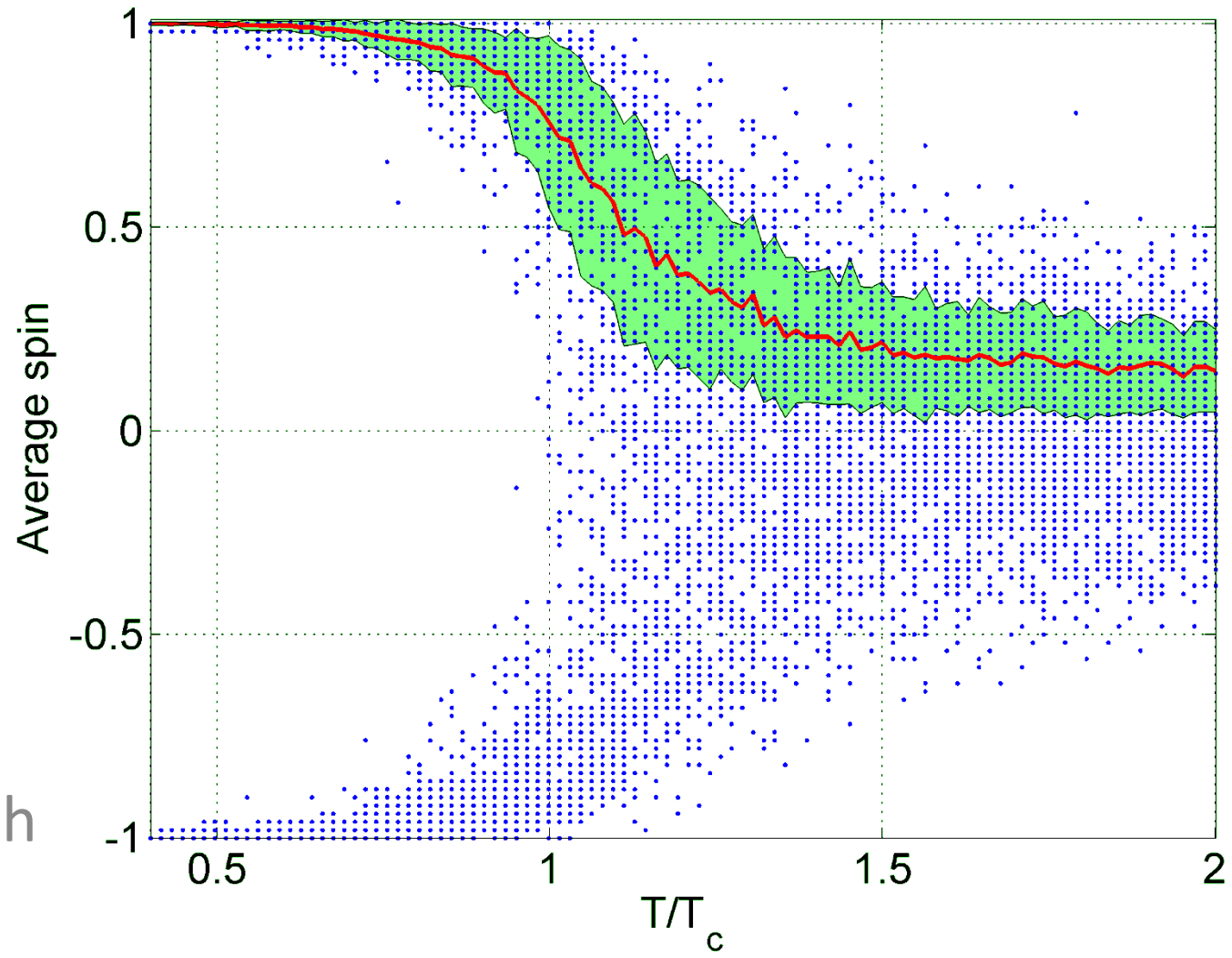


# The Ising model of Ferromagnetism

Average spin.  $J=0.04\text{eV}$ .  $T_c=1043\text{ K}$



# The Ising Model of Ferromagnetism

All atoms will respond in some fashion to **magnetic fields**. The angular momentum (and spin) properties of electrons imply a circulating charge, which means they will be subject to a Lorentz force in a magnetic field. **However the effects of *diamagnetism, paramagnetism and anti-ferromagnetism* are typically very small. Ferromagnetic materials** (iron, cobalt, nickel, some rare earth metal compounds) respond strongly to magnetic fields and can intensify them by orders of magnitude. i.e. the *relative permeability* can be tens or hundreds, or possibly thousands.

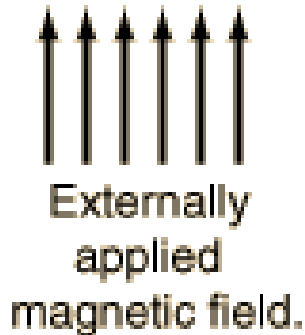
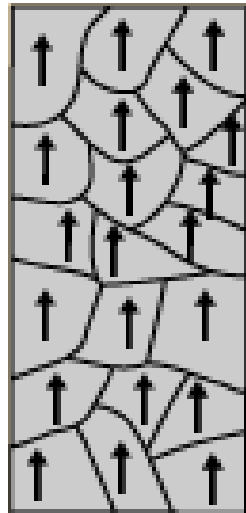
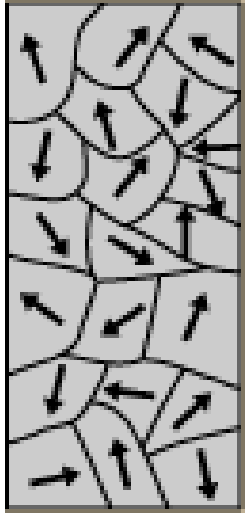
The Ising model is a simplified model of a **ferromagnet** which exhibits a **phase transition** above the **Curie temperature**. Below this, magnetic dipole alignment will tend to cluster into **domains**, and it is these micro-scale groupings which give rise to ferromagnetic behaviour.



Ernst Ising (1900-1998)

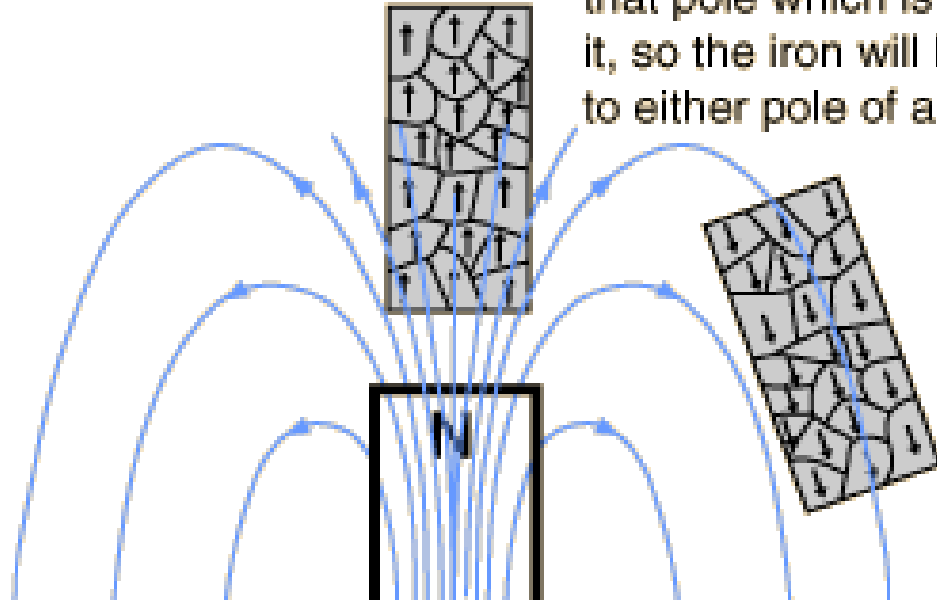


# “Soft” magnetism - Ferromagnets



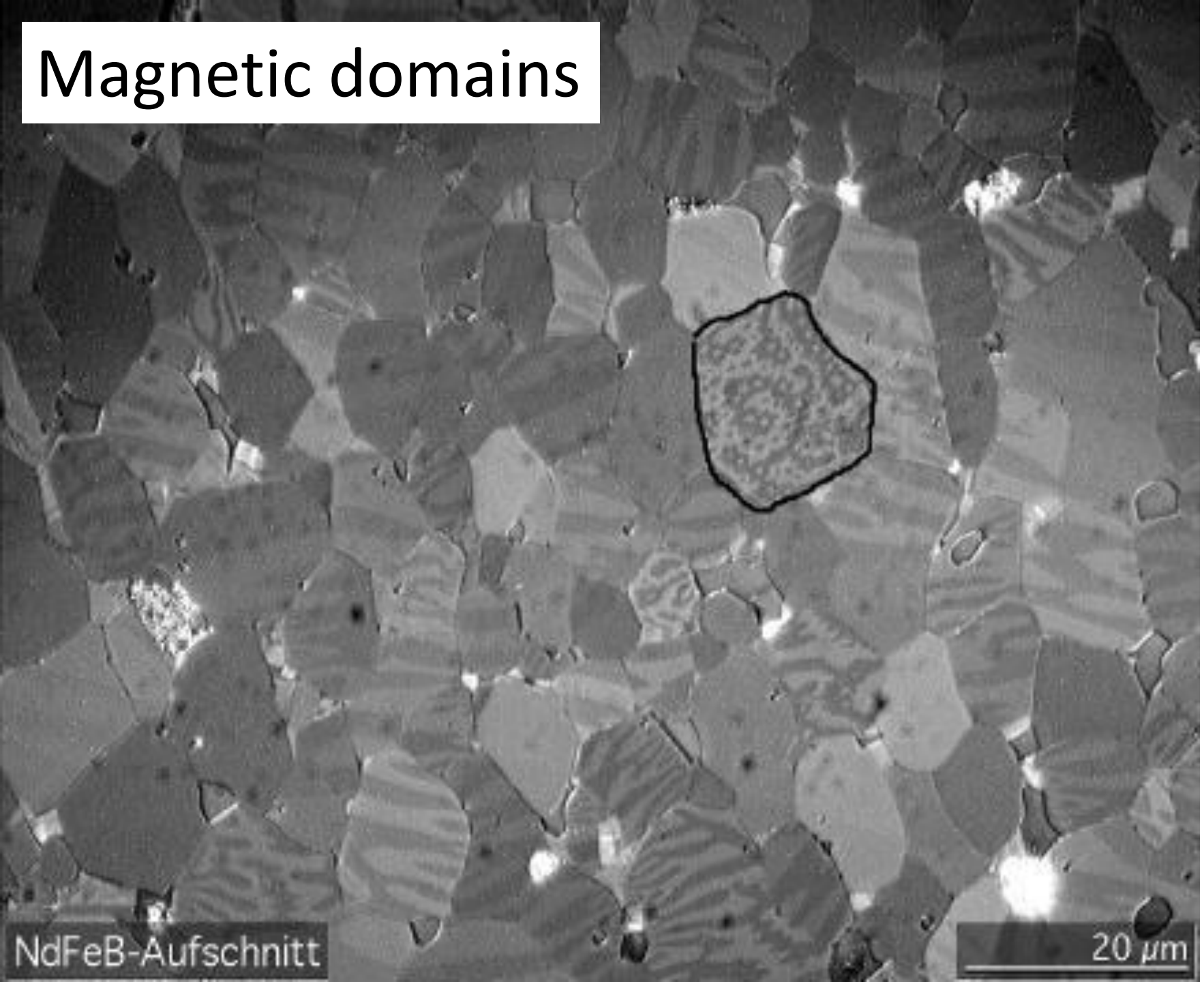
Externally applied magnetic field.

Iron will become magnetized in the direction of any applied magnetic field. This magnetization will produce a magnetic pole in the iron opposite to that pole which is nearest to it, so the iron will be attracted to either pole of a magnet.



Unlike permanent “hard” magnets, once the applied field is removed, the domain alignment will randomize again, effectively zeroing the net magnetism.

# Magnetic domains

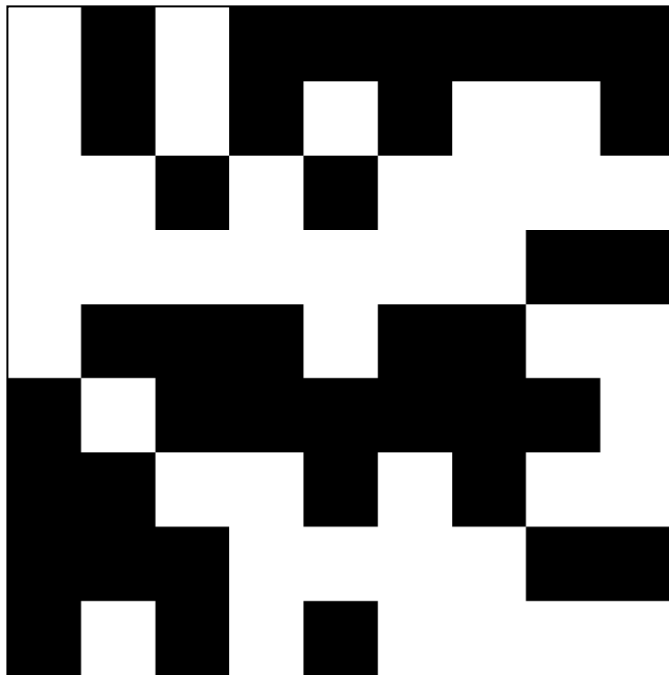


NdFeB-Aufschnitt

20 μm

The **Ising model** can be used to demonstrate spontaneous mass alignment of magnetic dipoles, and possibly a mechanism for domain formation.

Perhaps the simplest model which yields characteristic behaviour is an  $N \times N$  square grid, where each square is initially randomly assigned a +1 or -1 value, with equal probability. The +/-1 values correspond to a single direction of magnetic dipole moment in a rectangular lattice of ferromagnetic atoms, or in the case of individual electrons, *spin*.



10 x 10 grid

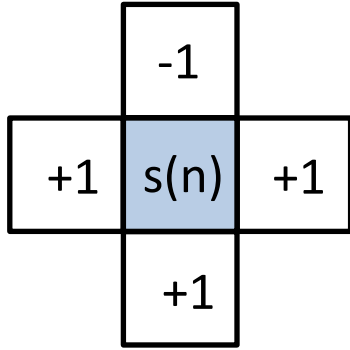
White squares  
represent +1  
**Black** squares  
represent -1



100 x 100 grid

## Original Metropolis algorithm

1. Choose one square at random from the  $N \times N$  grid. Let its spin be  $s(n) = +1$  or  $-1$ .
2. Find the spins of the nearest neighbours. Use *circular boundary conditions* e.g. if  $s(n)$  is at the edge of the grid, use the nearest neighbour to be that of the other end.



3. Compute a sum of **spin-coupling energies** for  $s(n)$  and its neighbours, and work out the energy change if  $s(n)$  were to **change sign**

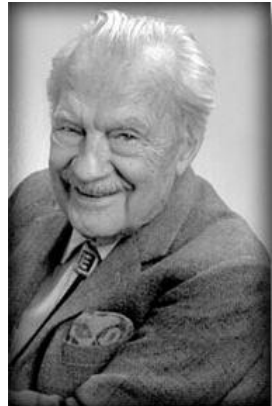
$$\Delta E = 2 \times \left( F + J \sum_{k=1}^4 s_n(k) \right) s(n)$$

$J$  is the spin coupling energy in eV and  $F$  is the energy in eV associated with the alignment of spin  $s(n)$  with an applied external magnetic field. Let us ignore any energy contributions from non-nearest neighbours.

$$r \sim U(0,1)$$

Now change the sign of spin  $s(n)$  according to the following rule:

$$s(n) \rightarrow -s(n) \quad \text{if} \quad e^{-\frac{\Delta E}{k_B T}} \geq r \quad \text{or} \quad \Delta E < 0$$



Nicholas  
Metropolis  
1915-1999

Apply the Metropolis method for  $L \times N \times N$  iterations, and then compute from the  $N \times N$  grid the following parameters

$$\langle s \rangle = \frac{1}{N^2} \sum_{n=1}^{N^2} s(n) \quad \text{Mean spin}$$

$$\langle E \rangle = -\frac{1}{2} \frac{1}{N^2} \sum_{n=1}^{N^2} \left( J \sum_{k=1}^4 s_n(k) + F \right) s(n) \quad \text{Mean energy per spin}$$

$$k_B T^2 \langle C \rangle = \frac{1}{4} \frac{1}{N^2} \sum_{n=1}^{N^2} \left( J s(n) \sum_{k=1}^4 s_n(k) + F s(n) \right)^2 - \langle E \rangle^2$$

Heat capacity in eV per K

This is a well known result in Statistical Thermodynamics

$$k_B T^2 \langle C \rangle = \text{Var}[E]$$

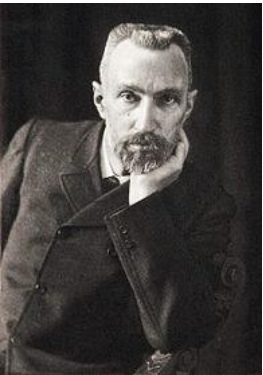
For a 2D Ising model, Lars Onsager determined in 1944 the relationship between the phase transition Curie temperature and coupling energy  $J$

$$\frac{2J}{\ln\left(1 + \sqrt{2}\right)} = k_B T_C$$

(Note this expression assumes coupling energy  $J$  is in joules)

Boltzmann's constant

$$k_B = 1.38 \times 10^{-23} \text{ JK}^{-1}$$



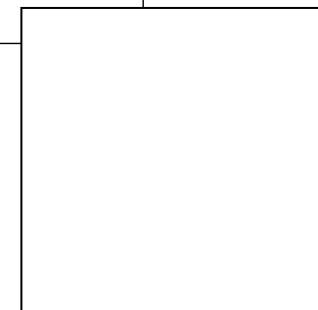
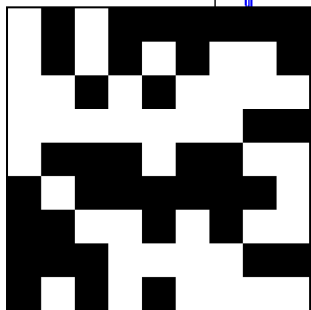
Peter Curie  
(1859-1906)



Lars Onsager  
(1903-1976)

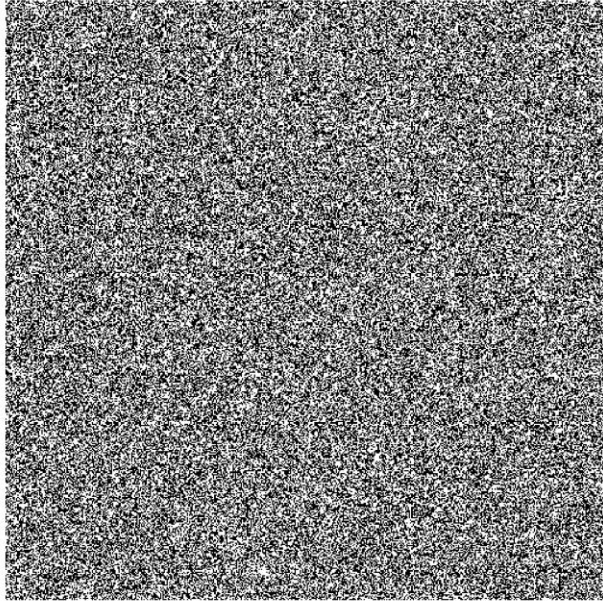


Mean spin vs iteration.  $N=10$ ,  $T/T_c=0.5$



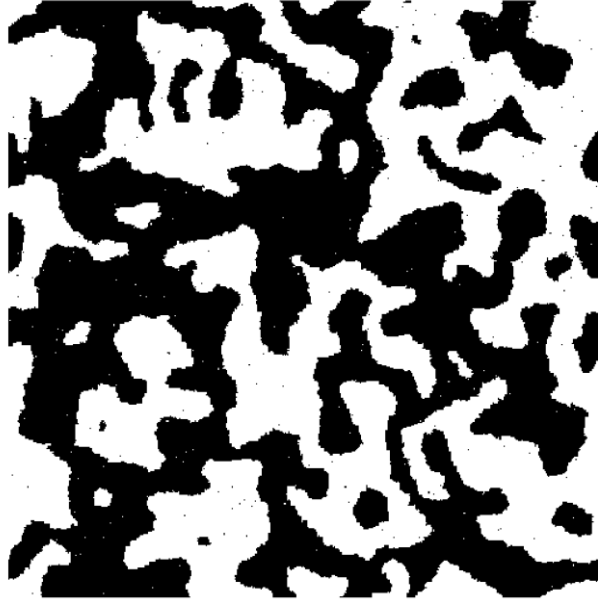
iteration = 1/2000

Mean spin=0.002568,  $T/T_c=0.5$



iteration = 2000/2000

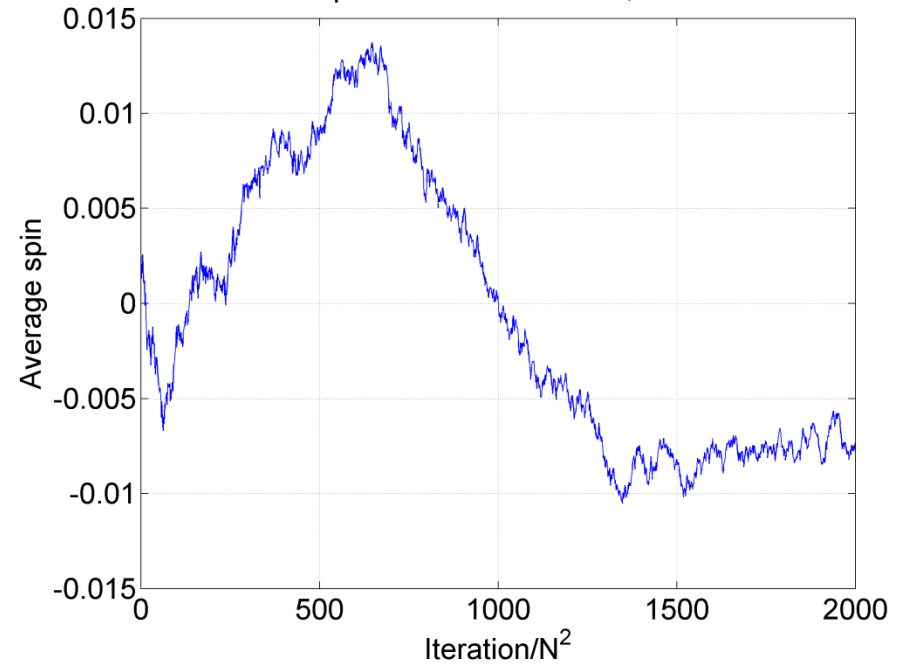
Mean spin=-0.007728,  $T/T_c=0.5$



For a 500 x 500 grid, a similar equilibrium is not yet reached, even after  $I = 2000 \times 500 \times 500$  iterations.

However, domain-like structures are clearly visible in this intermediate state.

Mean spin vs iteration.  $N=500$ ,  $T/T_c=0.5$



# Results of a MATLAB simulation:

10 x 10 grid

I = 2000 (x 10 x 10) iterations of Metropolis algorithm

R = 100 repeats for each temperature

100 different temperatures from  $T/T_c = 0.0 \dots 2.0$

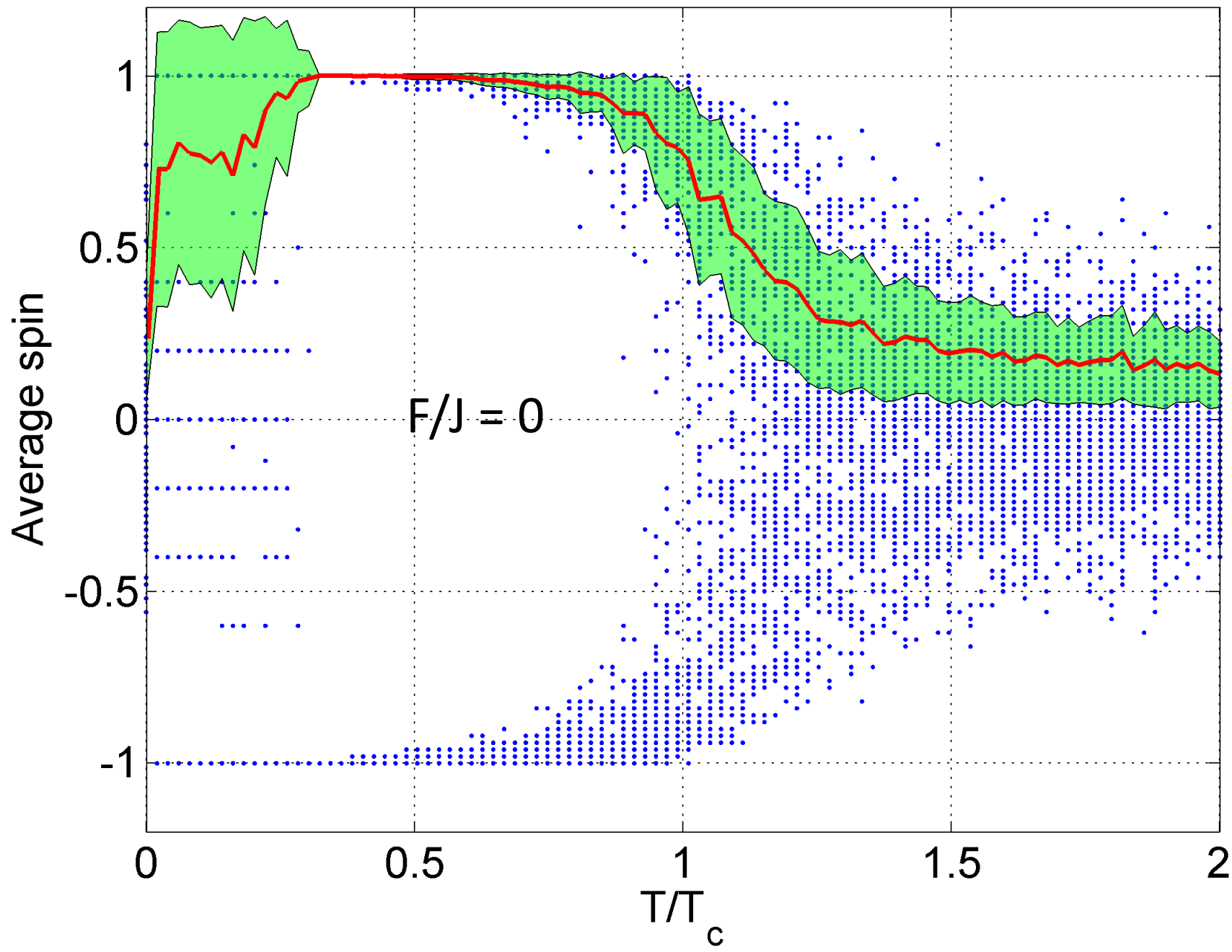
21 different F/J values from -2 to 2

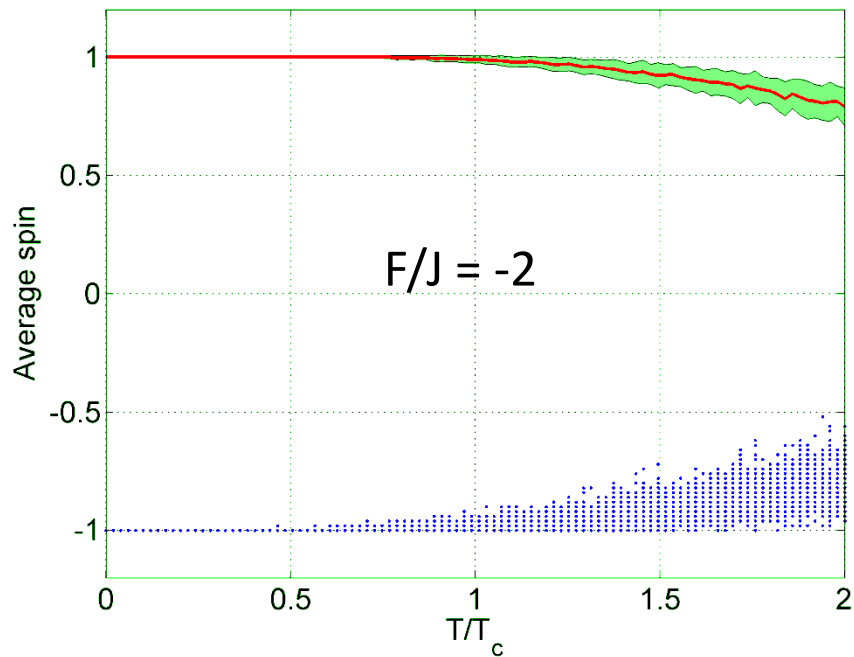
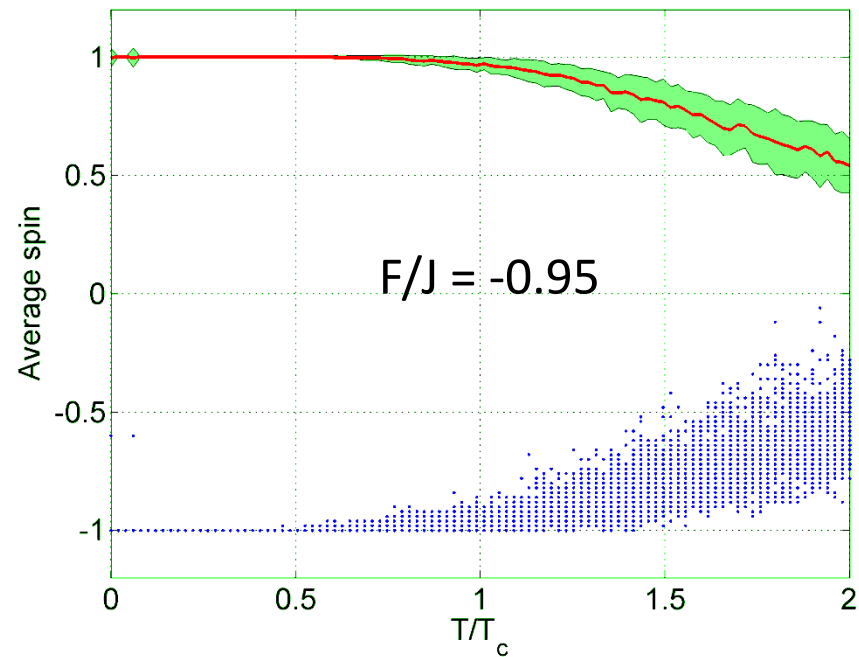
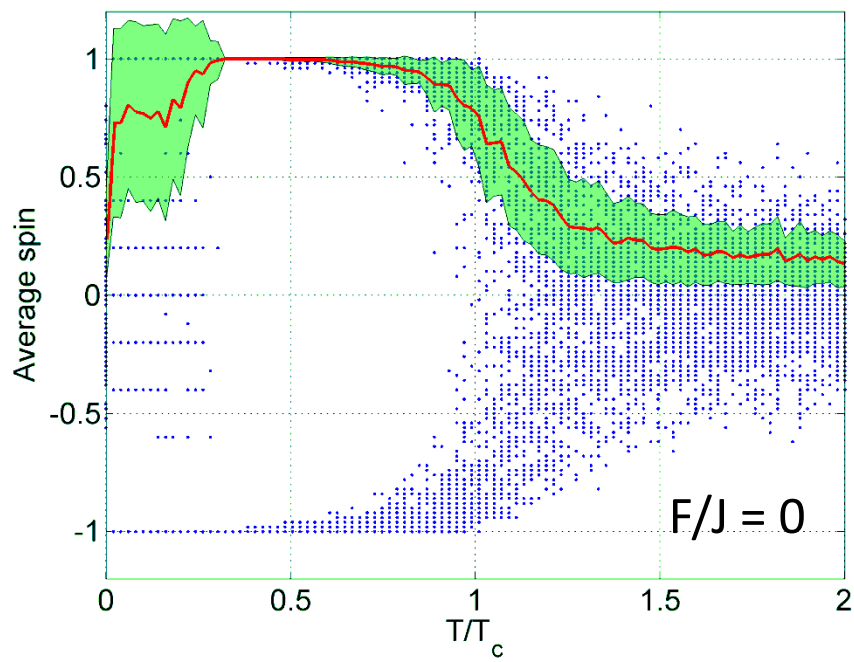
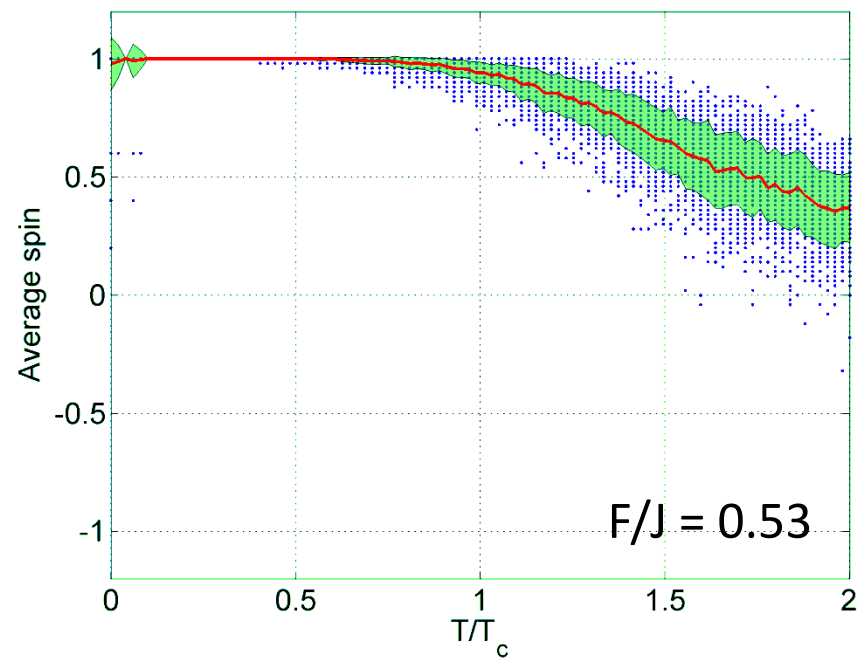
i.e.  $2000 \times 10 \times 10 \times 100 \times 100 \times 21 =$

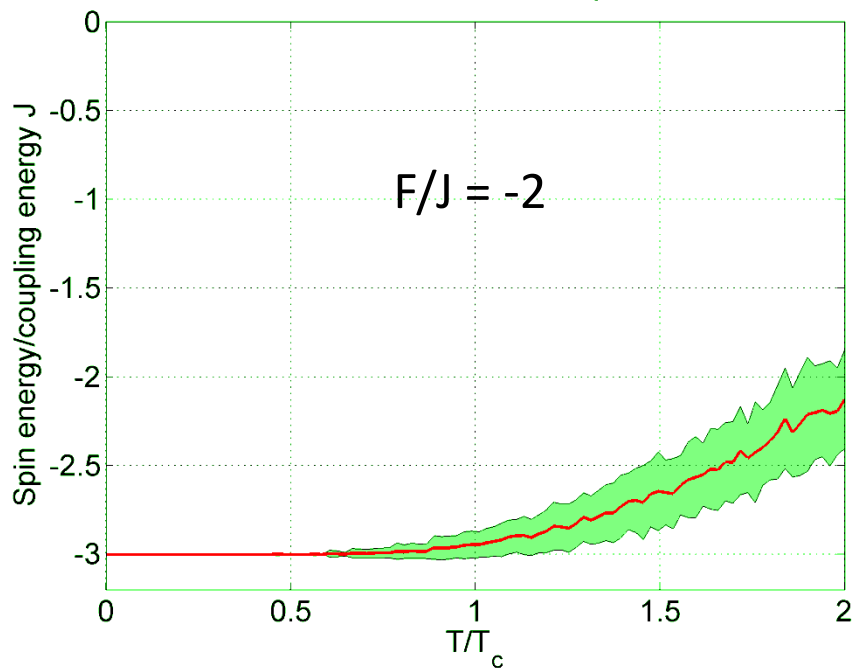
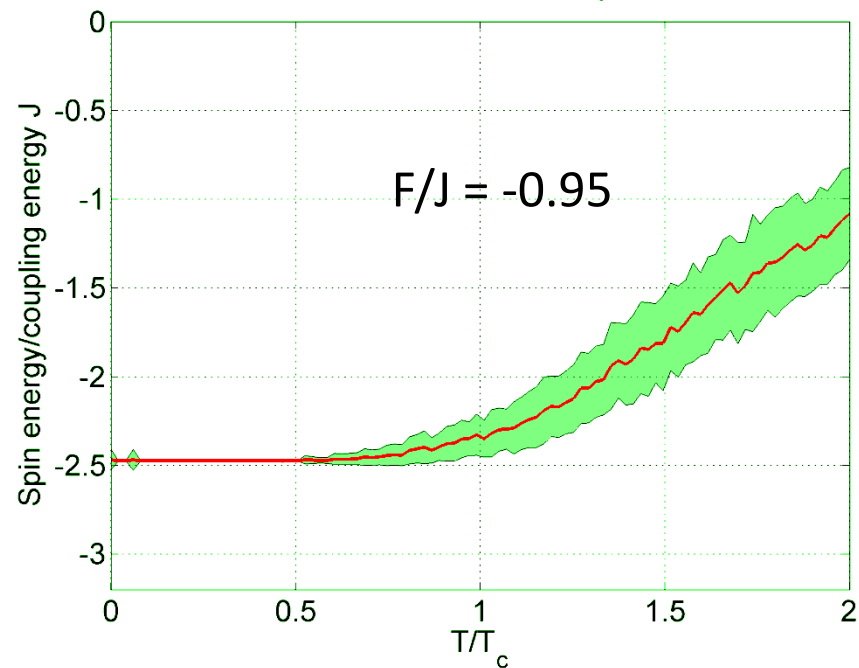
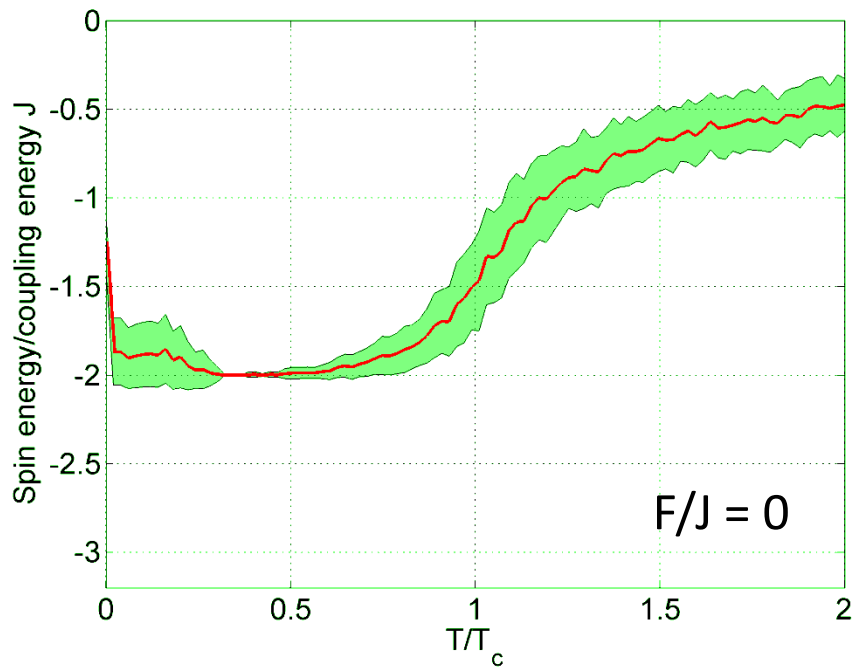
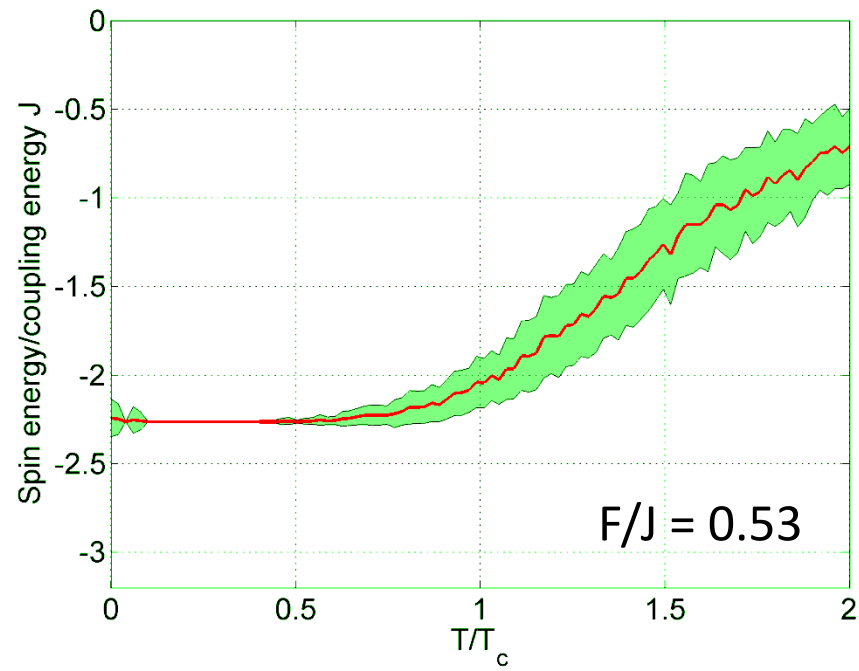
**42 billion iterations** of the Metropolis algorithm

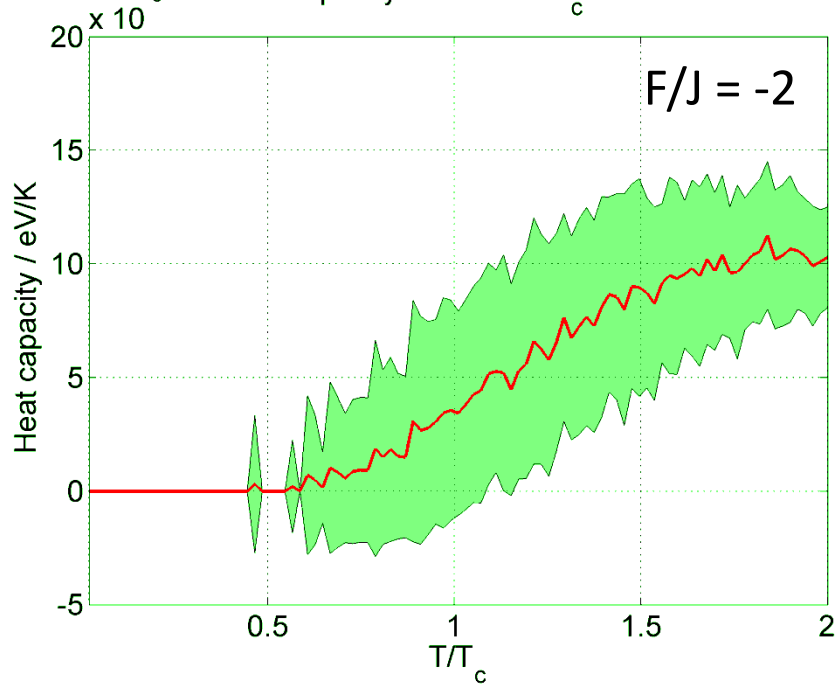
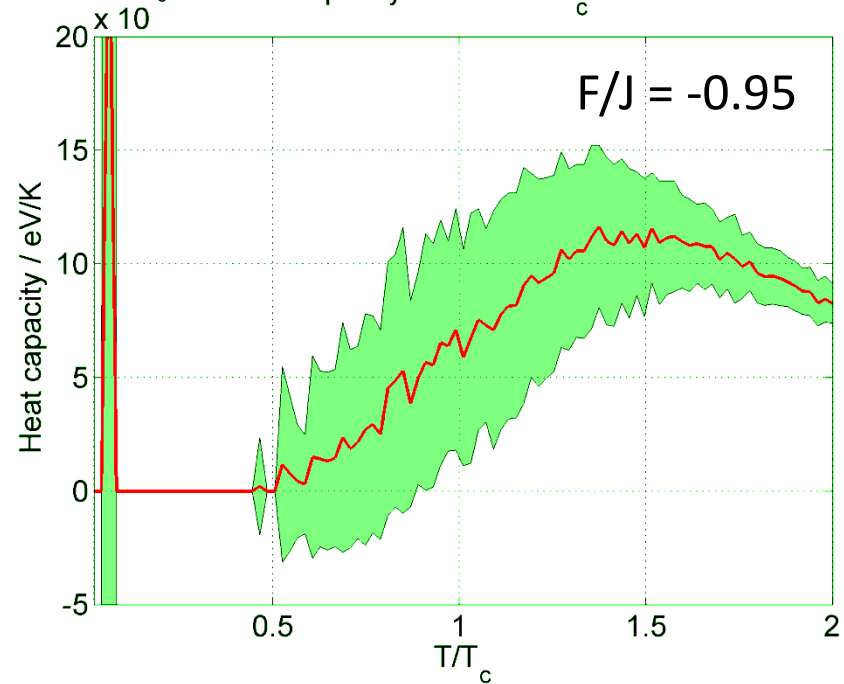
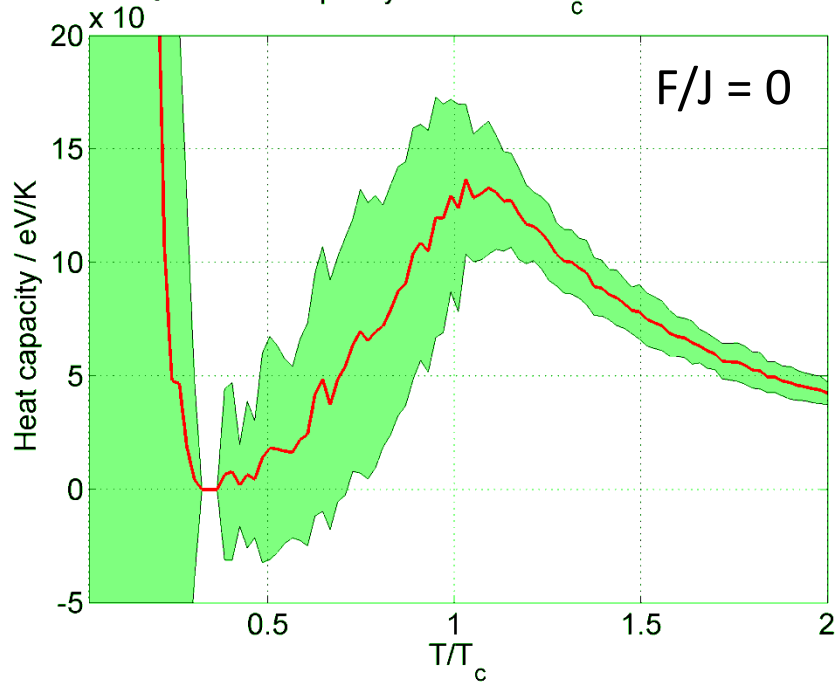
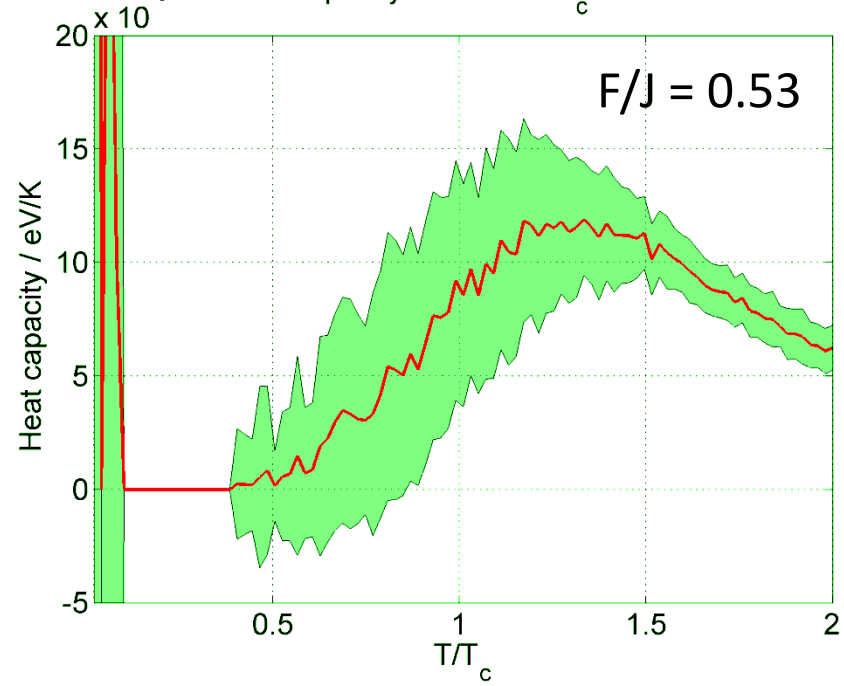
Running time on an i5 PC was about five days! Opportunity for *parallel processing*.

Average spin.  $J=0.04\text{eV}$ .  $T_c=1043\text{ K}$

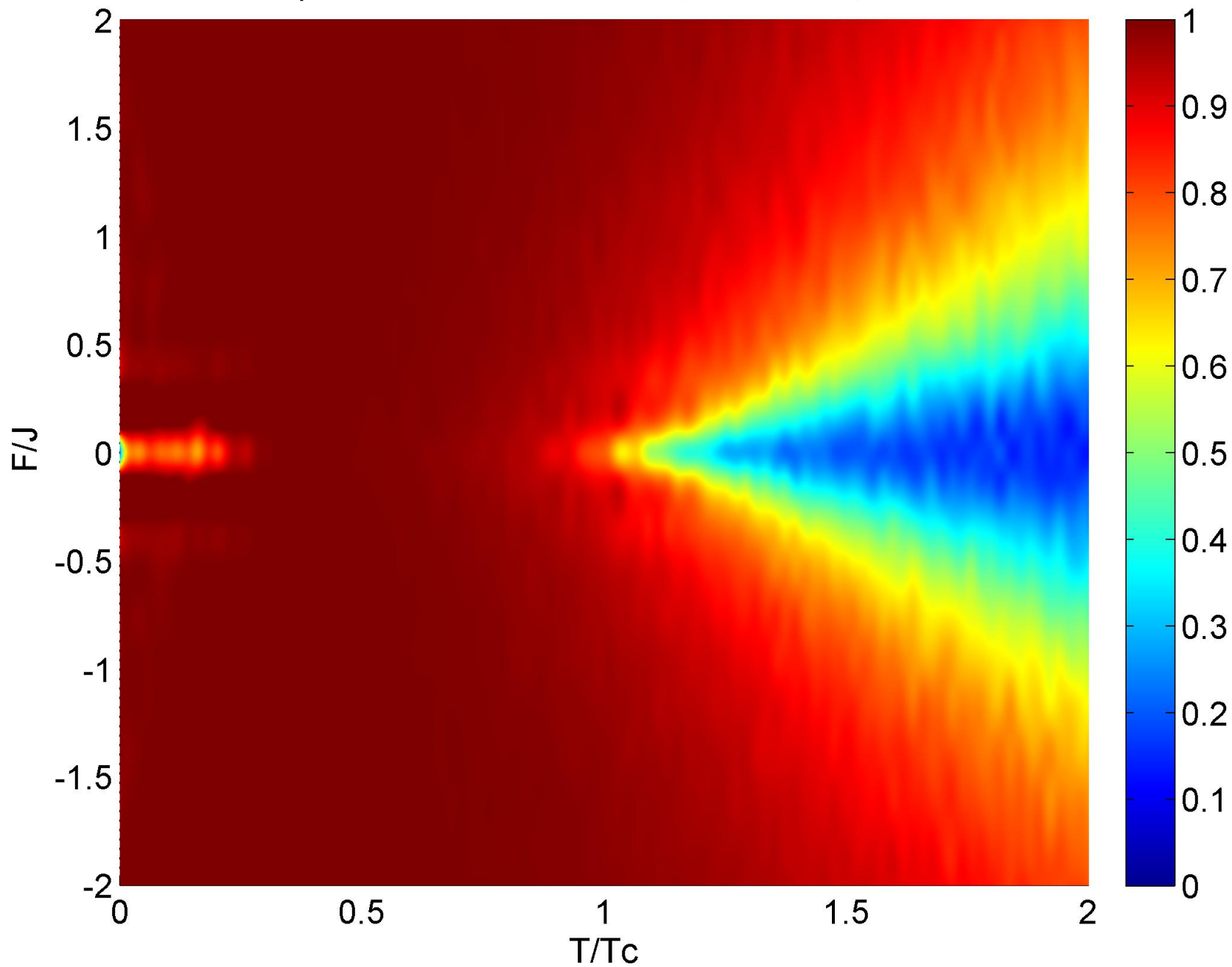


Average spin.  $J=0.04\text{eV}$ .  $T_c=1043\text{ K}$ Average spin.  $J=0.04\text{eV}$ .  $T_c=1043\text{ K}$ Average spin.  $J=0.04\text{eV}$ .  $T_c=1043\text{ K}$ Average spin.  $J=0.04\text{eV}$ .  $T_c=1043\text{ K}$ 

Spin energy/J.  $J=0.04\text{eV}$ .  $T_c=1043\text{ K}$ Spin energy/J.  $J=0.04\text{eV}$ .  $T_c=1043\text{ K}$ Spin energy/J.  $J=0.04\text{eV}$ .  $T_c=1043\text{ K}$ Spin energy/J.  $J=0.04\text{eV}$ .  $T_c=1043\text{ K}$ 

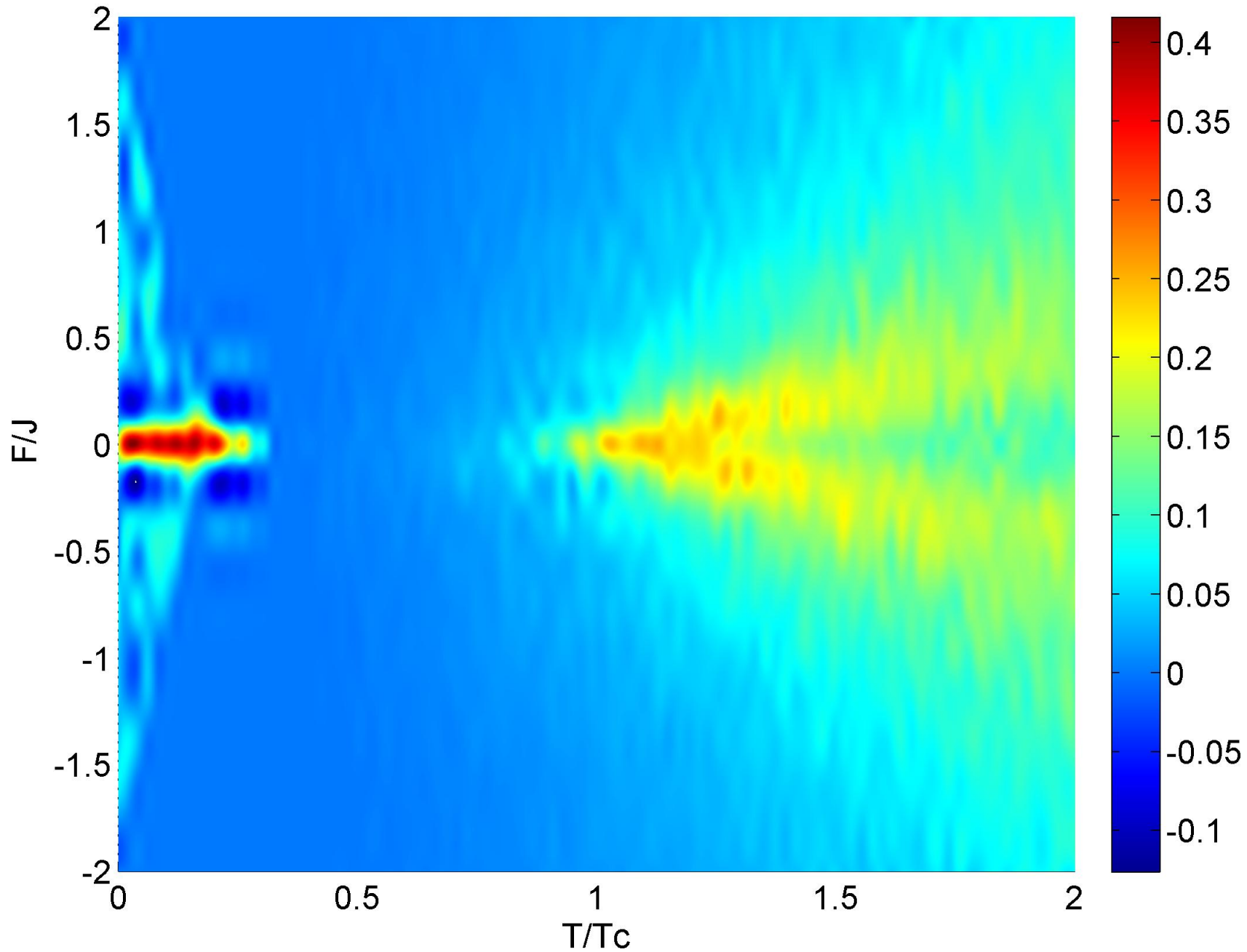
Heat capacity.  $J=0.04\text{eV}$ .  $T_c=1043\text{ K}$ Heat capacity.  $J=0.04\text{eV}$ .  $T_c=1043\text{ K}$ Heat capacity.  $J=0.04\text{eV}$ .  $T_c=1043\text{ K}$ Heat capacity.  $J=0.04\text{eV}$ .  $T_c=1043\text{ K}$ 

Mean abs spin  $N=10$ ,  $R=100$ ,  $I=2000$ ,  $T_c=1043\text{K}$ ,  $J=0.039644\text{eV}$

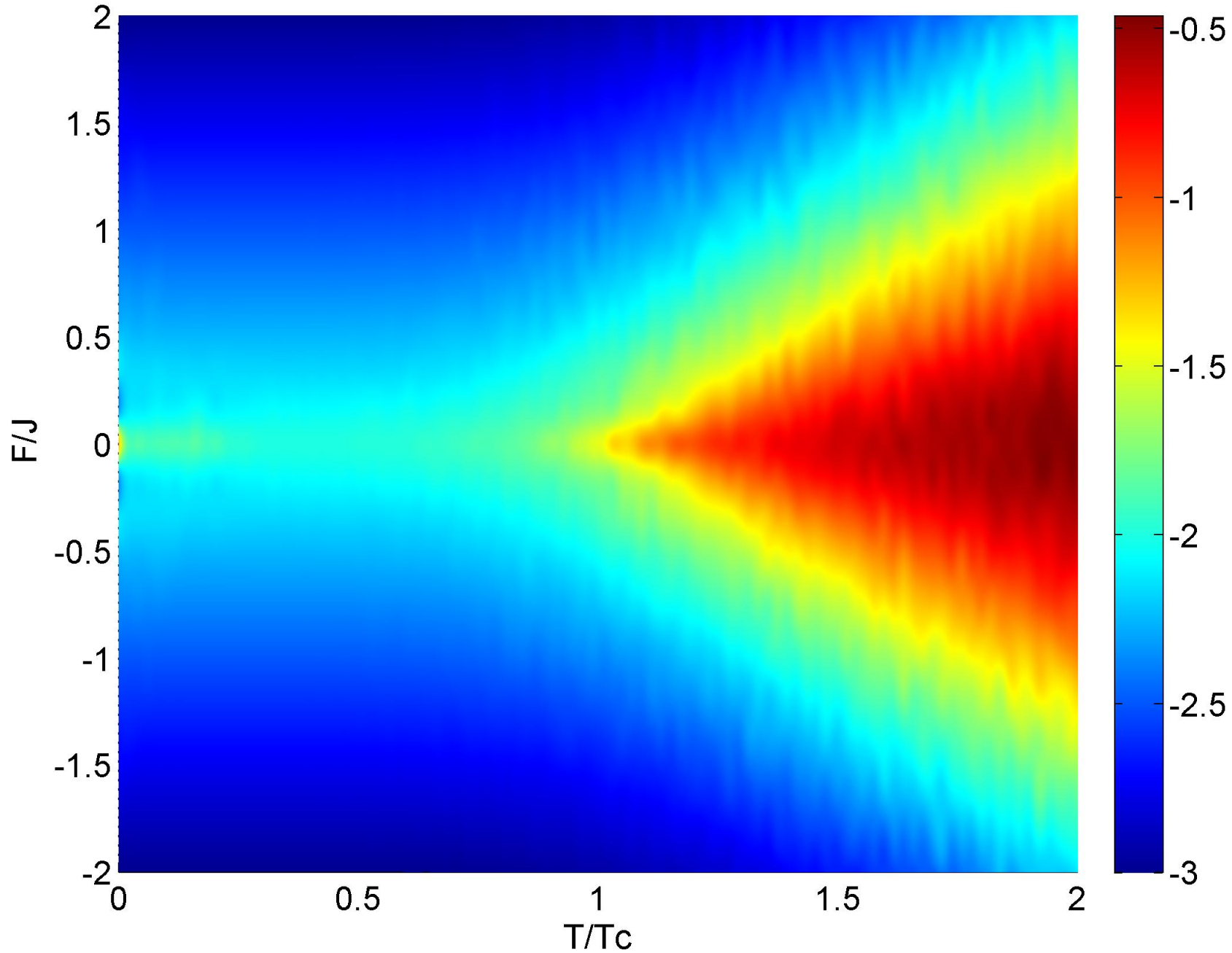




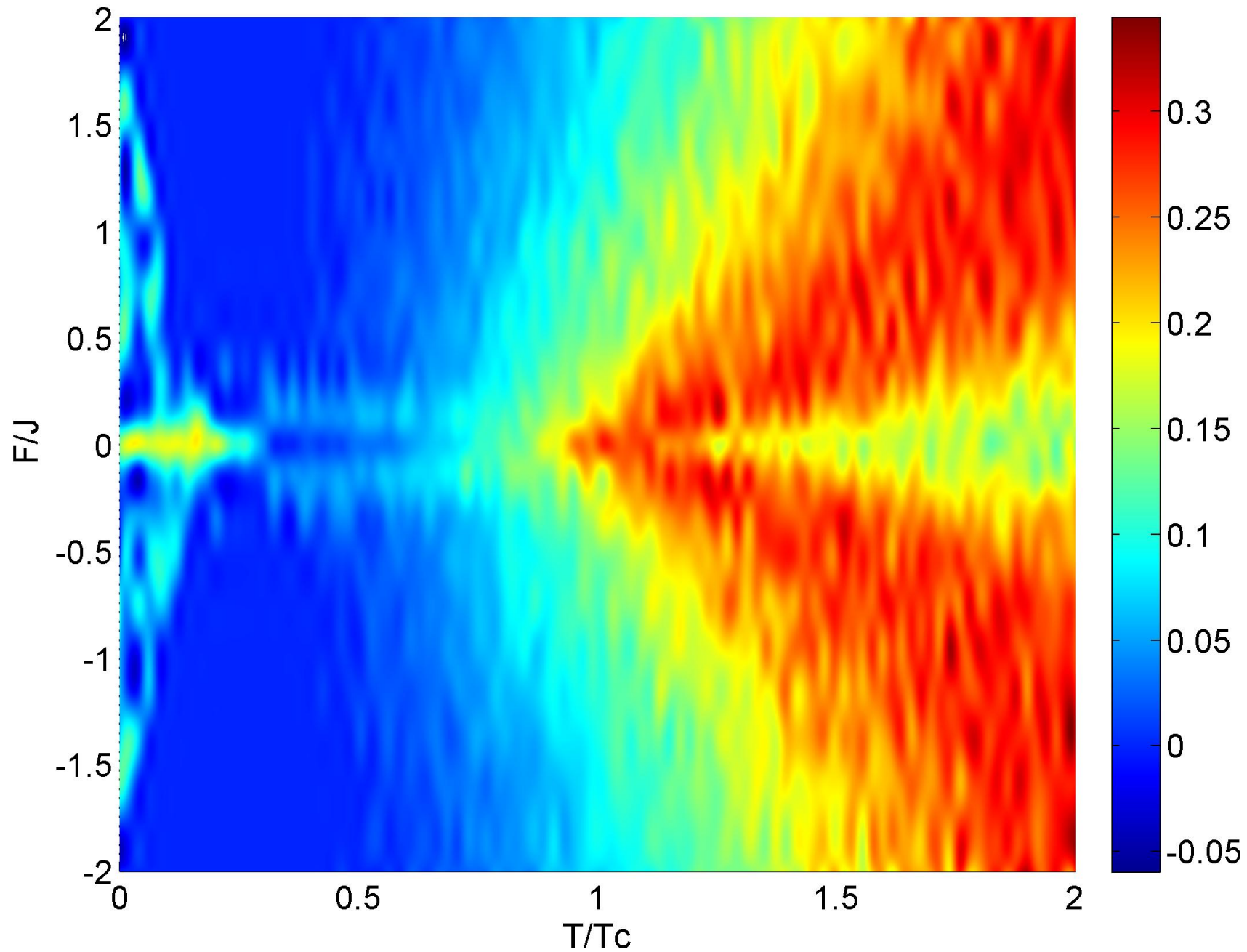
Mean abs spin sd N=10, R=100, I=2000, Tc=1043K, J=0.039644eV



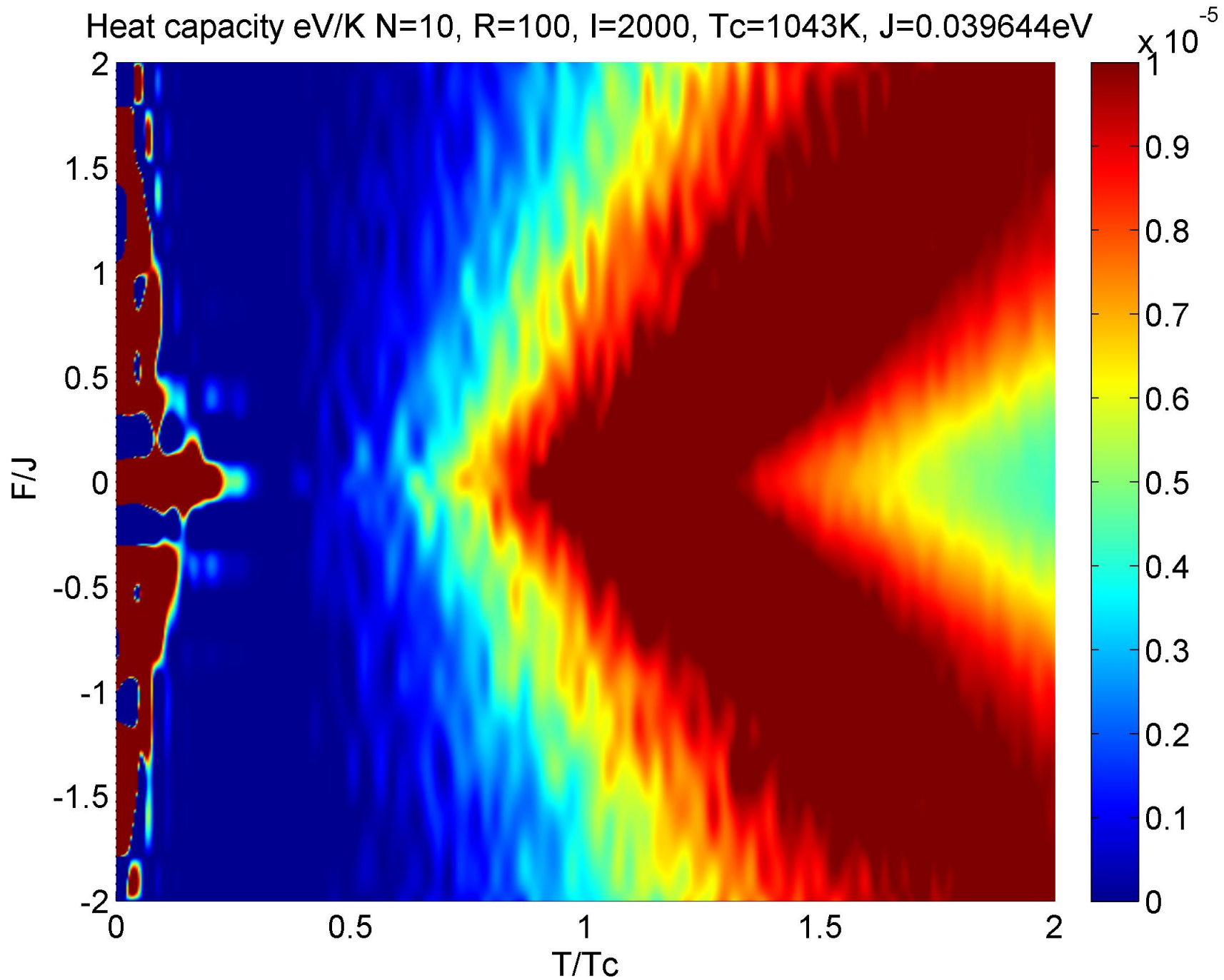
$E_{\text{mean}}(\text{eV})/J(\text{eV})$   $N=10$ ,  $R=100$ ,  $I=2000$ ,  $T_c=1043\text{K}$ ,  $J=0.039644\text{eV}$



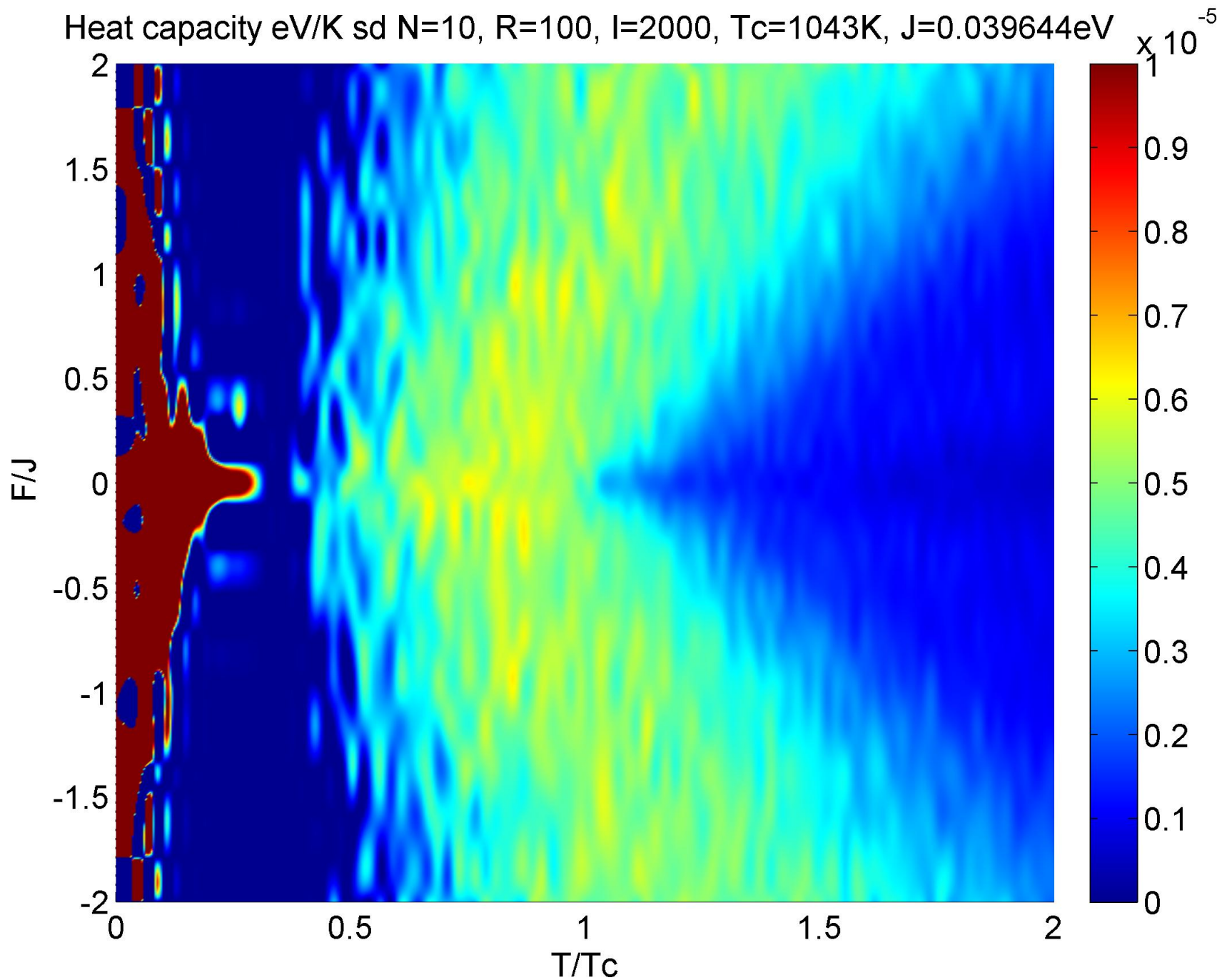
Em/J sd N=10, R=100, I=2000, Tc=1043K, J=0.039644eV



Heat capacity eV/K N=10, R=100, I=2000, Tc=1043K, J=0.039644eV



Heat capacity eV/K sd N=10, R=100, I=2000, Tc=1043K, J=0.039644eV



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[https://en.wikipedia.org/wiki/Ising\\_model](https://en.wikipedia.org/wiki/Ising_model)