac{1}{2} 2x}{\sqrt{x^2 + y^2}} + \frac{\frac{1}{2} 2(L-x)(-1)}{\sqrt{(L-x)^2 + y^2}} \right) \quad \frac{1}{\cos^2 \theta} = 1 + \tan^2 \theta$$

$$\tan \theta = \frac{\sin \theta}{\cos \theta}$$

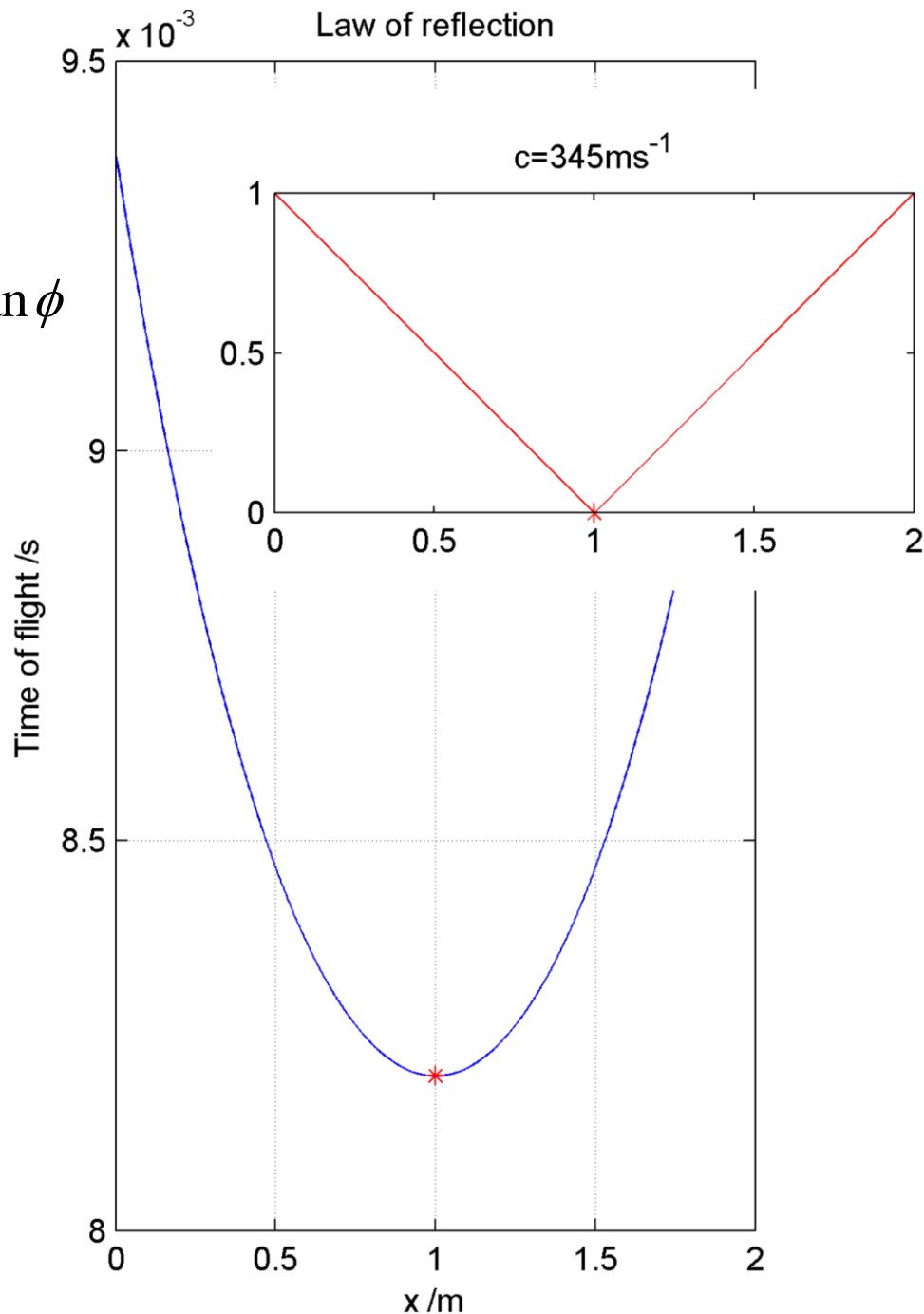
$$\therefore \frac{\partial t}{\partial x} = \frac{n}{c} \left( \frac{y \tan \theta}{\sqrt{y^2 \tan^2 \theta + y^2}} - \frac{y \tan \phi}{\sqrt{y^2 \tan^2 \phi + y^2}} \right)$$

$$\therefore \frac{\partial t}{\partial x} = \frac{n}{c} (\sin \theta - \sin \phi)$$

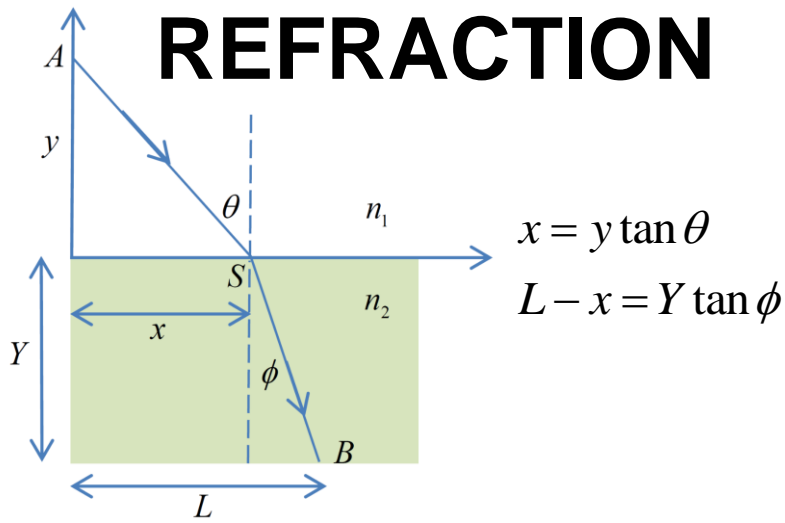
The travel time is minimized when  $\frac{\partial t}{\partial x} = 0$

$$\therefore \sin \theta = \sin \phi$$

$$\therefore \theta = \phi$$



# REFRACTION



$$t = \frac{\sqrt{x^2 + y^2}}{c/n_1} + \frac{\sqrt{(L-x)^2 + Y^2}}{c/n_2}$$

$$\frac{\partial t}{\partial x} = \frac{1}{c} \left( \frac{\frac{1}{2} 2xn_1}{\sqrt{x^2 + y^2}} + \frac{\frac{1}{2} 2(L-x)(-1)n_2}{\sqrt{(L-x)^2 + Y^2}} \right)$$

$$\frac{\partial t}{\partial x} = \frac{1}{c} \left( \frac{y \tan \theta n_1}{y \sqrt{\tan^2 \theta + 1}} + \frac{-Y \tan \phi n_2}{Y \sqrt{\tan^2 \phi + 1}} \right)$$

$$\frac{\partial t}{\partial x} = \frac{1}{c} (\cos \theta \tan \theta n_1 - \cos \phi \tan \phi n_2)$$

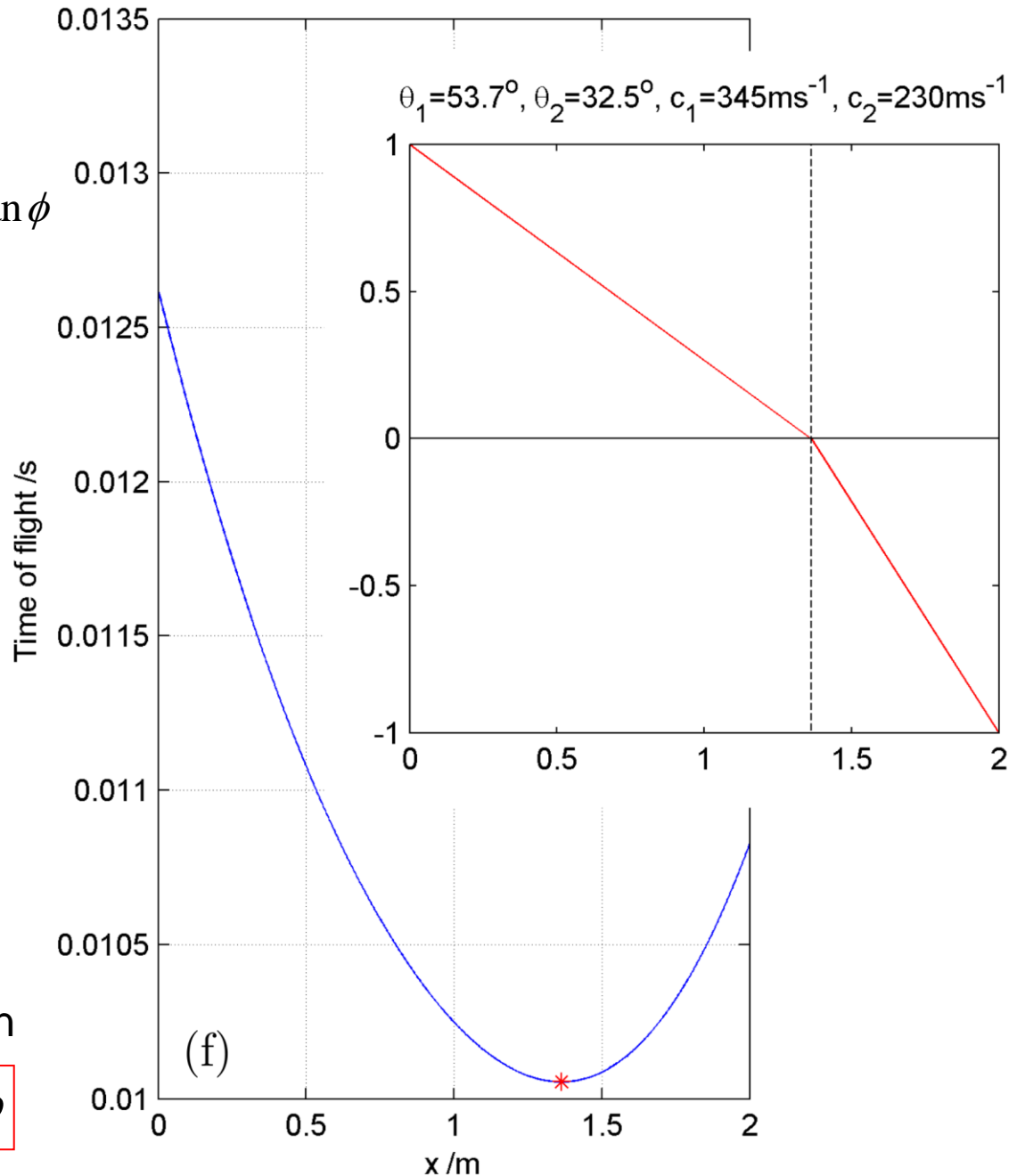
$$\frac{\partial t}{\partial x} = \frac{1}{c} (n_1 \sin \theta - n_2 \sin \phi)$$

The travel time is minimized when

$$\frac{\partial t}{\partial x} = 0$$

$$\therefore n_1 \sin \theta = n_2 \sin \phi$$

Refraction:  $\sin(\theta_1)/c_1 = 0.00234$ ,  $\sin(\theta_2)/c_2 = 0.00233$



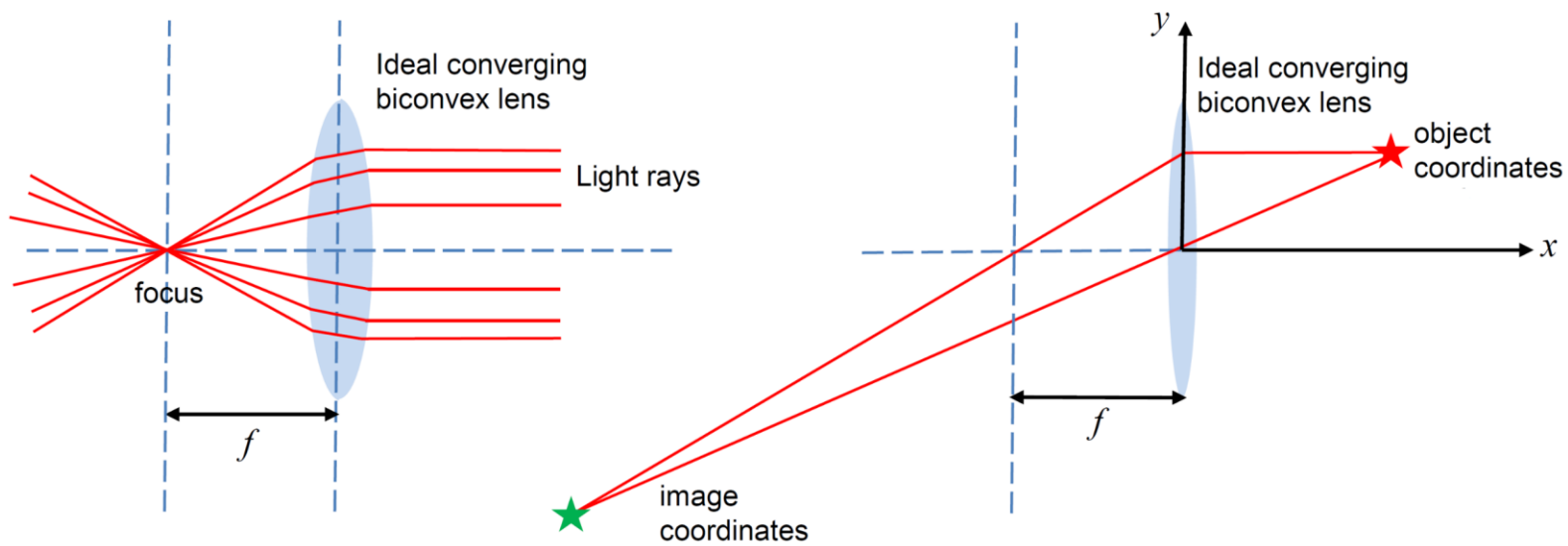
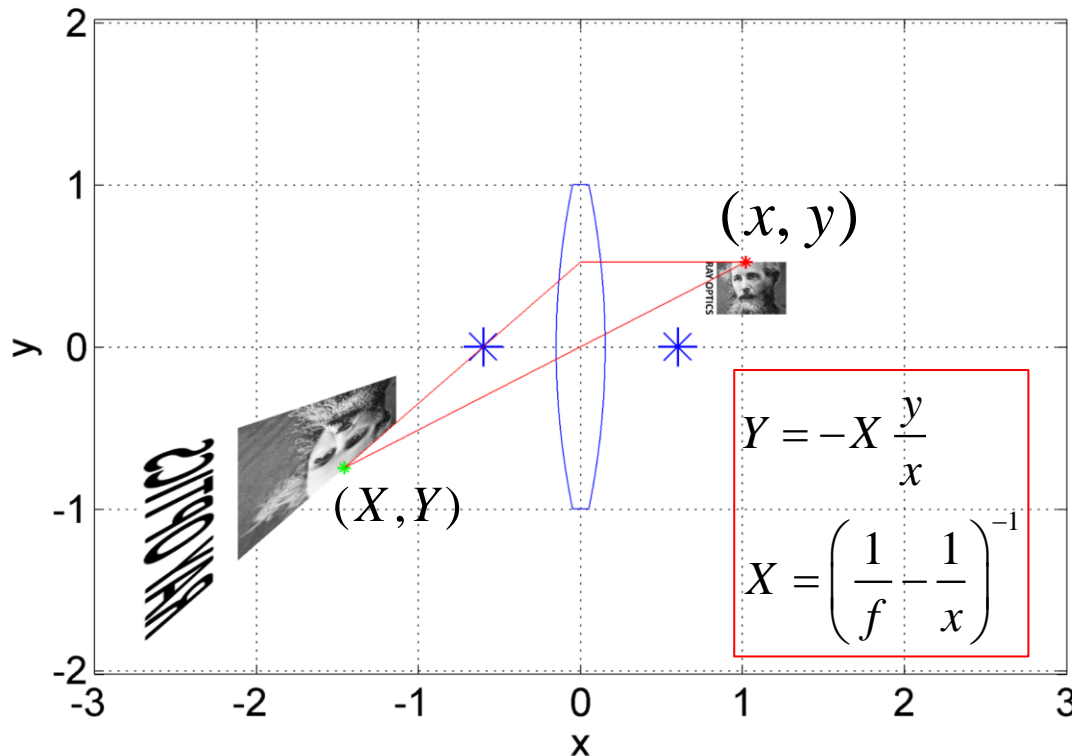
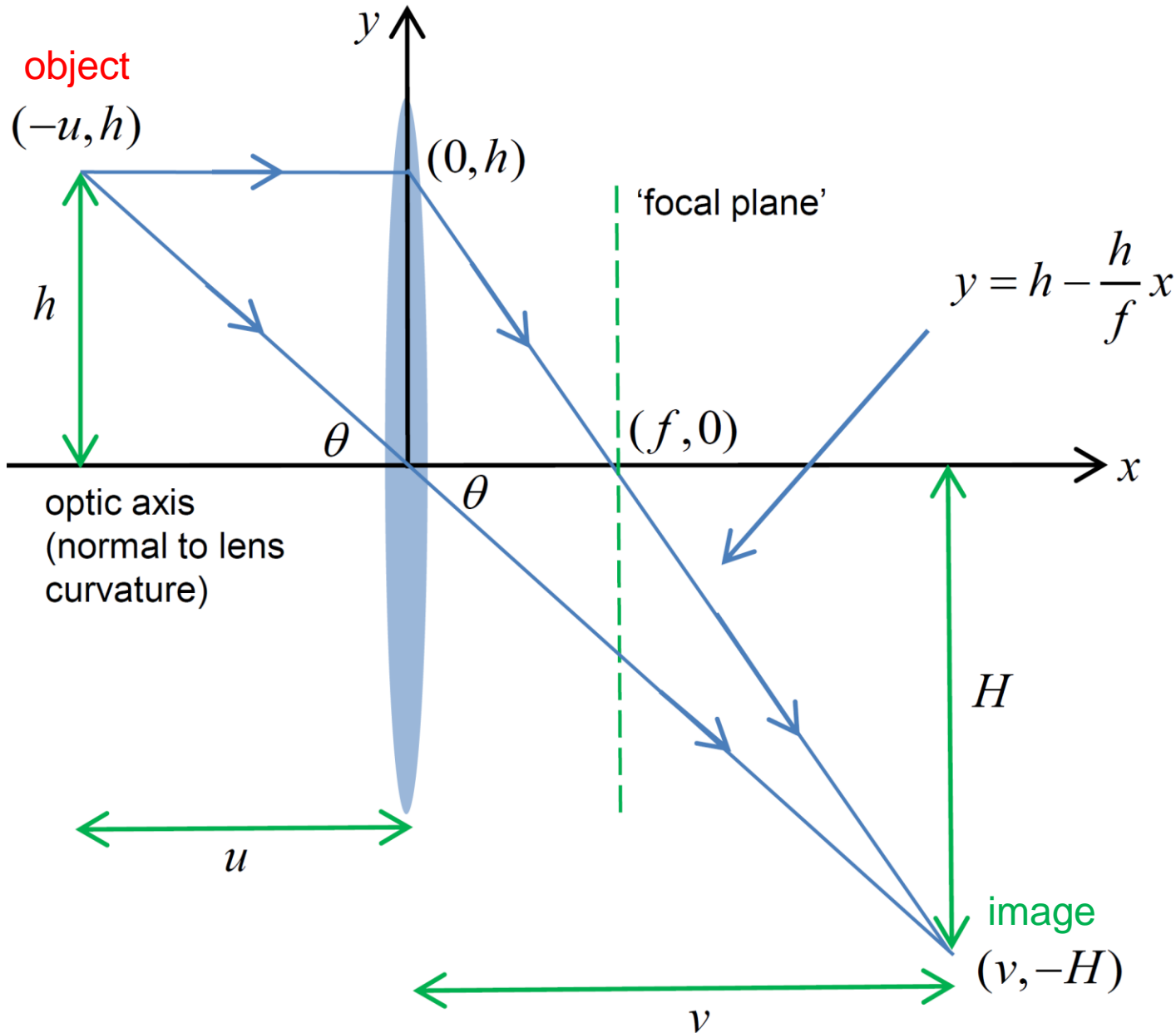


Image of object in a converging lens



# RAY OPTICS: THIN BICONVEX LENS

# THIN LENS EQUATION



$$\frac{1}{u} + \frac{1}{v} = \frac{1}{f}$$

$$\frac{H}{v} = \frac{h}{u}$$

$$v = \left( \frac{1}{f} - \frac{1}{u} \right)^{-1}$$

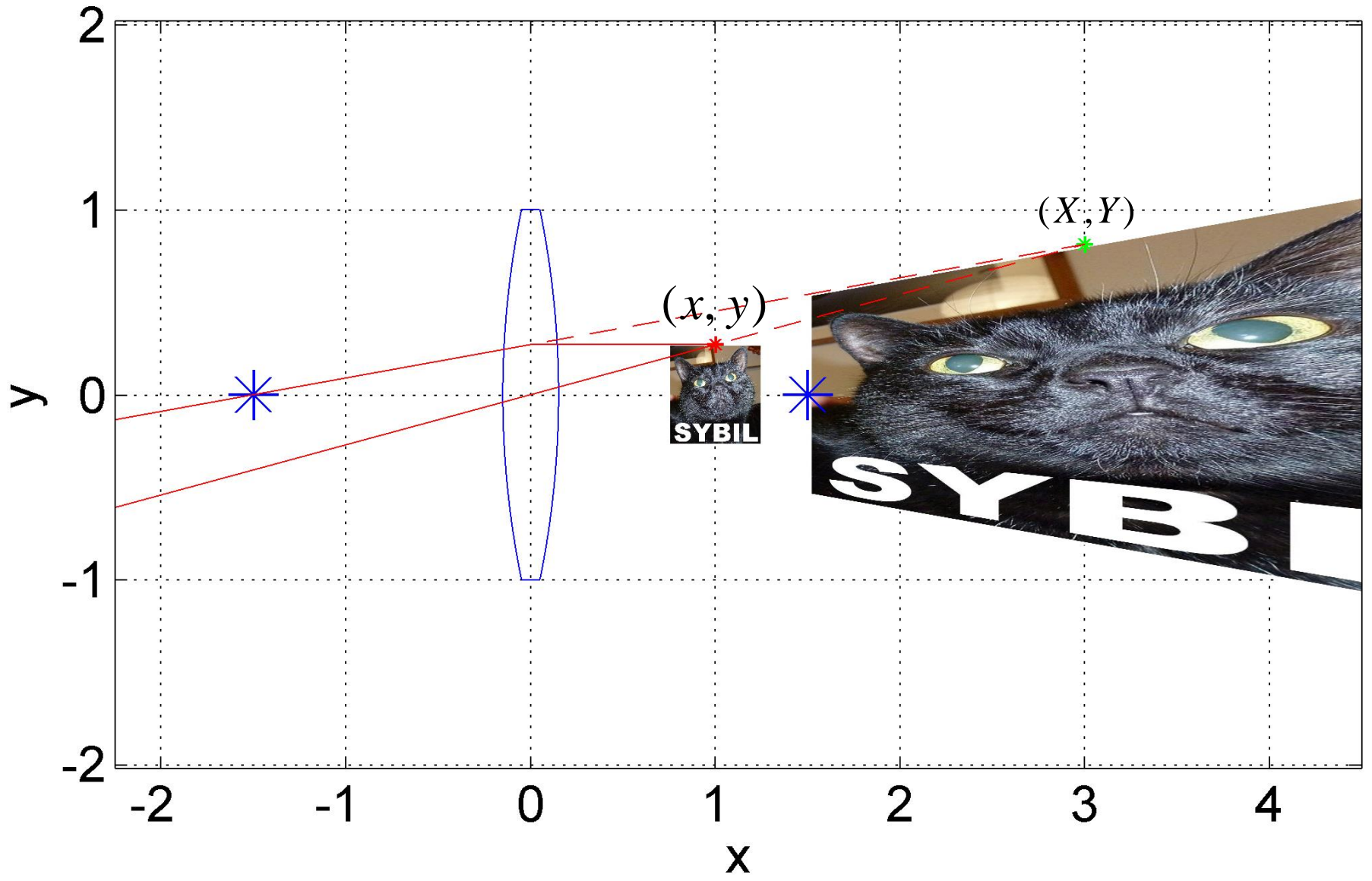


$$Y = -X \frac{y}{x}$$

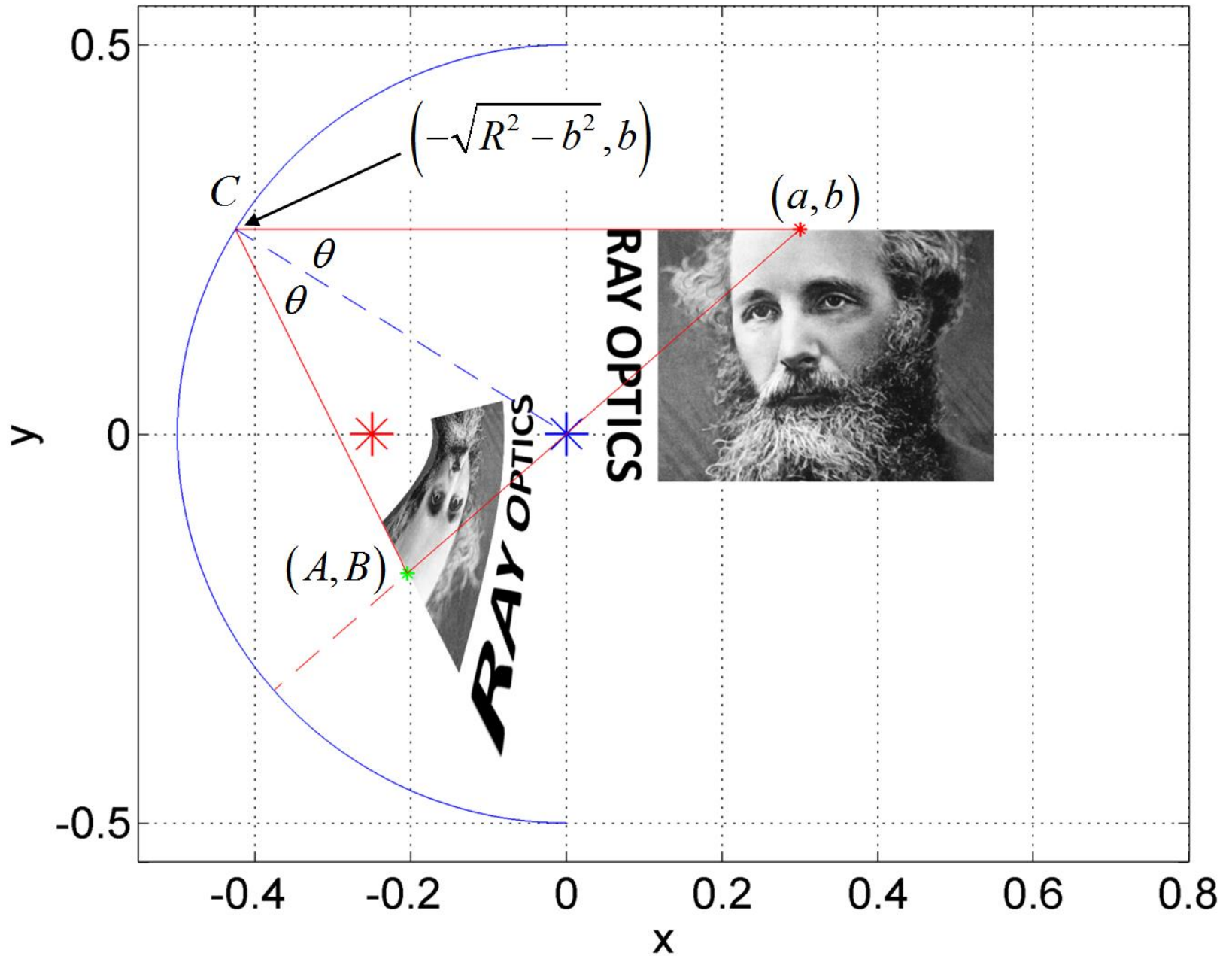
$$X = \left( \frac{1}{f} - \frac{1}{x} \right)^{-1}$$

Image is a  
coordinate  
transformation

# Virtual image of object in a magnifying lens



# Reflection in a concave mirror



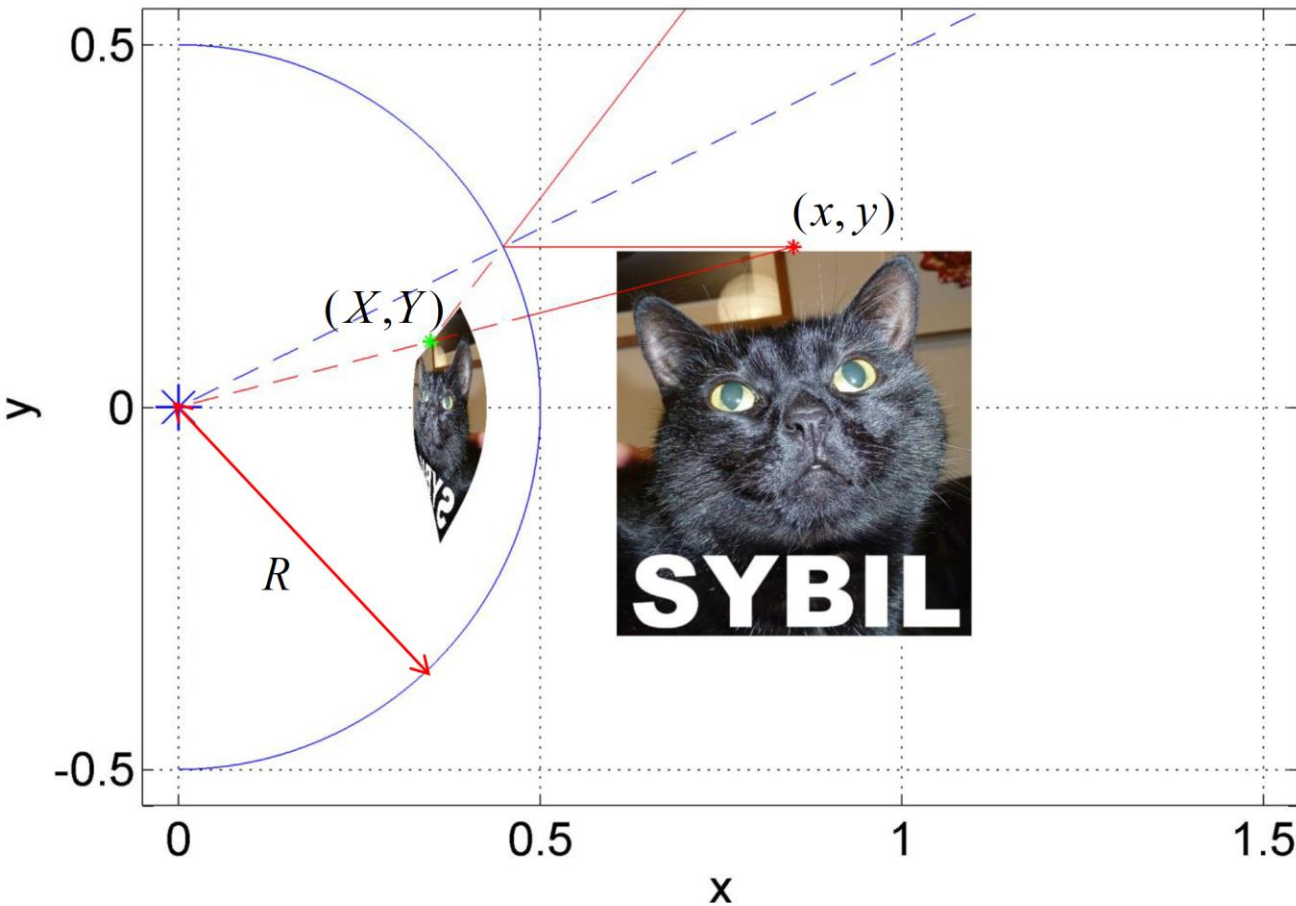
(a)



(b)



## Reflection in a convex mirror



$$\alpha = \frac{1}{2} \tan^{-1} \left( \frac{y}{x} \right)$$

$$k = \frac{x}{\cos(2\alpha)}$$

$$Y = \frac{k \sin \alpha}{\frac{k}{R} - \cos \alpha + \frac{x}{y} \sin \alpha}$$

$$X = x \frac{Y}{y}$$

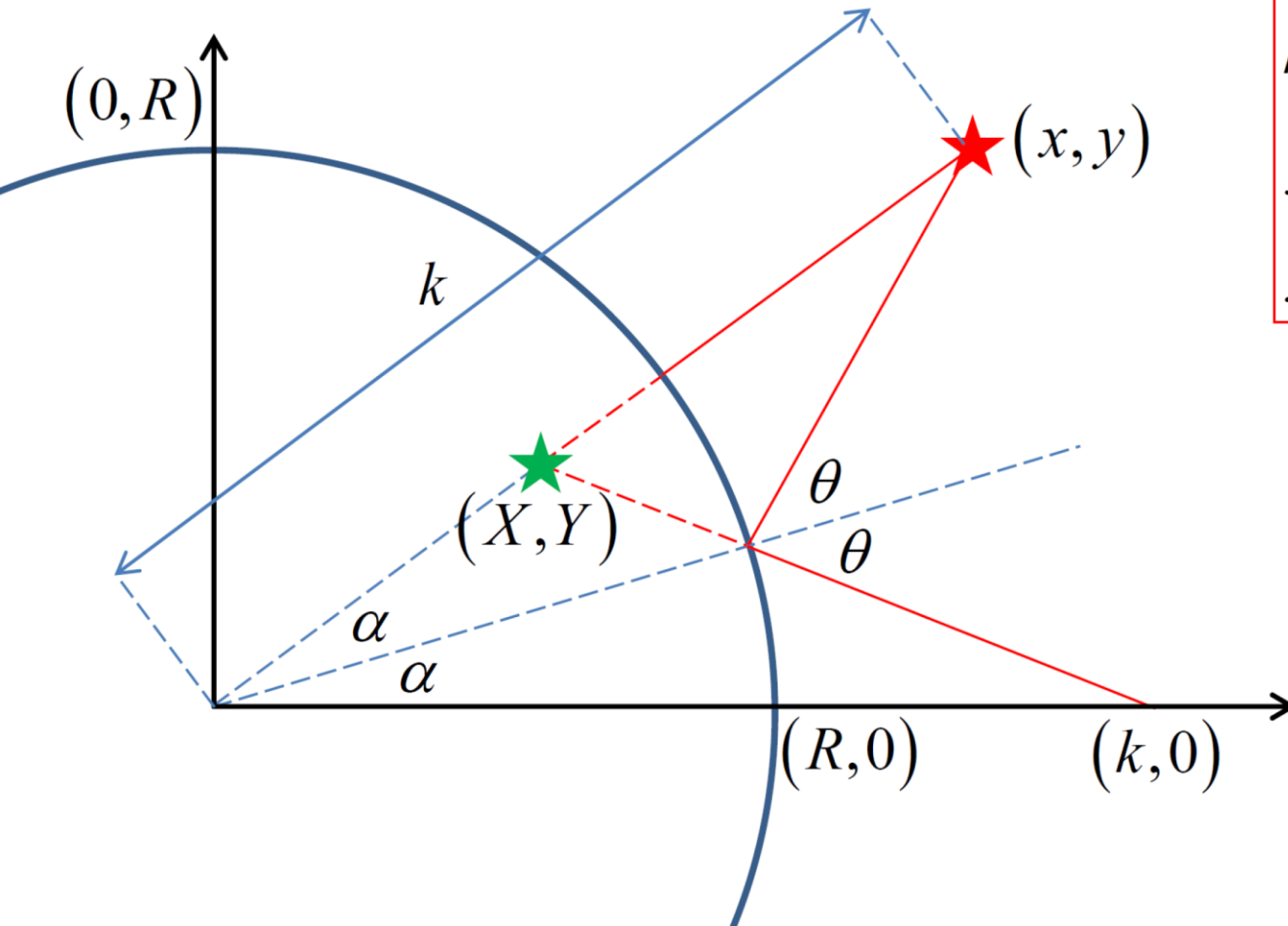
Virtual image from object coordinates

We see an upright, distorted *virtual image* in a cylindrical mirror.

i.e. the *apparent source* of (diverging) light rays from the mirror



# Convex mirror object to virtual image transformation (and reverse transformation)



$$\alpha = \frac{1}{2} \tan^{-1} \left( \frac{y}{x} \right)$$

$$k = \frac{R(Y \cos \alpha - X \sin \alpha)}{Y - R \sin \alpha}$$

$$x = k \cos(2\alpha)$$

$$y = k \sin(2\alpha)$$

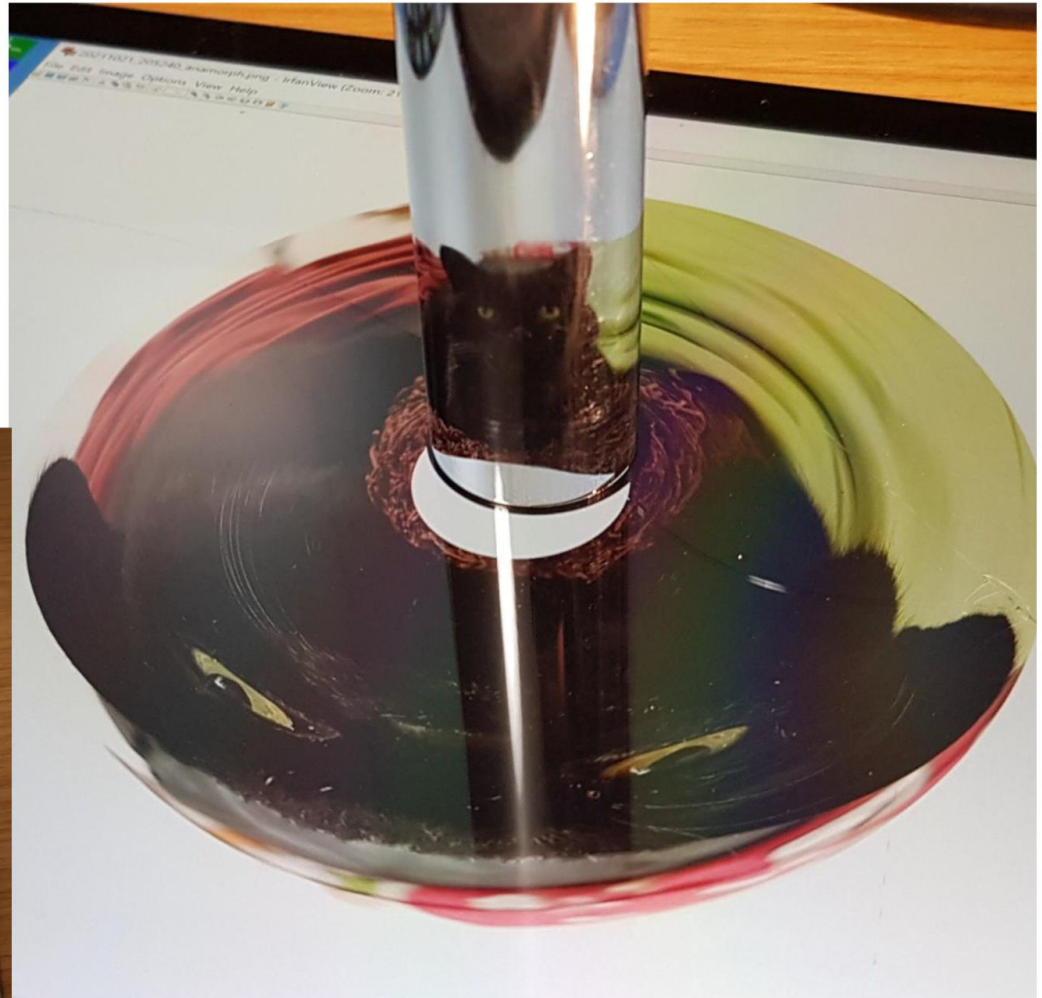
$$\alpha = \frac{1}{2} \tan^{-1} \left( \frac{y}{x} \right)$$

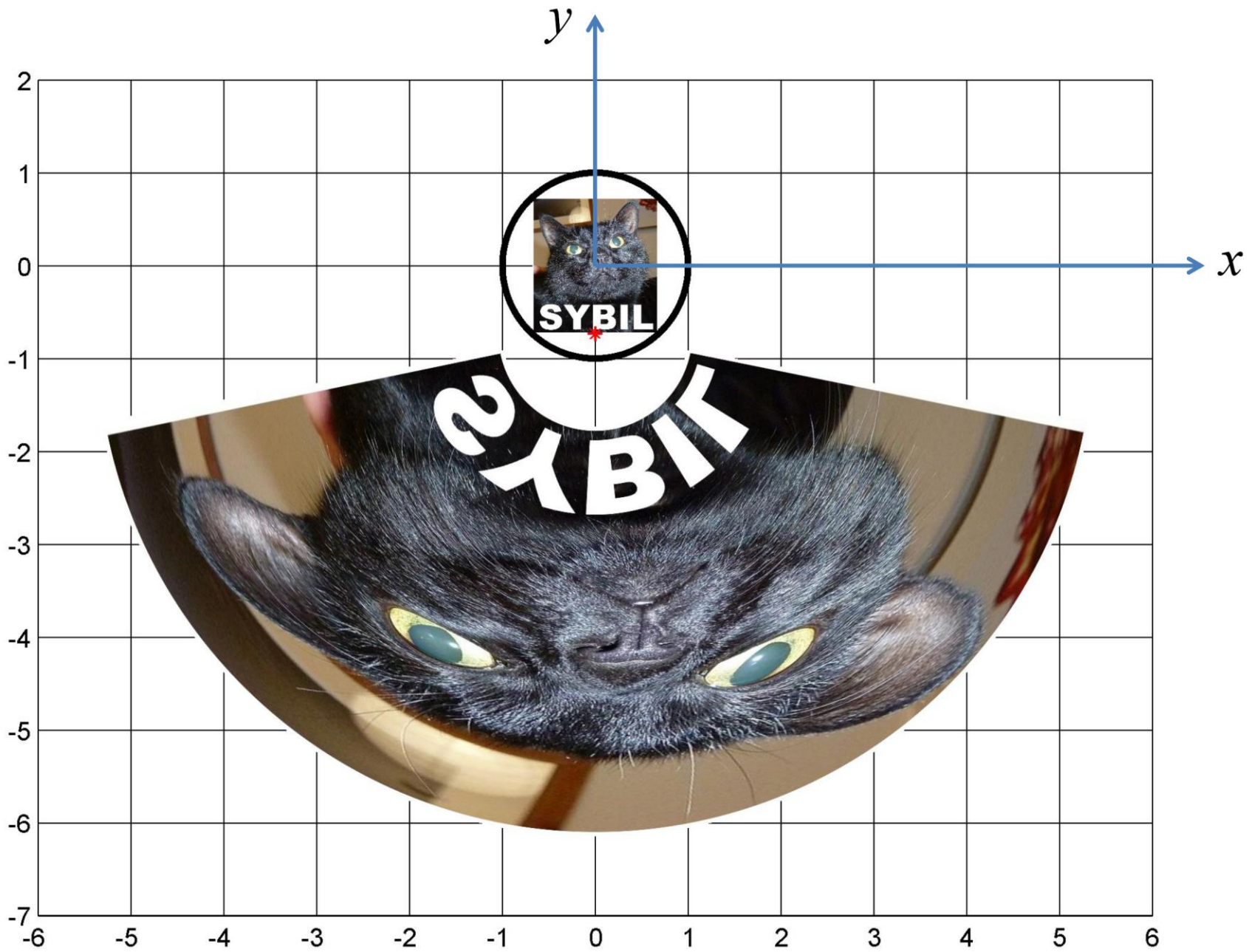
$$k = \frac{x}{\cos(2\alpha)}$$

$$Y = \frac{k \sin \alpha}{\frac{k}{R} - \cos \alpha + \frac{x}{y} \sin \alpha}$$

$$X = x \frac{Y}{y}$$

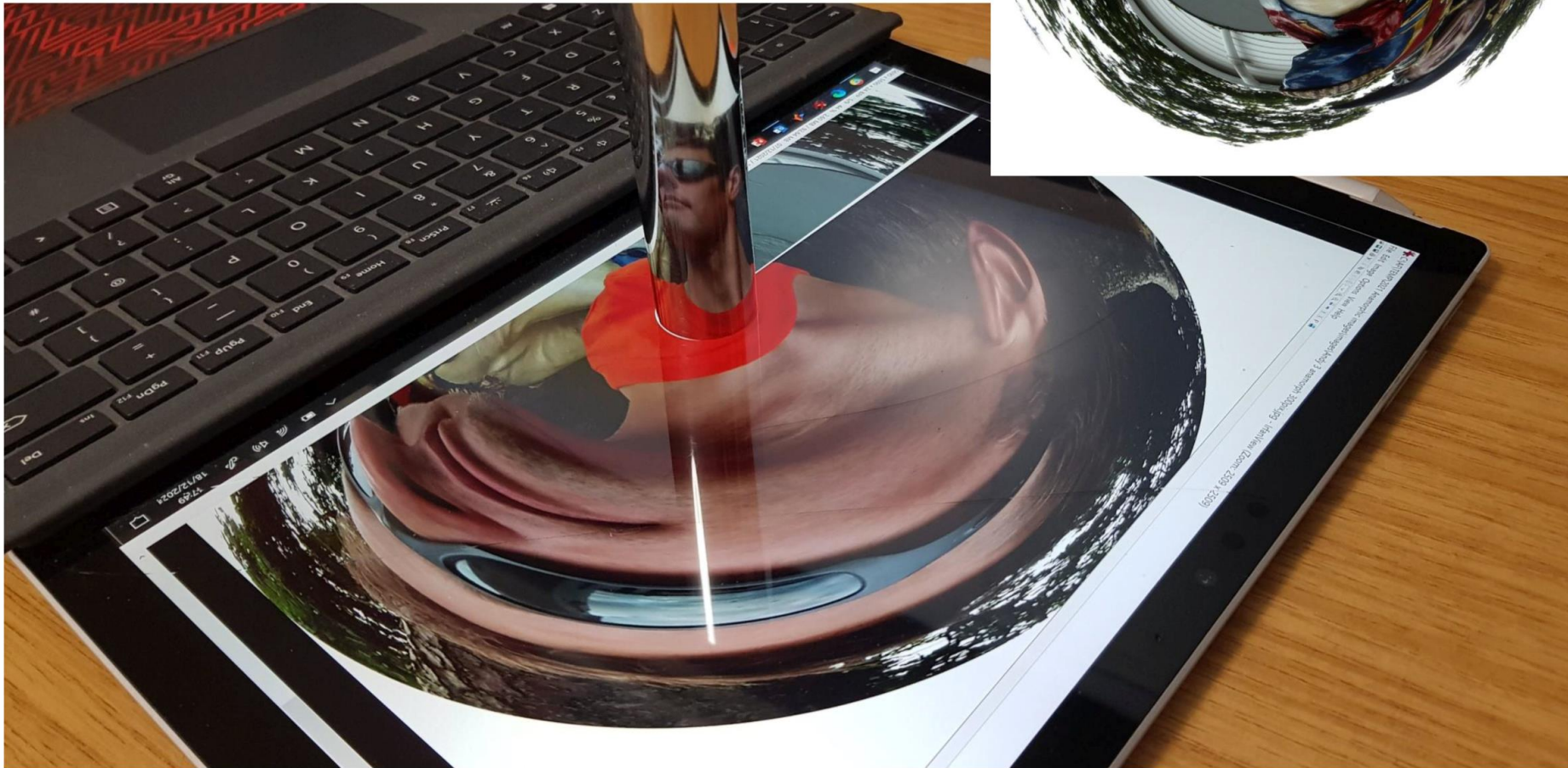
Sybil the cat was unperturbed by this anamorphic transformation.





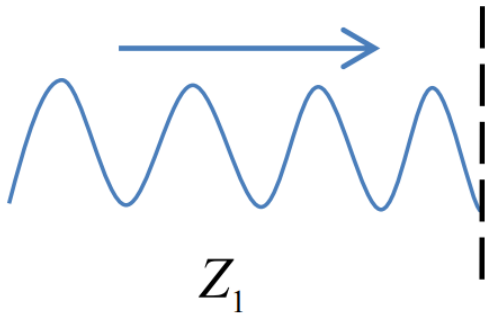
arc\_deg = 160      Rf = 3

'Fear and loathing' in Portmeirion ...



# Reflection and transmission of waves on boundaries

$$Z = \rho c \quad \text{acoustic impedance}$$



**Incident wave**

$$\psi_I = A_I e^{i(k_1 x - \omega t)}$$

$Z_2$

$$\frac{A_R}{A_I} = \frac{Z_1 - Z_2}{Z_2 + Z_1}$$

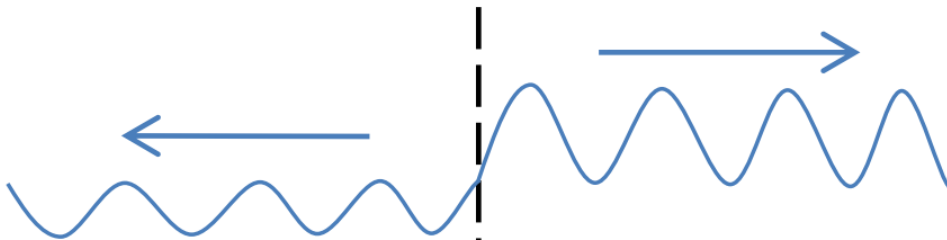
$$\frac{A_T}{A_I} = \frac{2Z_1}{Z_2 + Z_1}$$

$$P = \frac{1}{2} Z A^2 \omega^2$$

Wave power

$$\frac{P_R}{P_I} = \left| \frac{Z_1 - Z_2}{Z_1 + Z_2} \right|^2$$

$$\frac{P_T}{P_I} = 1 - \left| \frac{Z_1 - Z_2}{Z_1 + Z_2} \right|^2$$



**Reflected wave**

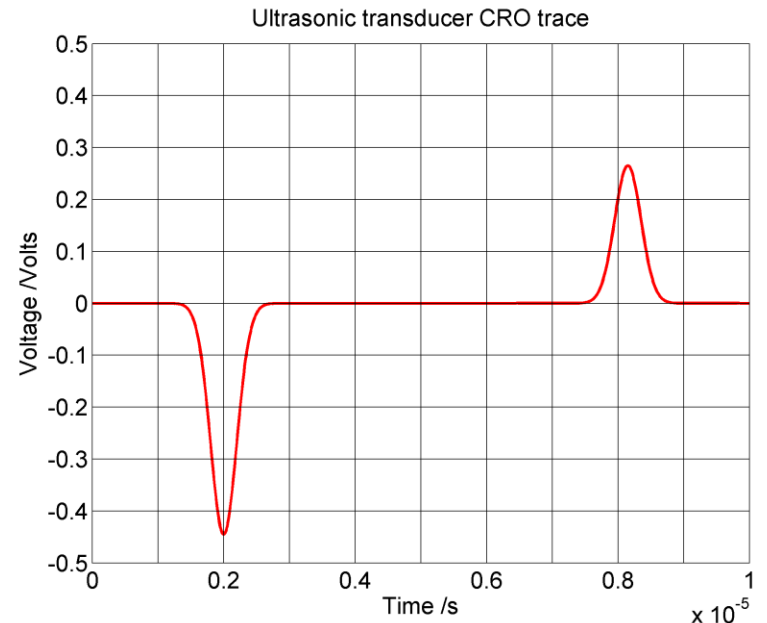
$$\psi_R = A_R e^{i(-k_1 x - \omega t)}$$

$Z_1$

**Transmitted wave**

$$\psi_T = A_T e^{i(k_2 x - \omega t)}$$

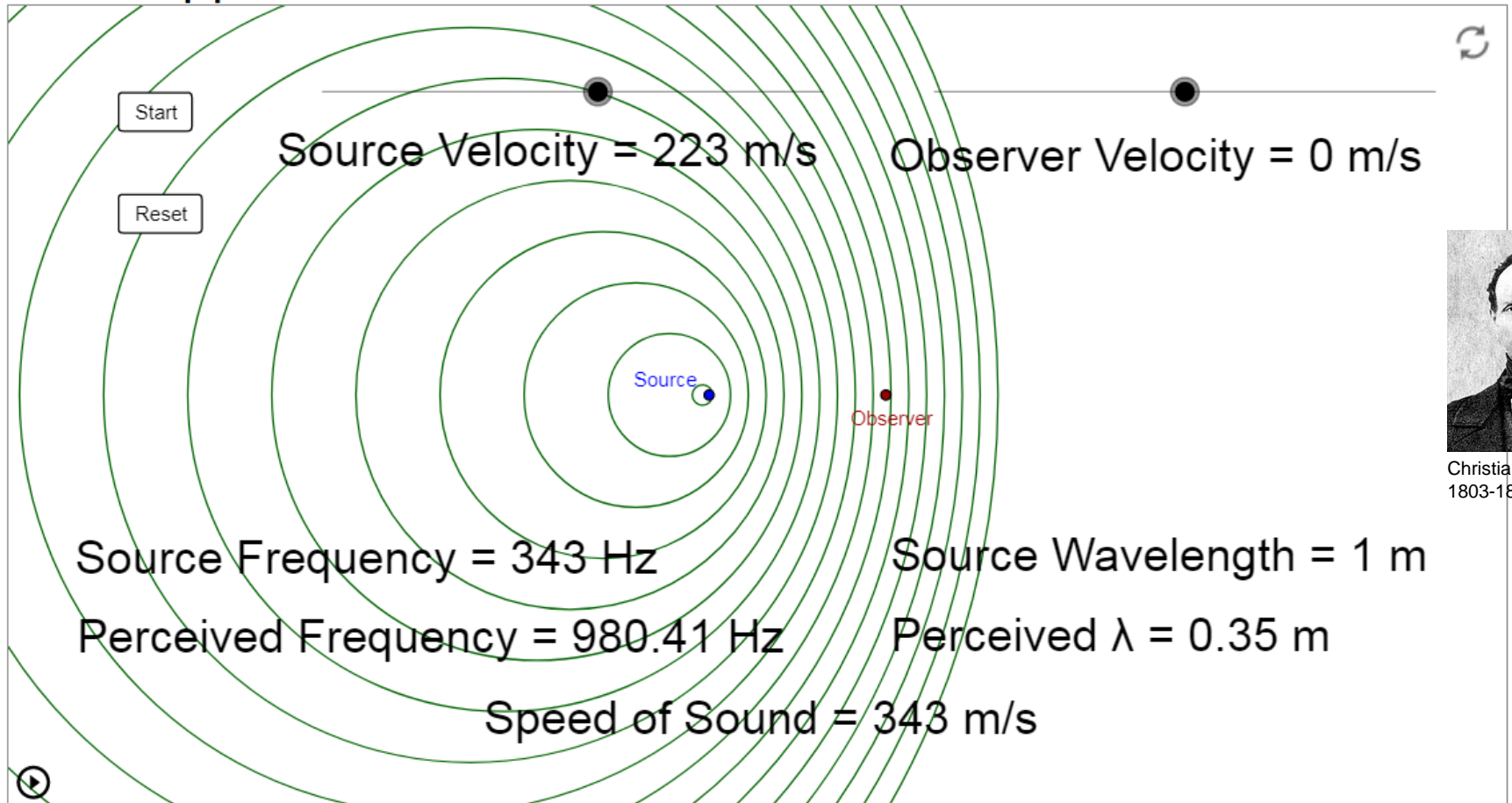
$Z_2$



# *oPhysics: Interactive Physics Simulations*

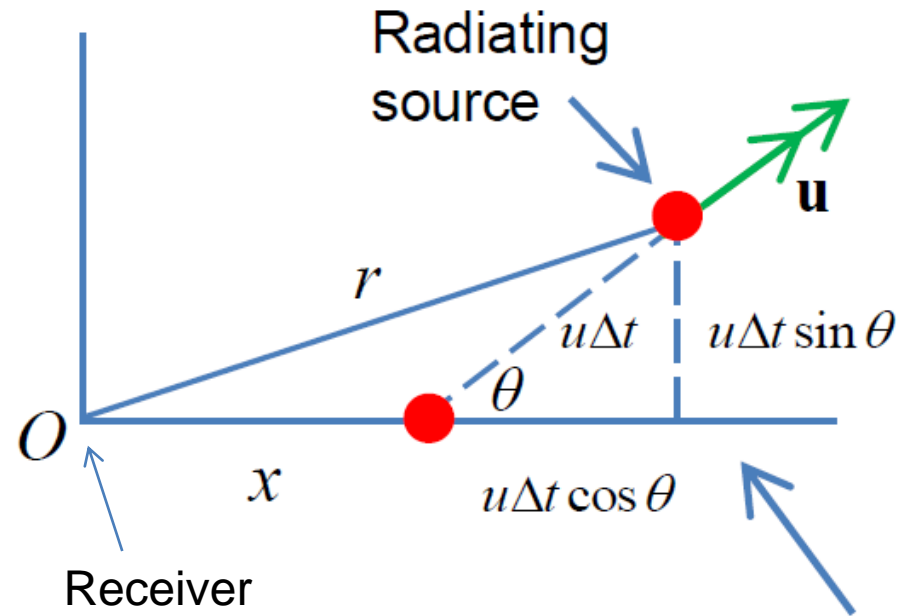
Home Kinematics Forces Conservation **Waves** Light E & M Rotation Fluids Modern Drawing Tools Fun Stuff

## The Doppler Effect & Sonic Boom



Consider a **receding** wave source of frequency  $f$ . It crosses the  $x$  axis of a Cartesian reference frame at angle  $\theta$  with speed  $u$ . The receiver of the waves is stationary at the origin of the Cartesian frame.

The speed of waves, relative to the observer, is  $w$ .



Cartesian reference frame with observer at the origin

Depending on the velocity of the wave source relative to the observer, the observer will experience a *frequency shift* from  $f$ . If the source *recedes*, the frequency *diminishes* and the *wavelength increases* ('**redshift**'). If the source is *approaching*, the observed *frequency will increase* and the *wavelength will decrease* ('**blueshift**').

# Excel model of Doppler Shift

- Plot the circular wavefront(s) when the source has got to position 8.
- i.e. you see a wavefront radius which *increases* as the *initial position* decreases, and also a *shift of wavefront centre*. The combination of these two effects causes the bunching up of wavefronts

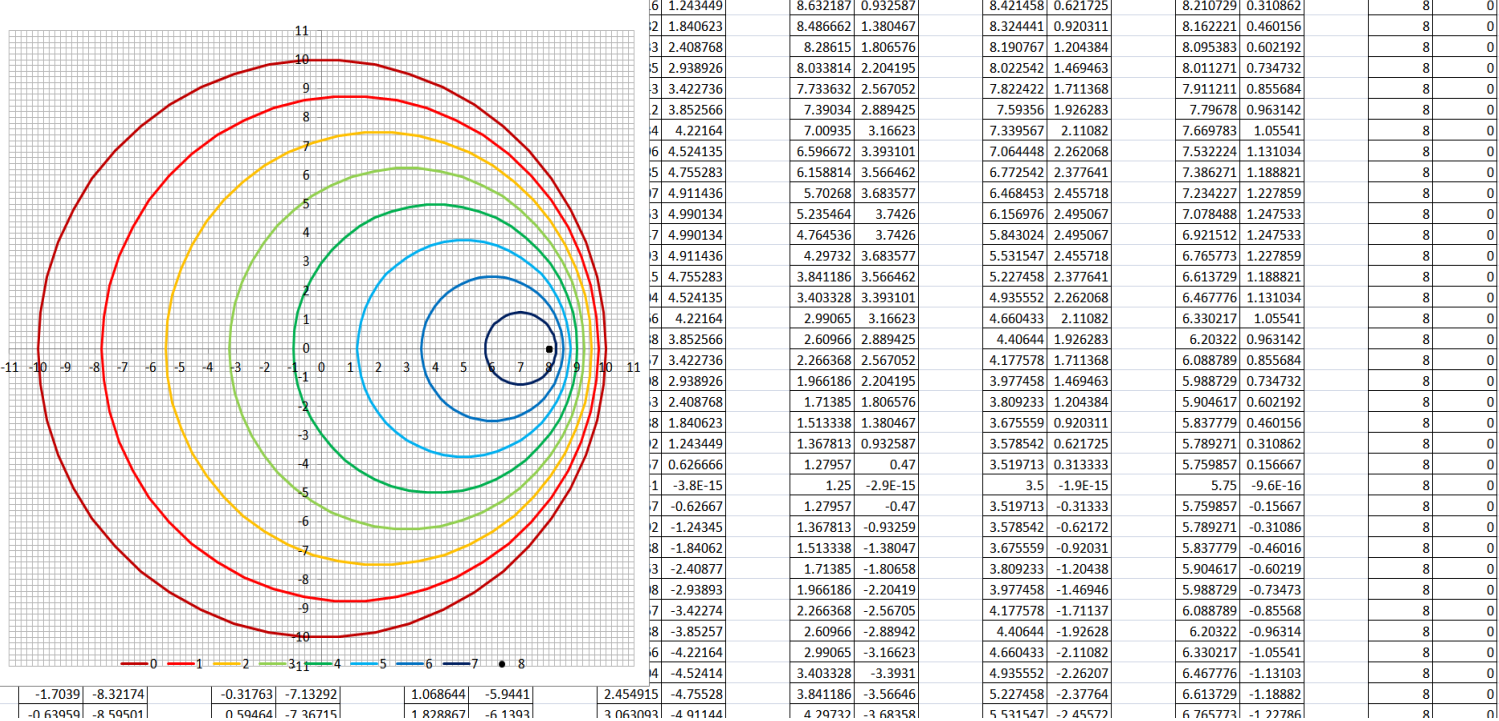
SIMULATION OF DOPPLER EFFECT OR MACH CONE	
A. French Dec 2021	
Wave source position:	0
Wavefront radius:	10.00

Mach number = speed of wave source / speed of waves

<b>M</b>	<b>0.8</b>
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0	1	2	3	4	5	6	7	8
10.00	8.75	7.50	6.25	5.00	3.75	2.50	1.25	0.00

	Wavefront circle		Wavefront circle		Wavefront circle		Wavefront circle		Wavefront circle		Wavefront circle		Wavefront circle		Wavefront circle		Wavefront circle					
Angle /rad	x	y	x	y	x	y	x	y	x	y	x	y	x	y	x	y	x	y				
0	10	0	9.75	0	9.5	0	9.25	0	9	0	8.75	0	8.5	0	8.25	0	8	0	8	0		
0.02	9.921147	1.253332	9.685832	2.486899	9.297765	3.681246	8.763067	4.817537	8.09017	5.877853	7.289686	6.845471	6.37424	7.705132	5.358268	8.443279	4.257793	9.048271	3.09017	9.510565	1.873813	9.822873
0.04	9.921147	1.253332	9.685832	2.486899	9.297765	3.681246	8.763067	4.817537	8.09017	5.877853	7.289686	6.845471	6.37424	7.705132	5.358268	8.443279	4.257793	9.048271	3.09017	9.510565	1.873813	9.822873

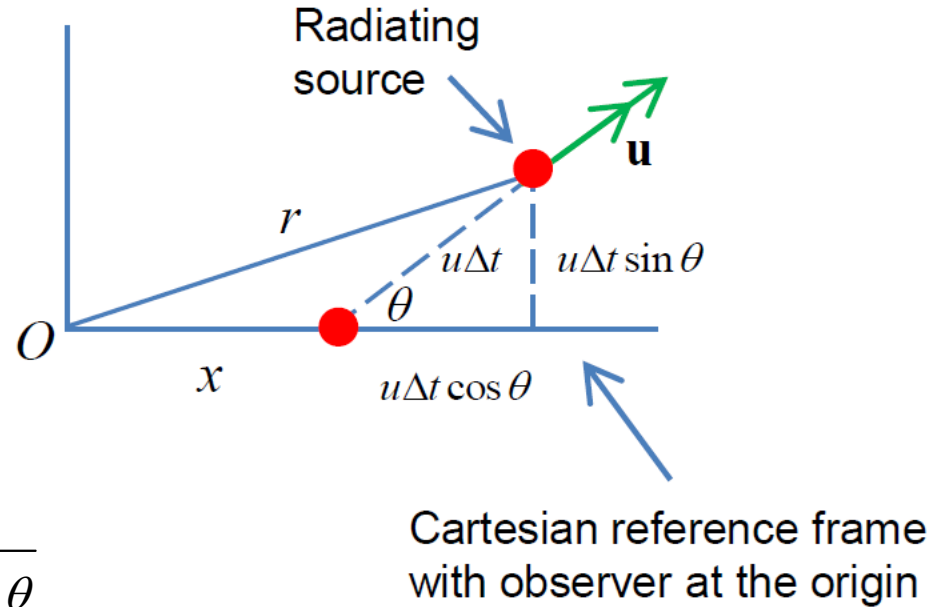




The period  $T$  of waves received by an observer (in the  $x$  direction) at the frame origin  $O$  is:

$$T = \Delta t + \frac{r - x}{w}$$

$\Delta t$ : time between wave crests at source  
 $w$ : wave speed  
 $r - x$ : extra distance travelled by source between wave crests



From geometry:

$$r = \sqrt{(x + u\Delta t \cos \theta)^2 + u^2 \Delta t^2 \sin^2 \theta}$$

$$r = \sqrt{x^2 + u^2 \Delta t^2 \cos^2 \theta + 2ux\Delta t \cos \theta + u^2 \Delta t^2 \sin^2 \theta}$$

$$r = \sqrt{x^2 + u^2 \Delta t^2 + 2ux\Delta t \cos \theta}$$

$$r = x \sqrt{1 + 2 \cos \theta \frac{u\Delta t}{x} + \left(\frac{u\Delta t}{x}\right)^2}$$

If  $u\Delta t \ll x$

$$r \approx x \sqrt{1 + 2 \cos \theta \frac{u\Delta t}{x}} \approx x \left(1 + \cos \theta \frac{u\Delta t}{x}\right) = x + u\Delta t \cos \theta$$

$$\therefore r - x \approx u\Delta t \cos \theta$$

Hence frequency of radiation received at O is  $F = 1/T$  where:

$$\frac{1}{F} = \Delta t + \frac{u\Delta t \cos \theta}{w} = \Delta t \left( 1 + \frac{u \cos \theta}{w} \right)$$

In a *Classical* scenario, where  $u, w$  are much less than the speed of light:

$$f = 1/\Delta t$$

Received frequency

$$\therefore F = \frac{1}{1 + \frac{u}{w} \cos \theta} f$$

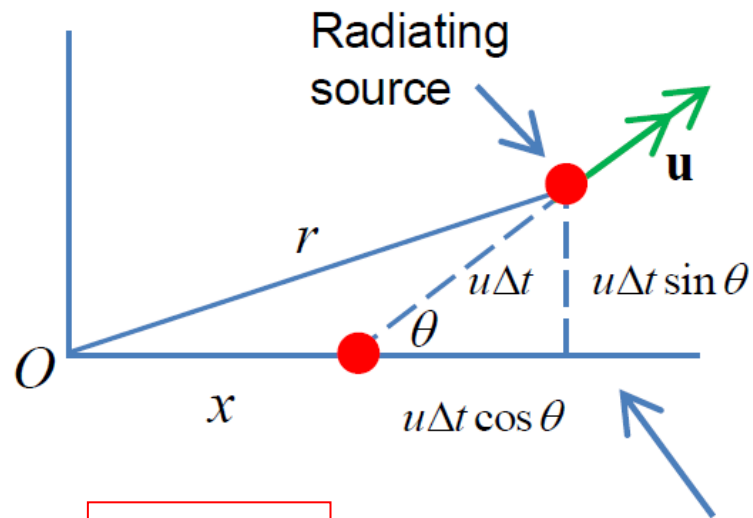
Emitted frequency

If  $w \gg u, \theta \approx 0$

$$F \approx f - \frac{u}{w} f$$

i.e. a **Doppler Shift** of:

$$\Delta f = \frac{u}{w} f$$



Cartesian reference frame with observer at the origin

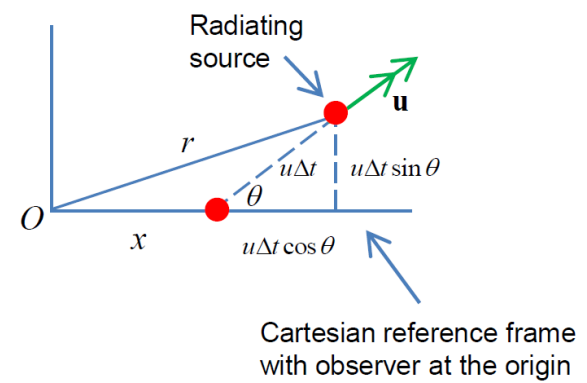
Recession velocity from observer

Observed wavelength

Emitted wavelength

$$\frac{\lambda_o}{\lambda_e} = 1 + \frac{u}{w} \cos \theta$$

wave speed

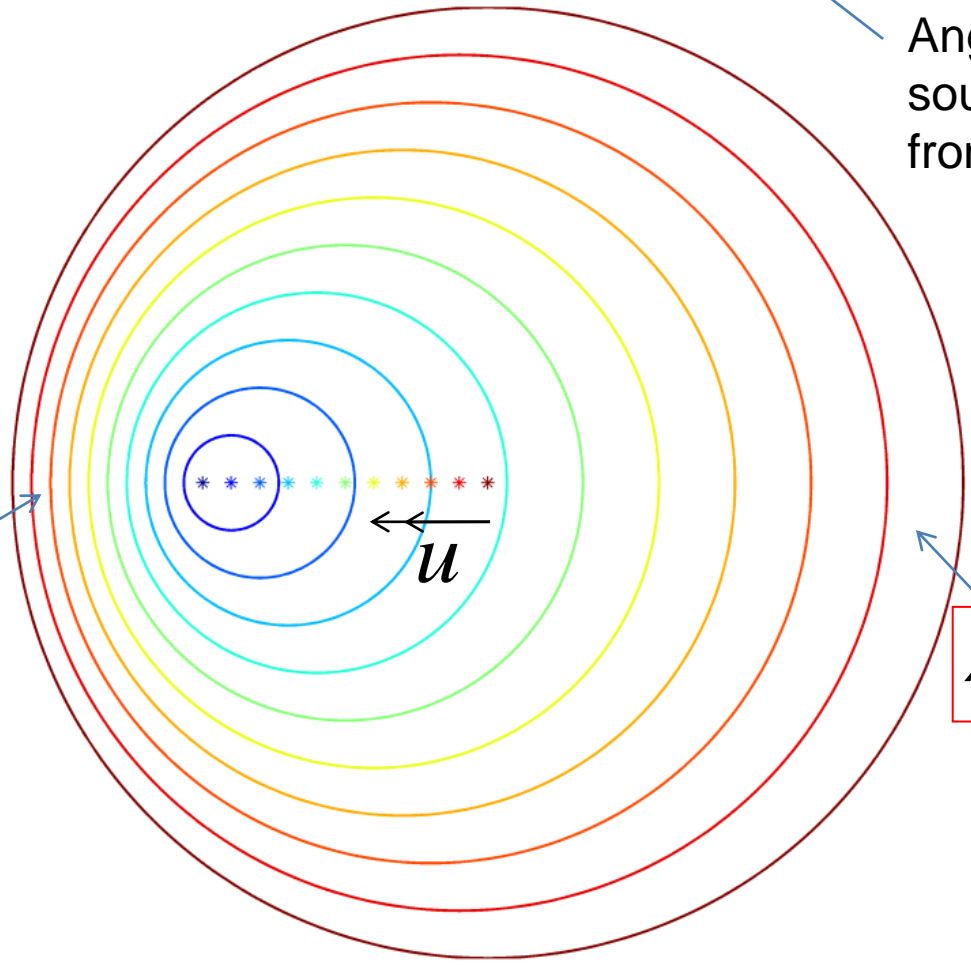


Angle of source velocity from horizontal

In this example:

$$\theta = 0$$

$$u = 0.6w$$



$$\lambda_o = \lambda_e (1 - 0.6)$$

$$\lambda_o = \lambda_e (1 + 0.6)$$

## Mach's construction

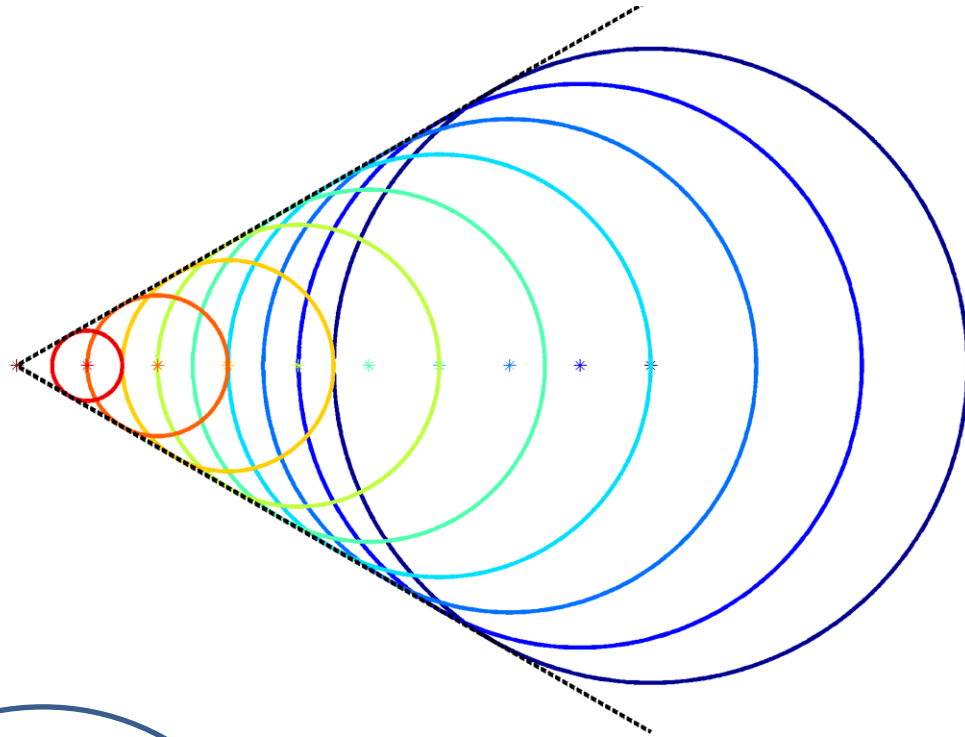
$c$  is the wave speed

$u$  is the velocity of the object making the waves

$$v/c = 2. \quad \sin^{-1}(c/v) = 30^\circ$$

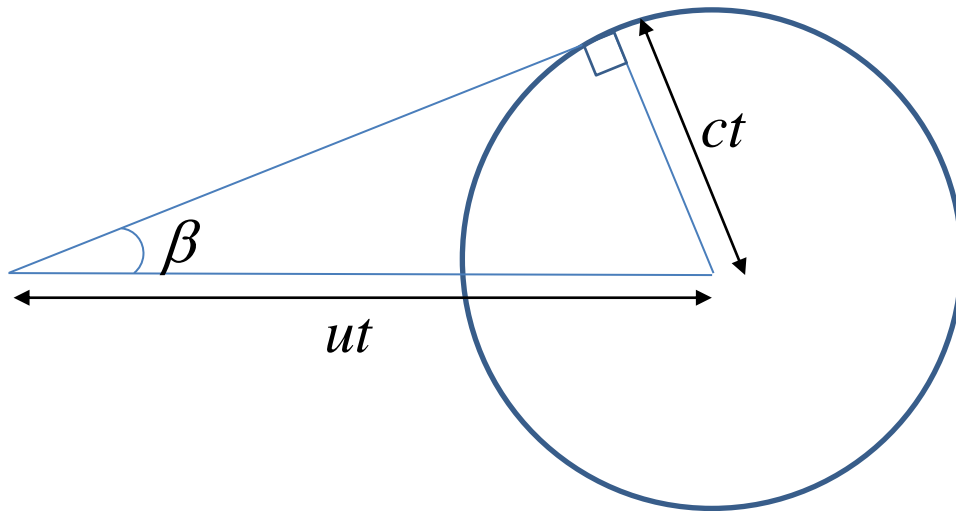


Ernst Mach 1838-1916



## Mach number

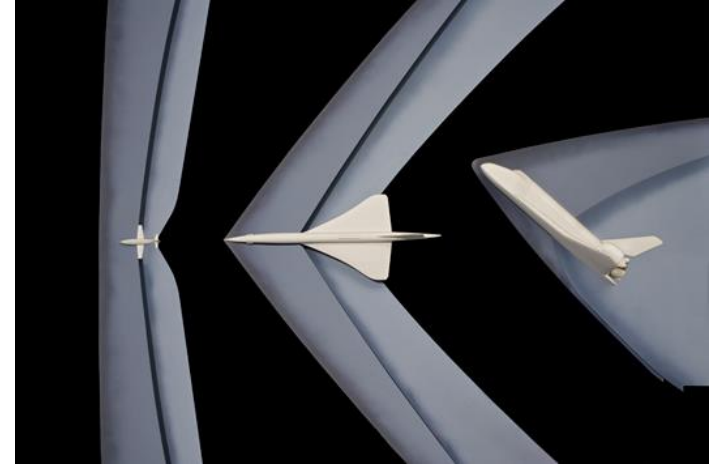
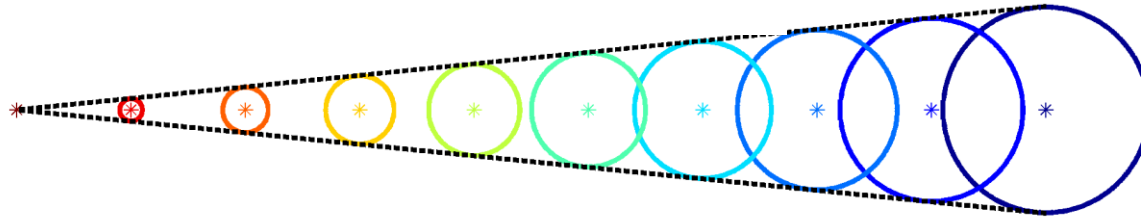
$$M = \frac{u}{c}$$



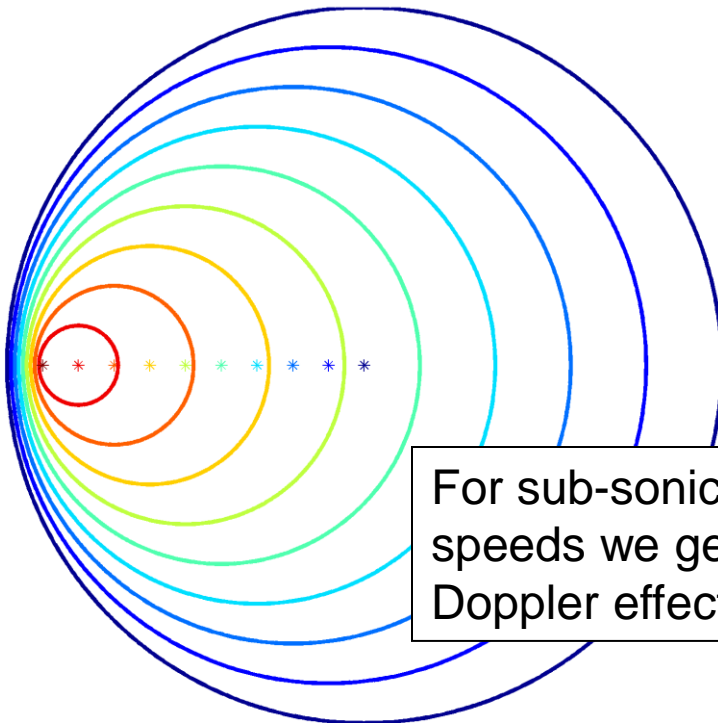
$$ut \sin \beta = ct$$

$$\therefore \beta = \sin^{-1} \left( \frac{c}{u} \right) = \sin^{-1} \frac{1}{M}$$

$$v/c = 10. \quad \sin^{-1}(c/v) = 5.7392^\circ$$

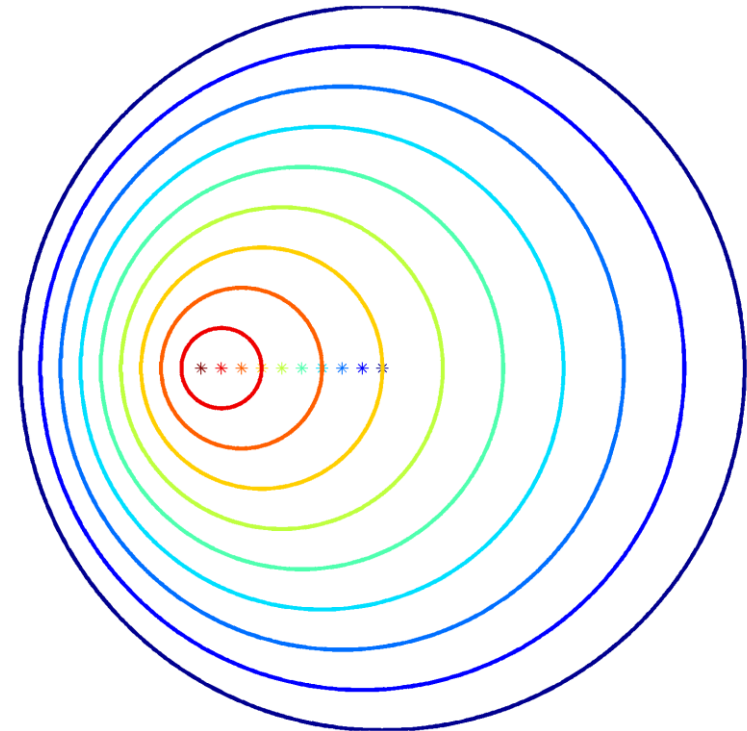


$$v/c = 0.9. \quad \sin^{-1}(c/v) = \text{NaN}^\circ$$



For sub-sonic speeds we get the Doppler effect!

$$v/c = 0.5. \quad \sin^{-1}(c/v) = \text{NaN}^\circ$$

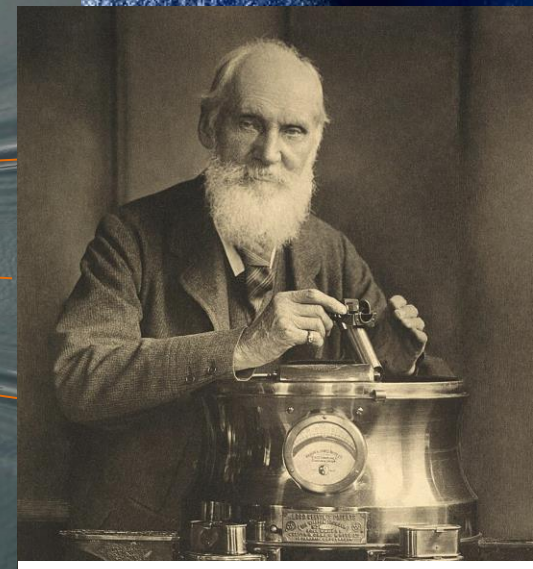
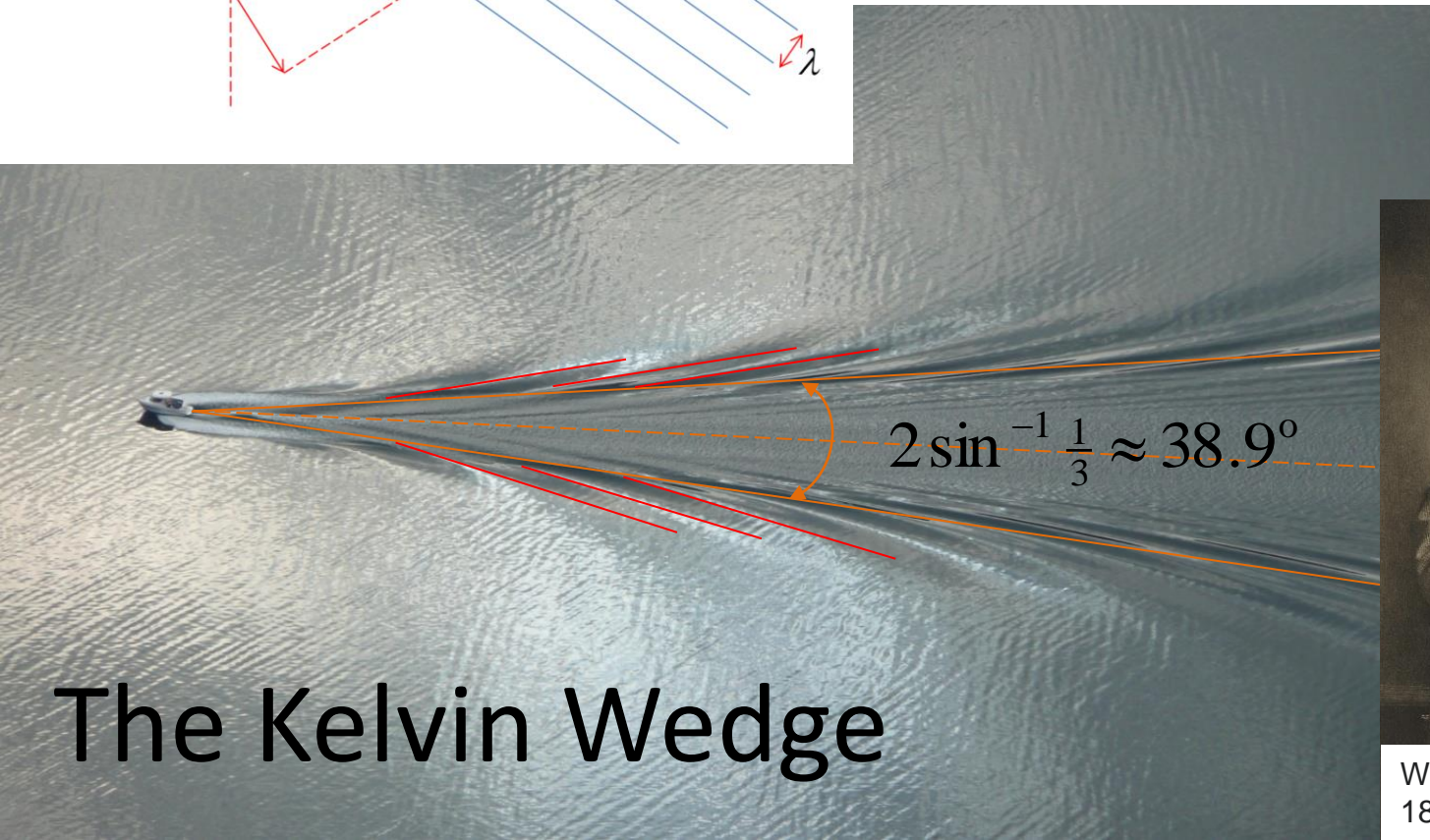
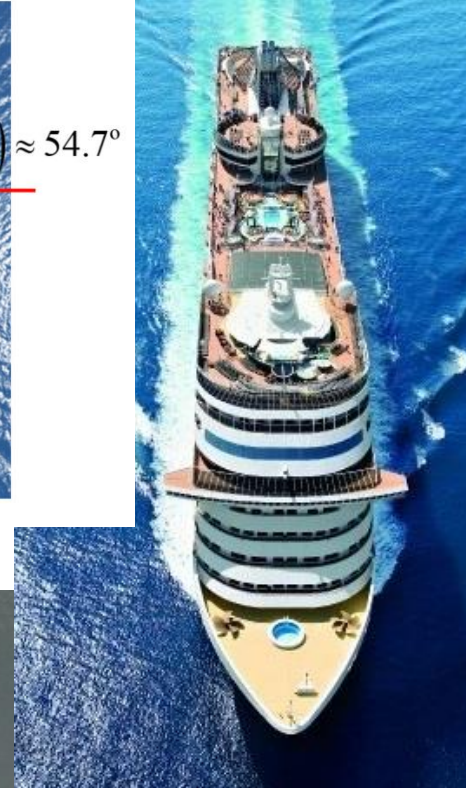
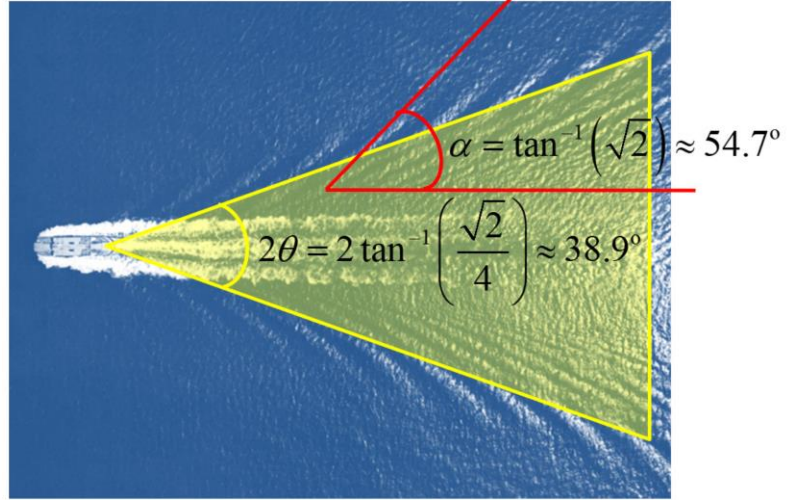
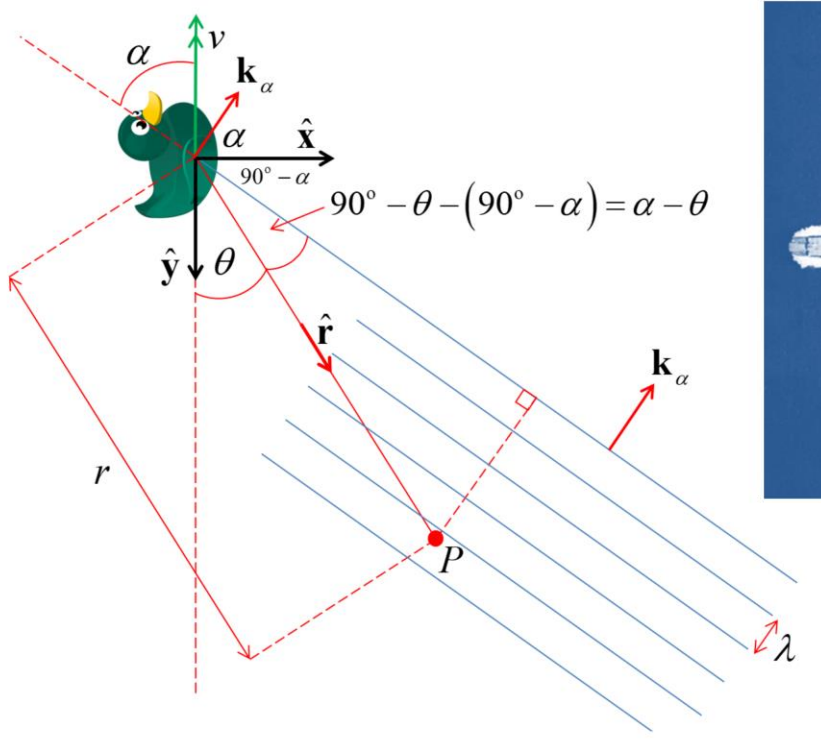




A wake is an *interference pattern* of waves formed by the motion of a body through a fluid. Intriguingly, the angular width of the wake produced by ships (and ducks!) in deep water is the same (about  $38.9^\circ$ ). A mathematical explanation for this phenomenon was first proposed by [Lord Kelvin](#) (1824-1907). The triangular envelope of the wake pattern has since been known as

<http://en.wikipedia.org/wiki/Wake>

# the Kelvin Wedge



William Thomson, 1<sup>st</sup> Baron Kelvin  
1824-1907

# The Kelvin Wedge