

$$7/ \quad (i) \quad p(E) = A e^{-E/k_B T} \quad E = \{0, \varepsilon\}$$

$$\therefore \text{Since } \sum_i p(E_i) = 1 \quad \text{Since a probability distribution}$$

$$\Rightarrow p(0) + p(\varepsilon) = 1$$

$$\therefore A (e^{-0} + e^{-\varepsilon/k_B T}) = 1$$

$$\therefore \boxed{A = \frac{1}{1 + e^{-\varepsilon/k_B T}}}$$

$$(ii) \quad \bar{E} = \sum_i E_i p(E_i) = \underset{\substack{\uparrow \\ E=0}}{0} + \boxed{\frac{\varepsilon e^{-\varepsilon/k_B T}}{1 + e^{-\varepsilon/k_B T}}} \quad \uparrow E=\varepsilon$$

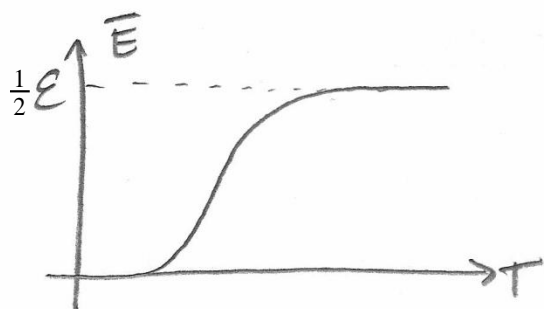
$$\text{Since } p(E_i) = \frac{e^{-E_i/k_B T}}{1 + e^{-\varepsilon/k_B T}}$$

$$\text{So: } \boxed{\bar{E} = \frac{\varepsilon}{e^{\varepsilon/k_B T} + 1}}$$

$$(iii) \quad \bar{E} = \frac{\varepsilon}{e^{\varepsilon/k_B T} + 1}$$

$$\text{when } T \rightarrow 0 \quad e^{\varepsilon/k_B T} \rightarrow \infty \quad \therefore \bar{E} \rightarrow 0$$

$$\text{when } T \rightarrow \infty \quad e^{\varepsilon/k_B T} \rightarrow 0 \quad \therefore \bar{E} \rightarrow \frac{1}{2} \varepsilon$$



(iv) Einstein's model of average energy of thermal motion of a solid

a)

$$\bar{E} = \frac{hf}{e^{hf/k_B T} - 1}$$

when $T \rightarrow 0$, $e^{hf/k_B T} \rightarrow \infty \therefore \bar{E} \rightarrow 0$

when $T \rightarrow \infty$, $e^{hf/k_B T} \rightarrow e^0 \rightarrow 1 \therefore \bar{E} \rightarrow \infty$

$$\text{Now } \frac{d\bar{E}}{dT} = \frac{-hf}{(e^{hf/k_B T} - 1)^2} \times e^{hf/k_B T} \times \left(-\frac{hf}{k_B T^2}\right)$$

$$= \left(\frac{hf}{k_B T}\right)^2 k_B \frac{e^{hf/k_B T}}{(e^{hf/k_B T} - 1)^2}$$

$$= \frac{k_B x^2 e^x}{(e^x - 1)^2}$$

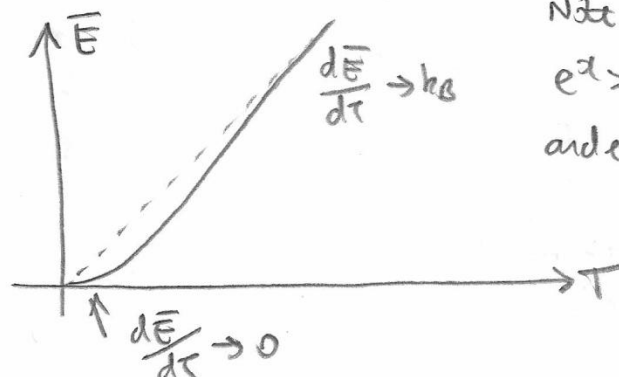
$$\text{where } x = \frac{hf}{k_B T}$$

when $T \rightarrow \infty$, $x \rightarrow 0$ so $e^x = 1 + x + \frac{x^2}{2} + \dots$
 $\rightarrow 1+x$. $\therefore e^x - 1 \rightarrow x$

$$\text{So } \frac{d\bar{E}}{dT} \rightarrow \frac{k_B x^2 (1+x)}{x^2} \rightarrow k_B (1+x)$$

and if $x \rightarrow 0$, $\frac{d\bar{E}}{dT} \rightarrow k_B$ constant gradient

So:



Note when $T \rightarrow 0$, $x \rightarrow \infty$
 $e^x \gg 1 \therefore \frac{d\bar{E}}{dT} \rightarrow \frac{k_B x^2}{e^x}$
 and $e^x \gg x^2$
 $\rightarrow 0$

b) $\therefore C = 3N_A \frac{d\bar{E}}{dT}$ Molar heat capacity

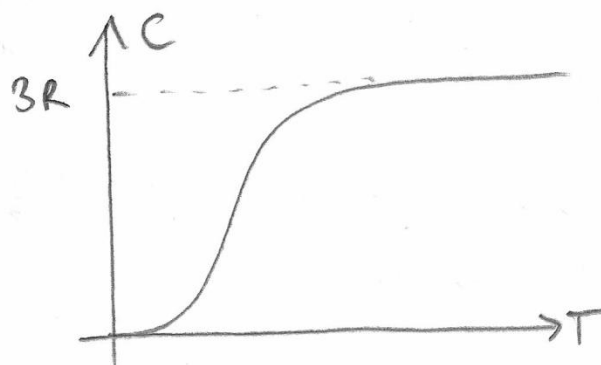
$$\therefore C = 3N_A k_B \frac{x^2 e^x}{(e^x - 1)^2} \quad x = \frac{hf}{k_B T}$$

Now $R = N_A k_B$

$$\therefore \boxed{C = \frac{3R x^2 e^x}{(e^x - 1)^2}}$$

As discussed above, as $T \rightarrow \infty$, $x \rightarrow 0$
and $e^x \rightarrow 1 + x \quad \therefore \boxed{C \rightarrow 3R}$

when $T \rightarrow 0$, $x \rightarrow \infty$ and $\boxed{C \rightarrow 0}$



$C = 3R$ is the 'Dulong & Petit law', which is true at $x \rightarrow 0$

i.e. $\frac{hf}{k_B T} \ll 1$

$$\therefore \boxed{T \gg \frac{hf}{k_B}}$$

For various elements the 'Einstein frequency' is:

Au	$0.286 \times 10^{13} \text{ Hz}$
Cu	0.577 ''
Fe	0.789 ''
C	3.745 ''
Si	1.083 ''

{ Note Einstein's model is only approximate.
Debye theory assumes a spectrum of possible vibration frequencies }