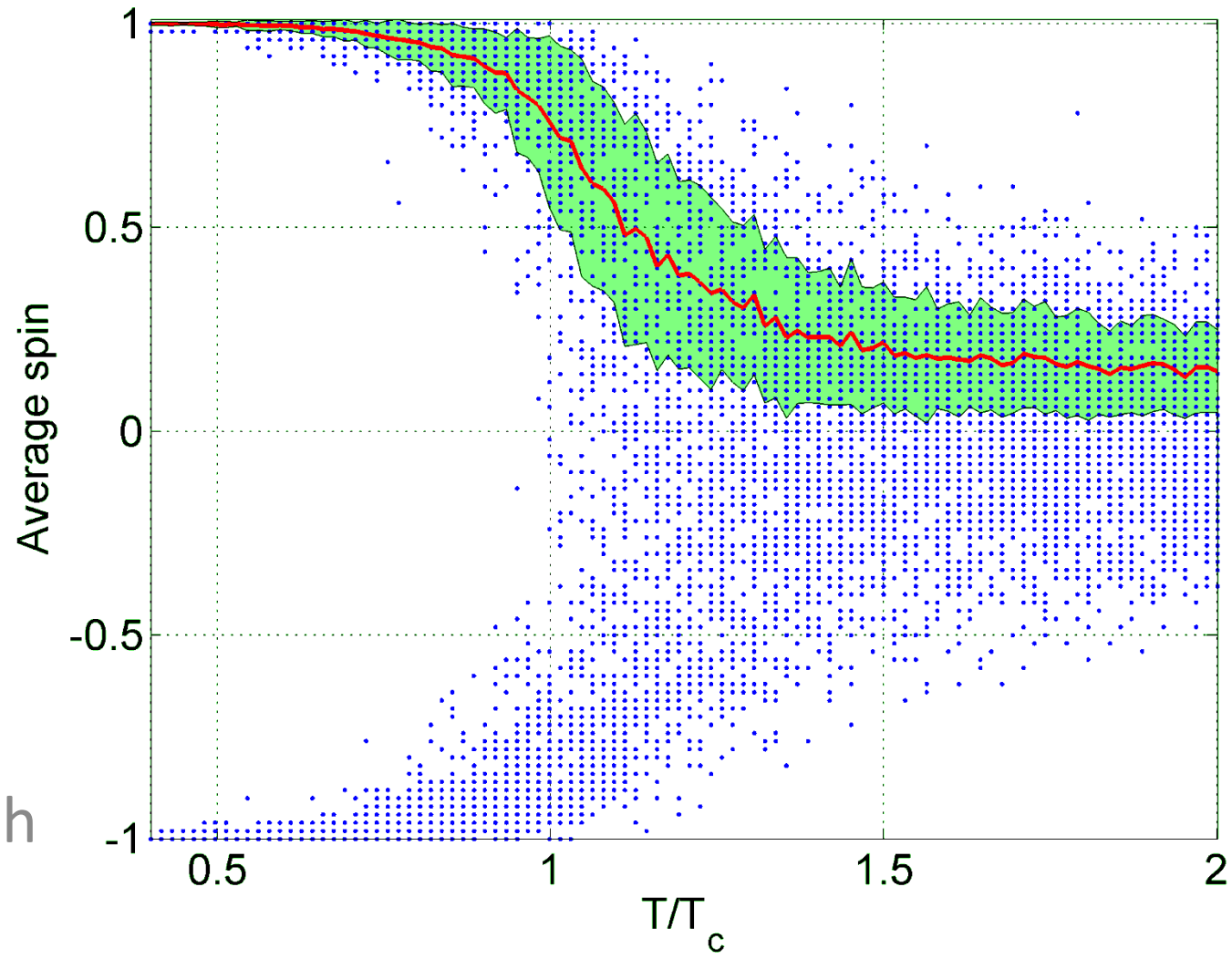


The Ising model of Ferromagnetism

Average spin. $J=0.04\text{eV}$. $T_c=1043\text{ K}$



The Ising Model of Ferromagnetism

All atoms will respond in some fashion to **magnetic fields**. The angular momentum (and spin) properties of electrons imply a circulating charge, which means they will be subject to a Lorentz force in a magnetic field. **However the effects of *diamagnetism*, *paramagnetism* and *anti-ferromagnetism* are typically very small.** **Ferromagnetic materials** (iron, cobalt, nickel, some rare earth metal compounds) respond strongly to magnetic fields and can intensify them by orders of magnitude. i.e. the *relative permeability* can be tens or hundreds, or possibly thousands.

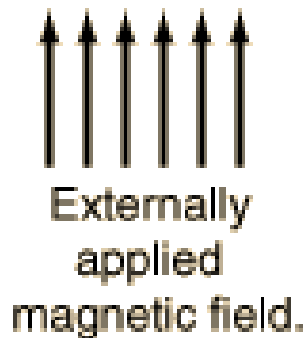
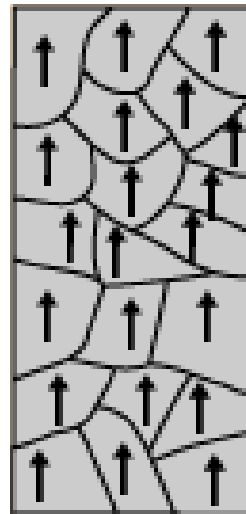
The Ising model is a simplified model of a **ferromagnet** which exhibits a **phase transition** above the **Curie temperature**. Below this, magnetic dipole alignment will tend to cluster into **domains**, and it is these micro-scale groupings which give rise to ferromagnetic behaviour.



Ernst Ising (1900-1998)

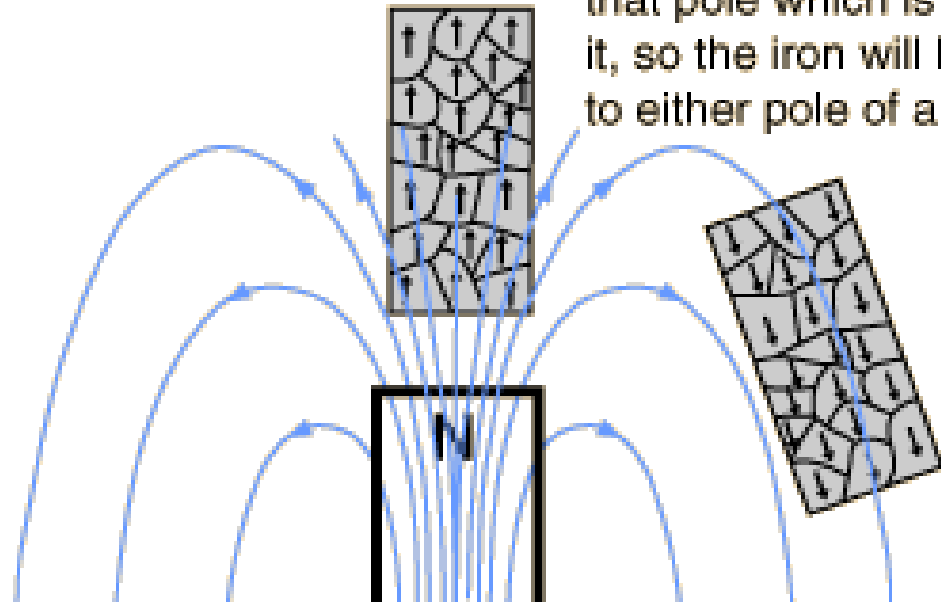


“Soft” magnetism - Ferromagnets



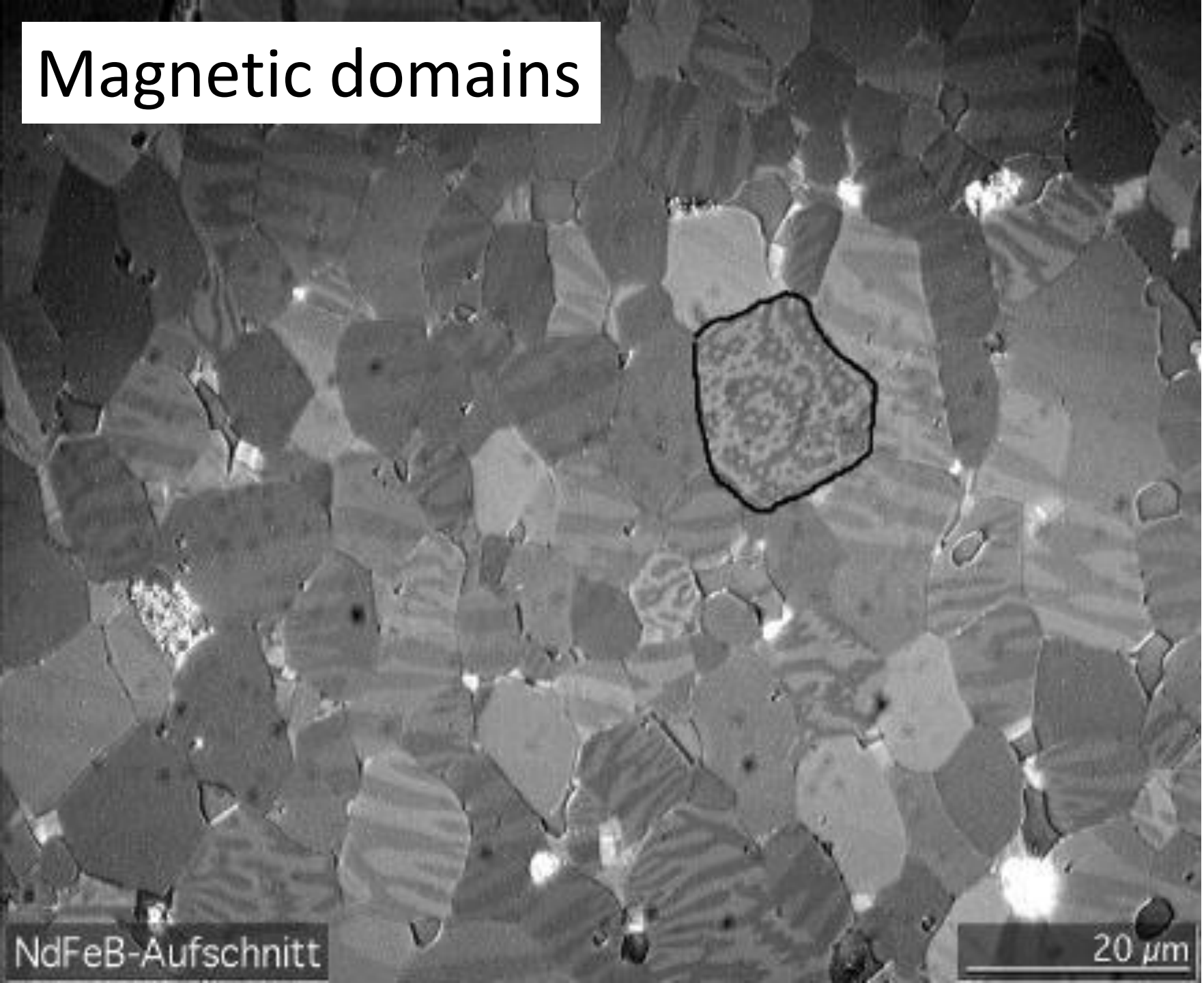
Externally
applied
magnetic field.

Iron will become magnetized in the direction of any applied magnetic field. This magnetization will produce a magnetic pole in the iron opposite to that pole which is nearest to it, so the iron will be attracted to either pole of a magnet.



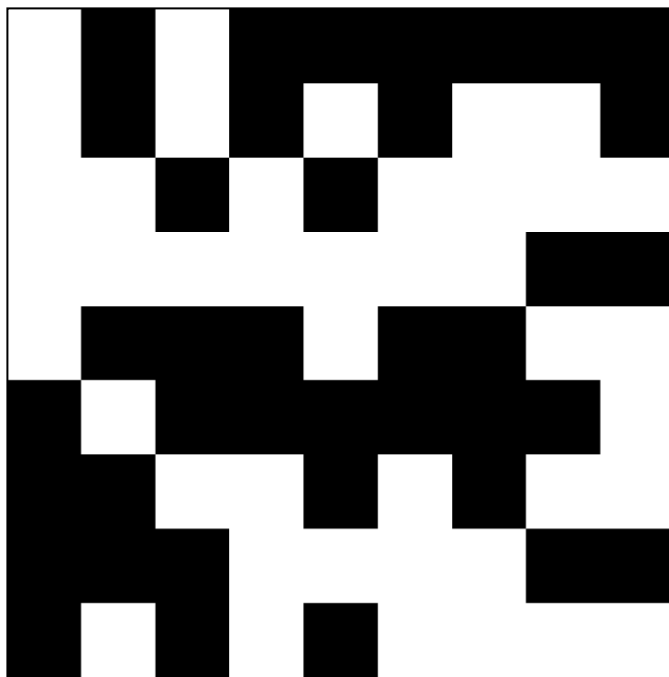
Unlike permanent “hard” magnets, once the applied field is removed, the domain alignment will randomize again, effectively zeroing the net magnetism.

Magnetic domains



The **Ising model** can be used to demonstrate spontaneous mass alignment of magnetic dipoles, and possibly a mechanism for domain formation.

Perhaps the simplest model which yields characteristic behaviour is an $N \times N$ square grid, where each square is initially randomly assigned a +1 or -1 value, with equal probability. The +/-1 values correspond to a single direction of magnetic dipole moment in a rectangular lattice of ferromagnetic atoms, or in the case of individual electrons, *spin*.



10 x 10 grid

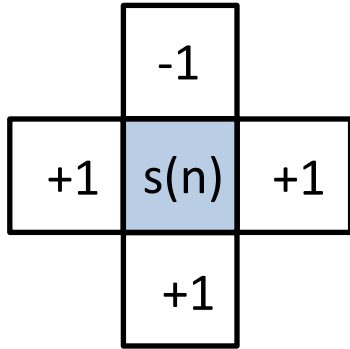
White squares
represent +1
Black squares
represent -1



100 x 100 grid

Original Metropolis algorithm

1. Choose one square at random from the $N \times N$ grid. Let its spin be $s(n) = +1$ or -1 .
2. Find the spins of the nearest neighbours. Use *circular boundary conditions* e.g. if $s(n)$ is at the edge of the grid, use the nearest neighbour to be that of the other end.



3. Compute a sum of **spin-coupling energies** for $s(n)$ and its neighbours, and work out the energy change if $s(n)$ were to **change sign**

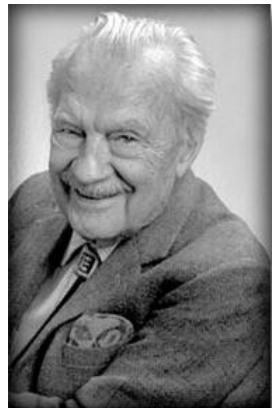
$$\Delta E = 2 \times \left(F + J \sum_{k=1}^4 s_n(k) \right) s(n)$$

J is the spin coupling energy in eV and F is the energy in eV associated with the alignment of spin $s(n)$ with an applied external magnetic field. Let us ignore any energy contributions from non-nearest neighbours.

$$r \sim U(0,1)$$

Now change the sign of spin $s(n)$ according to the following rule:

$$s(n) \rightarrow -s(n) \quad \text{if} \quad e^{-\frac{\Delta E}{k_B T}} \geq r \quad \text{or} \quad \Delta E < 0$$



Nicholas
Metropolis
1915-1999

Apply the Metropolis method for $L \times N \times N$ iterations, and then compute from the $N \times N$ grid the following parameters

$$\langle s \rangle = \frac{1}{N^2} \sum_{n=1}^{N^2} s(n) \quad \text{Mean spin}$$

$$\langle E \rangle = -\frac{1}{2} \frac{1}{N^2} \sum_{n=1}^{N^2} \left(s(n) J \sum_{k=1}^4 s_n(k) + F \right) \quad \text{Mean energy per spin}$$

$$k_B T^2 \langle C \rangle = \frac{1}{4} \frac{1}{N^2} \sum_{n=1}^{N^2} \left(s(n) J \sum_{k=1}^4 s_n(k) + F \right)^2 - \langle E \rangle^2$$

Heat capacity in eV per K

This is a well known result in Statistical Thermodynamics

$$k_B T^2 \langle C \rangle = \text{Var}[E]$$

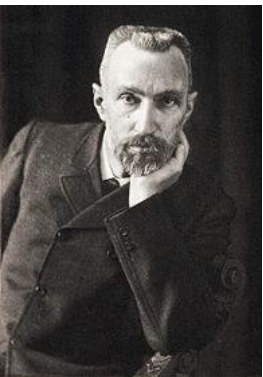
For a 2D Ising model, Lars Onsager determined in 1944 the relationship between the phase transition Curie temperature and coupling energy J

$$\frac{2J}{\ln\left(1 + \sqrt{2}\right)} = k_B T_C$$

(Note this expression assumes coupling energy J is in joules)

Boltzmann's constant

$$k_B = 1.38 \times 10^{-23} \text{ JK}^{-1}$$

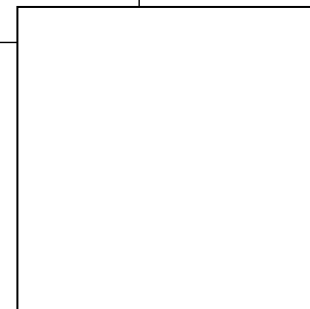
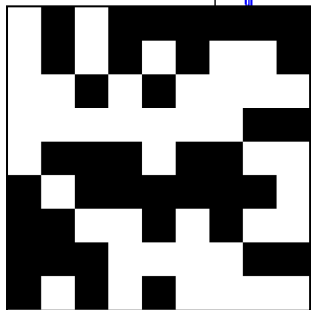


Peter Curie
(1859-1906)



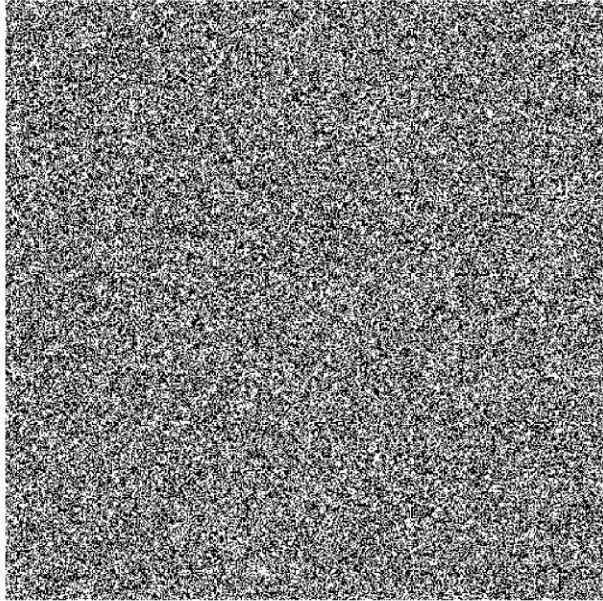
Lars Onsager
(1903-1976)

Mean spin vs iteration. $N=10$, $T/T_c=0.5$



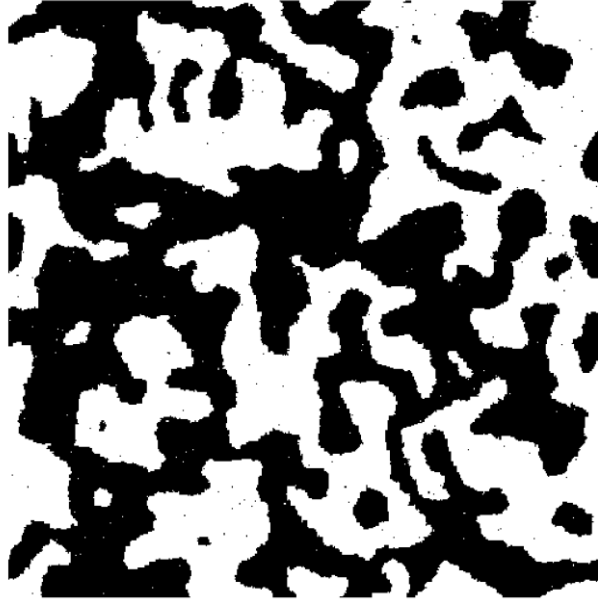
iteration = 1/2000

Mean spin=0.002568, $T/T_c=0.5$



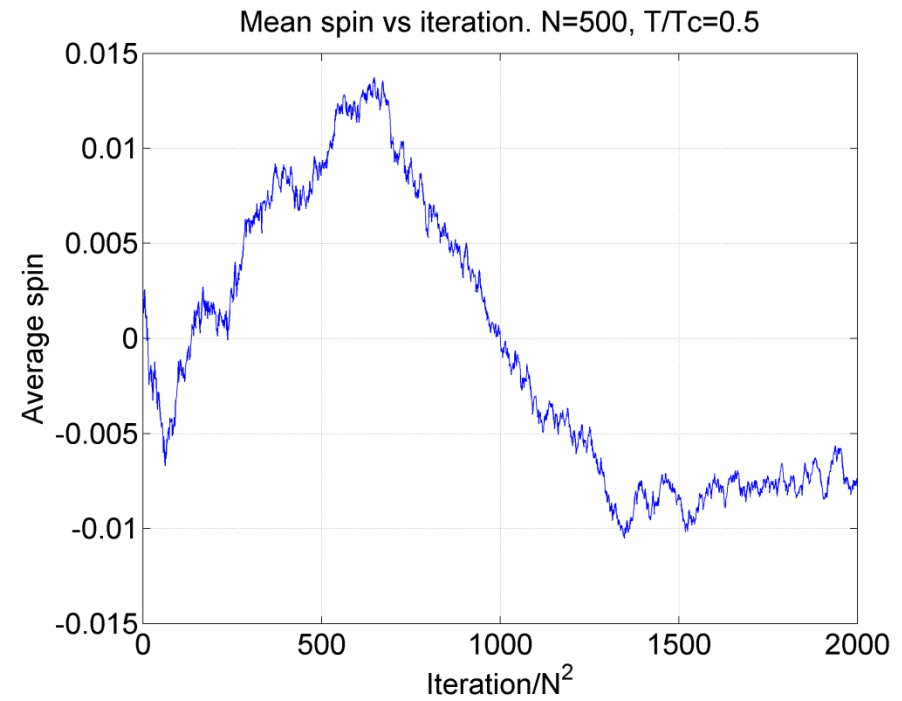
iteration = 2000/2000

Mean spin=-0.007728, $T/T_c=0.5$



For a 500 x 500 grid, a similar equilibrium is not yet reached, even after $I = 2000 \times 500 \times 500$ iterations.

However, domain-like structures are clearly visible in this intermediate state.



Results of a MATLAB simulation:

10 x 10 grid

$I = 2000$ (x 10 x 10) iterations of Metropolis algorithm

$R = 100$ repeats for each temperature

100 different temperatures from $T/T_c = 0.0 \dots 2.0$

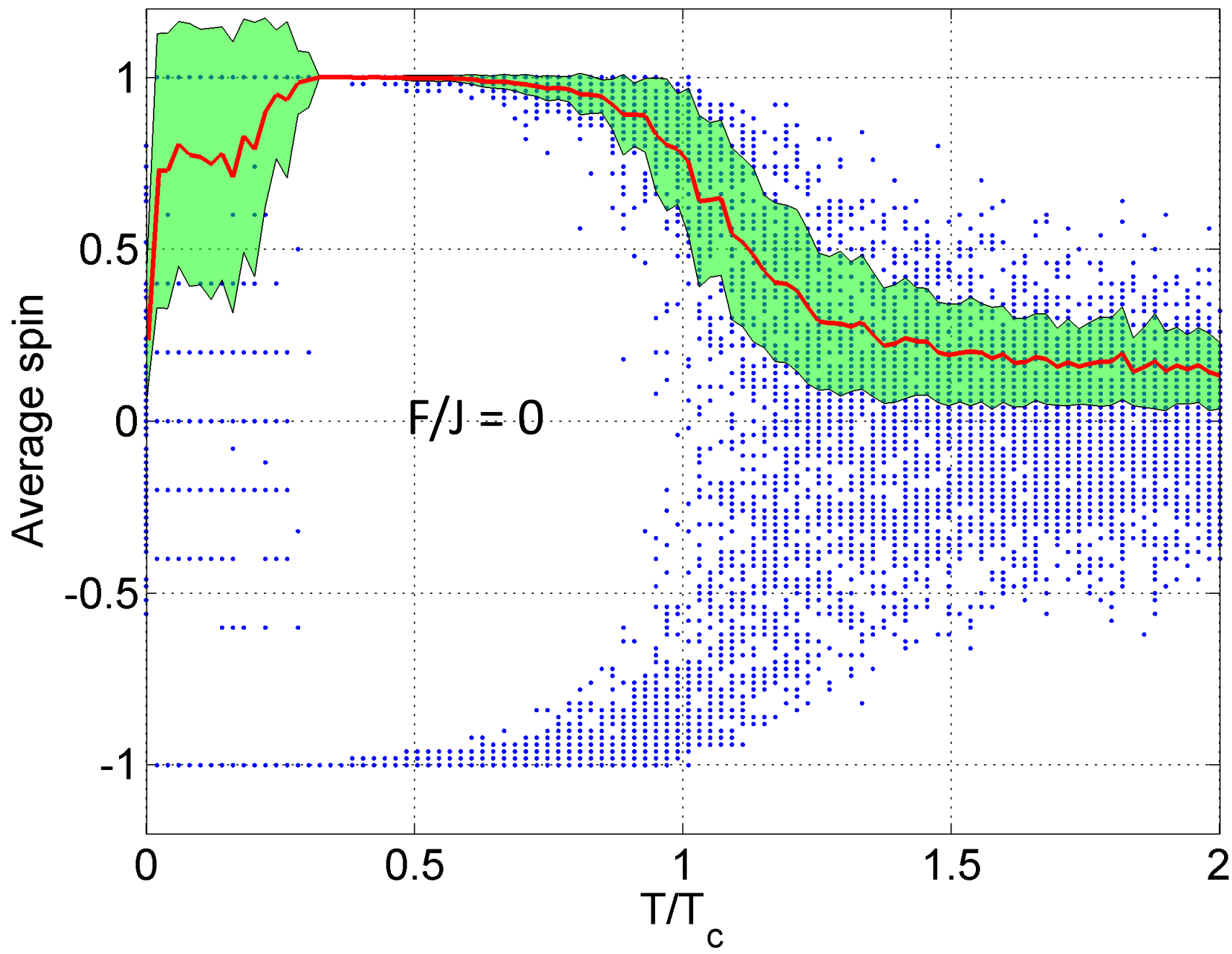
21 different F/J values from -2 to 2

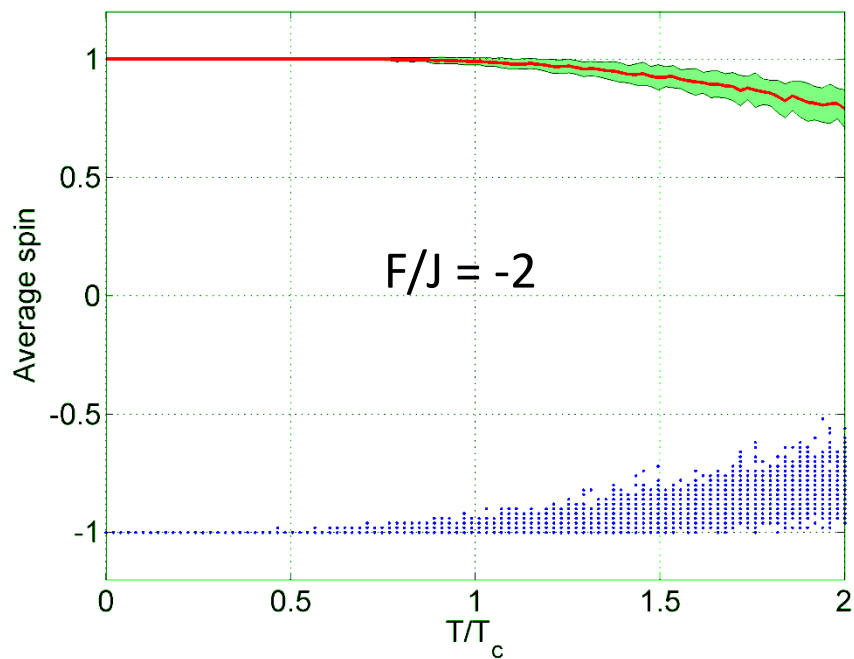
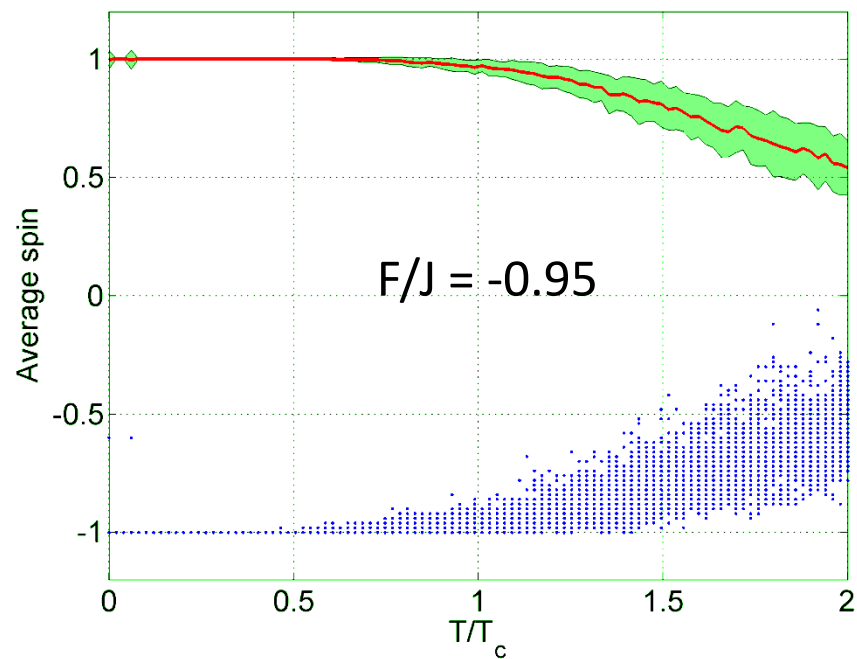
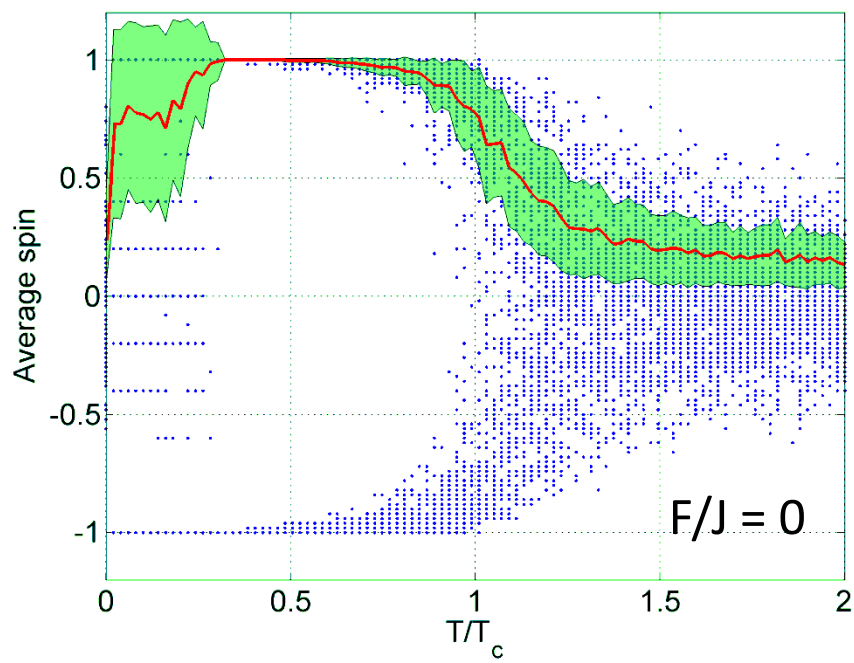
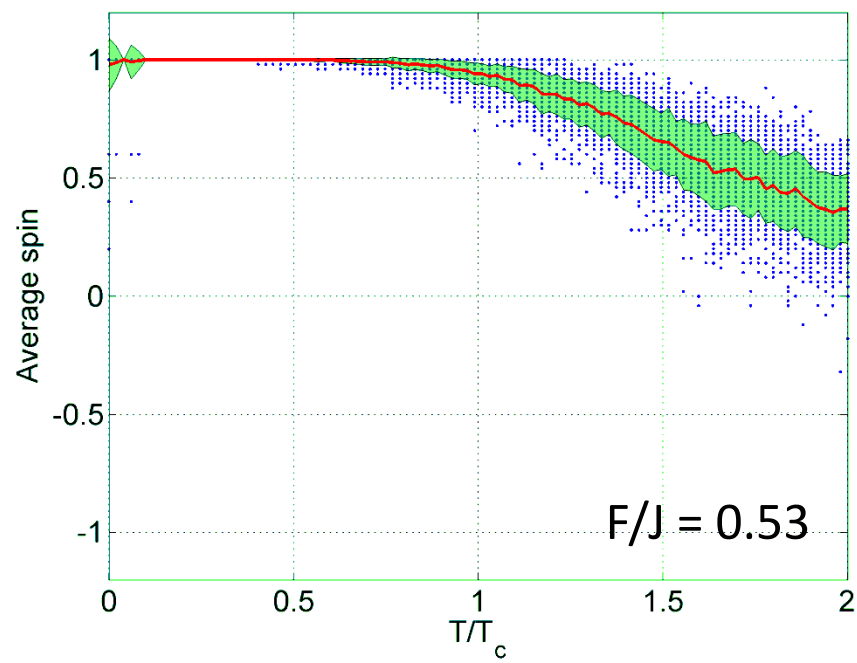
i.e. $2000 \times 10 \times 10 \times 100 \times 100 \times 21 =$

42 billion iterations of the Metropolis algorithm

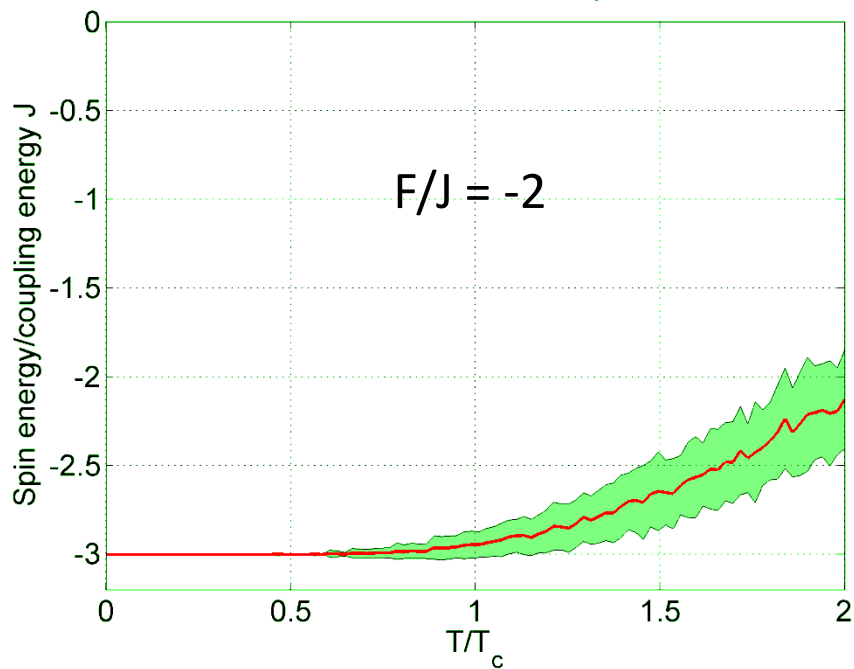
Running time on an i5 PC was about five days! Opportunity for *parallel processing*.

Average spin. $J=0.04\text{eV}$. $T_c=1043\text{ K}$

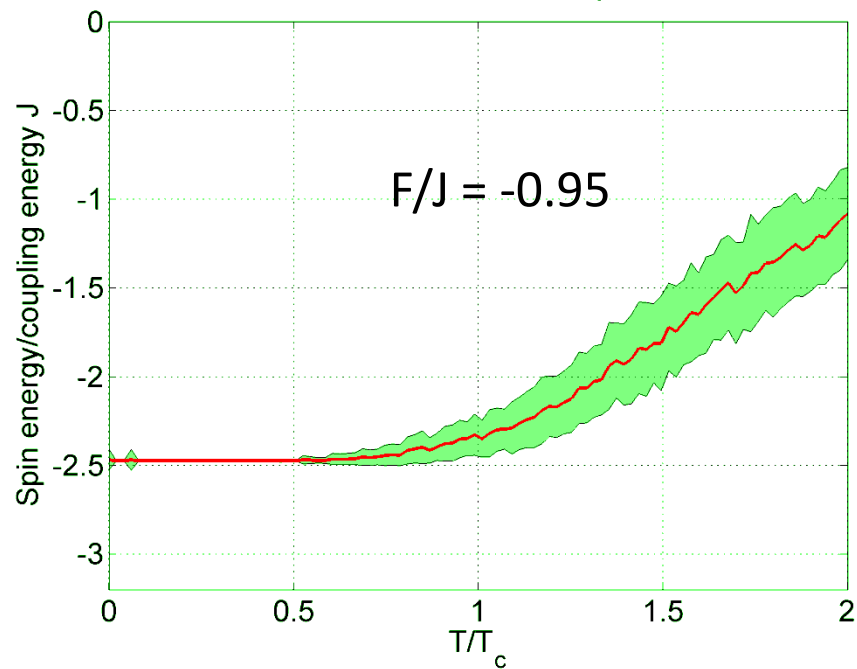


Average spin. $J=0.04\text{eV}$. $T_c=1043\text{ K}$ Average spin. $J=0.04\text{eV}$. $T_c=1043\text{ K}$ Average spin. $J=0.04\text{eV}$. $T_c=1043\text{ K}$ Average spin. $J=0.04\text{eV}$. $T_c=1043\text{ K}$ 

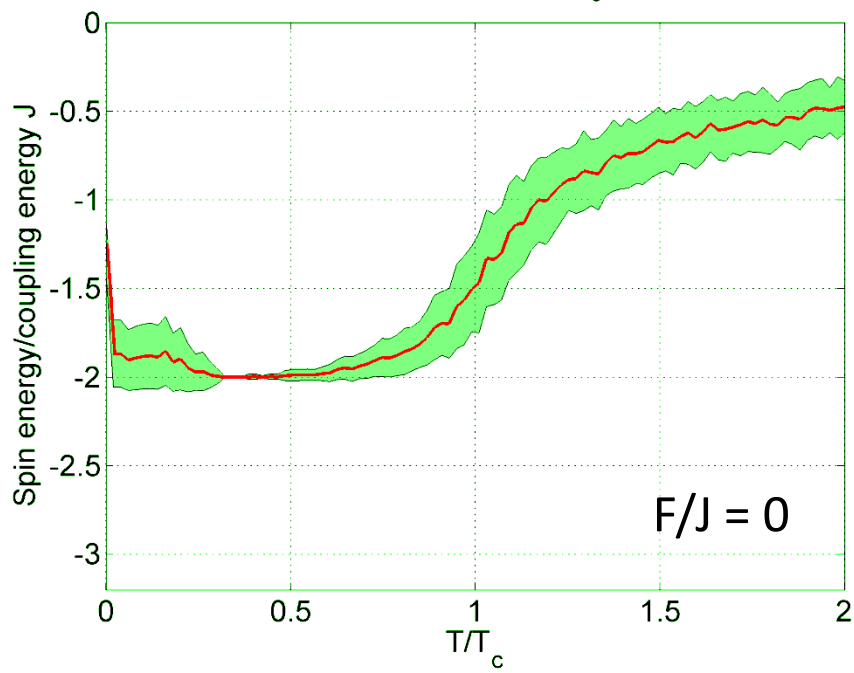
Spin energy/J. $J=0.04\text{eV}$. $T_c=1043\text{ K}$



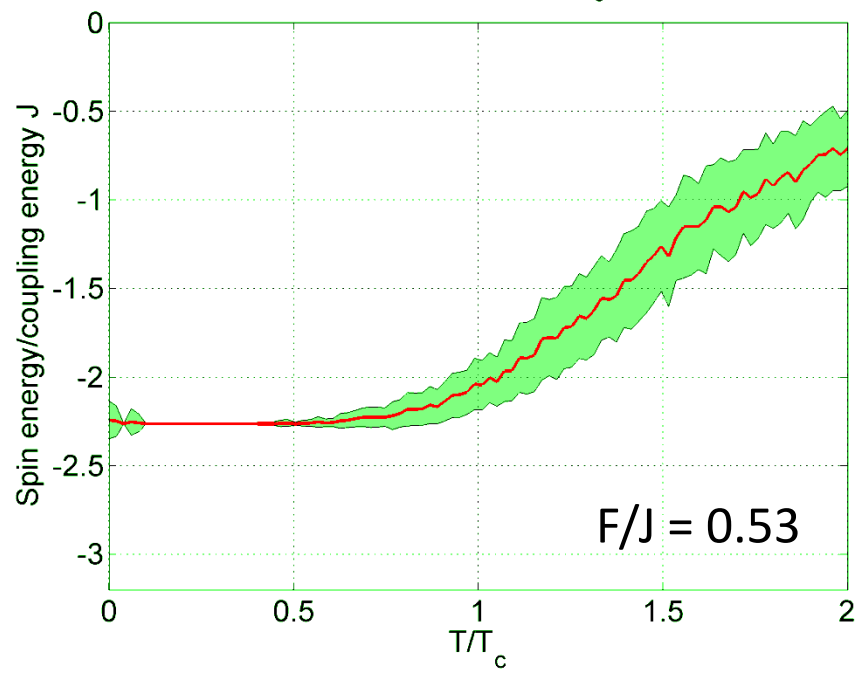
Spin energy/J. $J=0.04\text{eV}$. $T_c=1043\text{ K}$

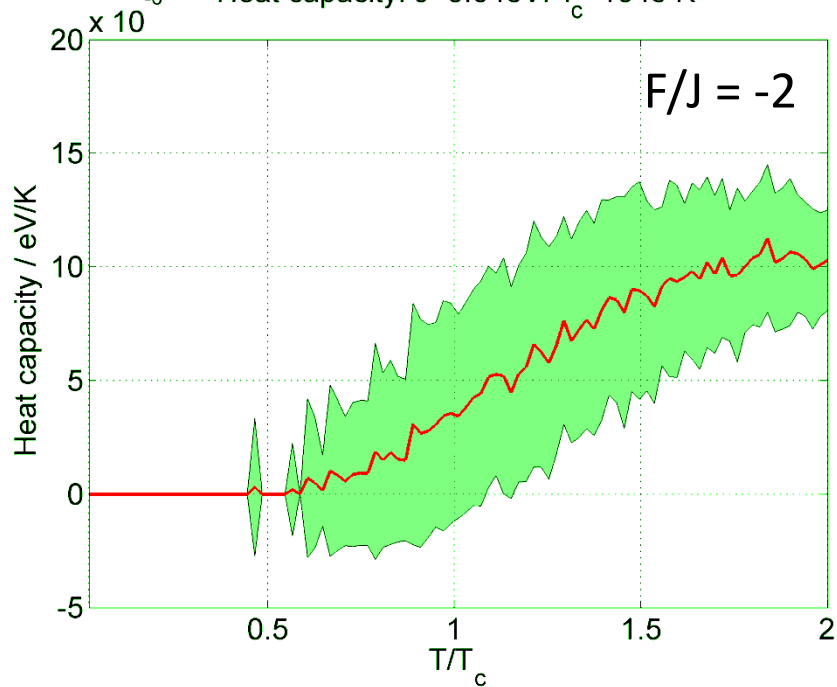
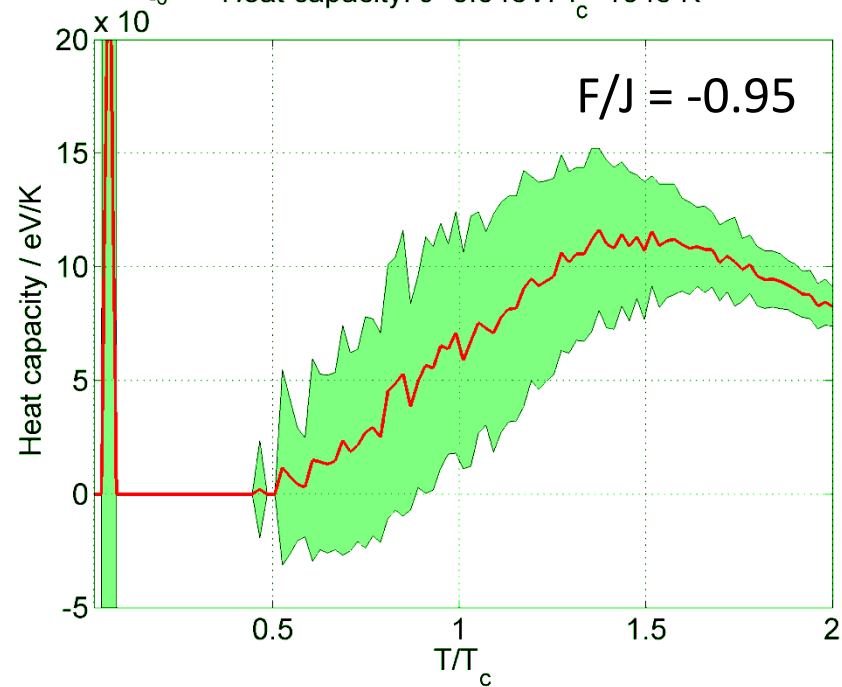
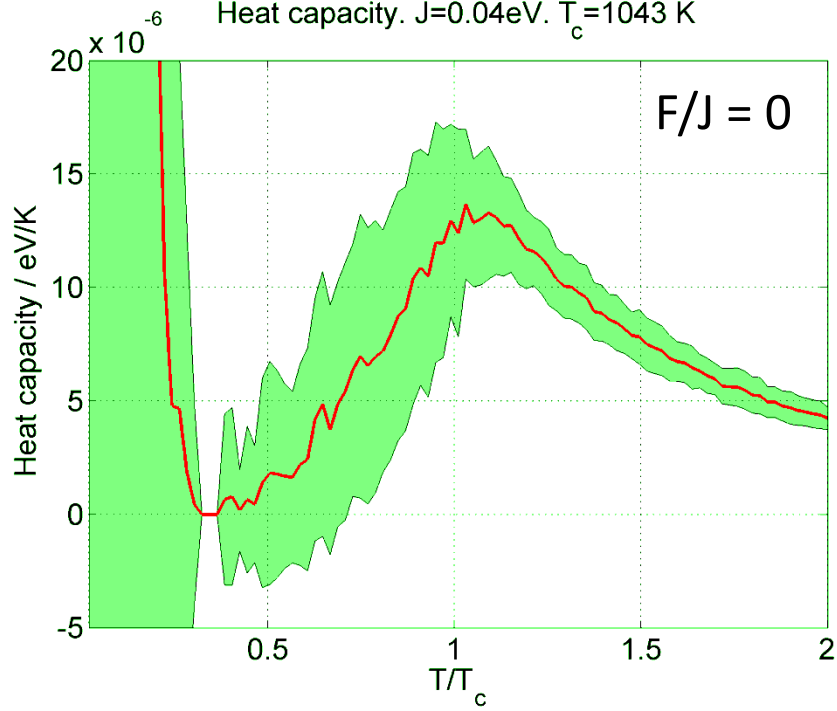
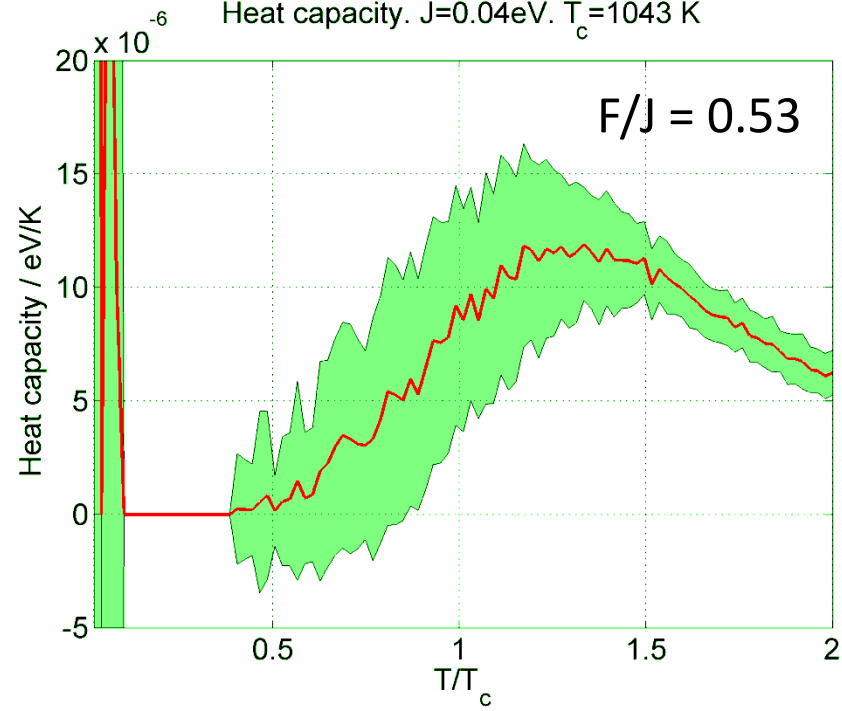


Spin energy/J. $J=0.04\text{eV}$. $T_c=1043\text{ K}$

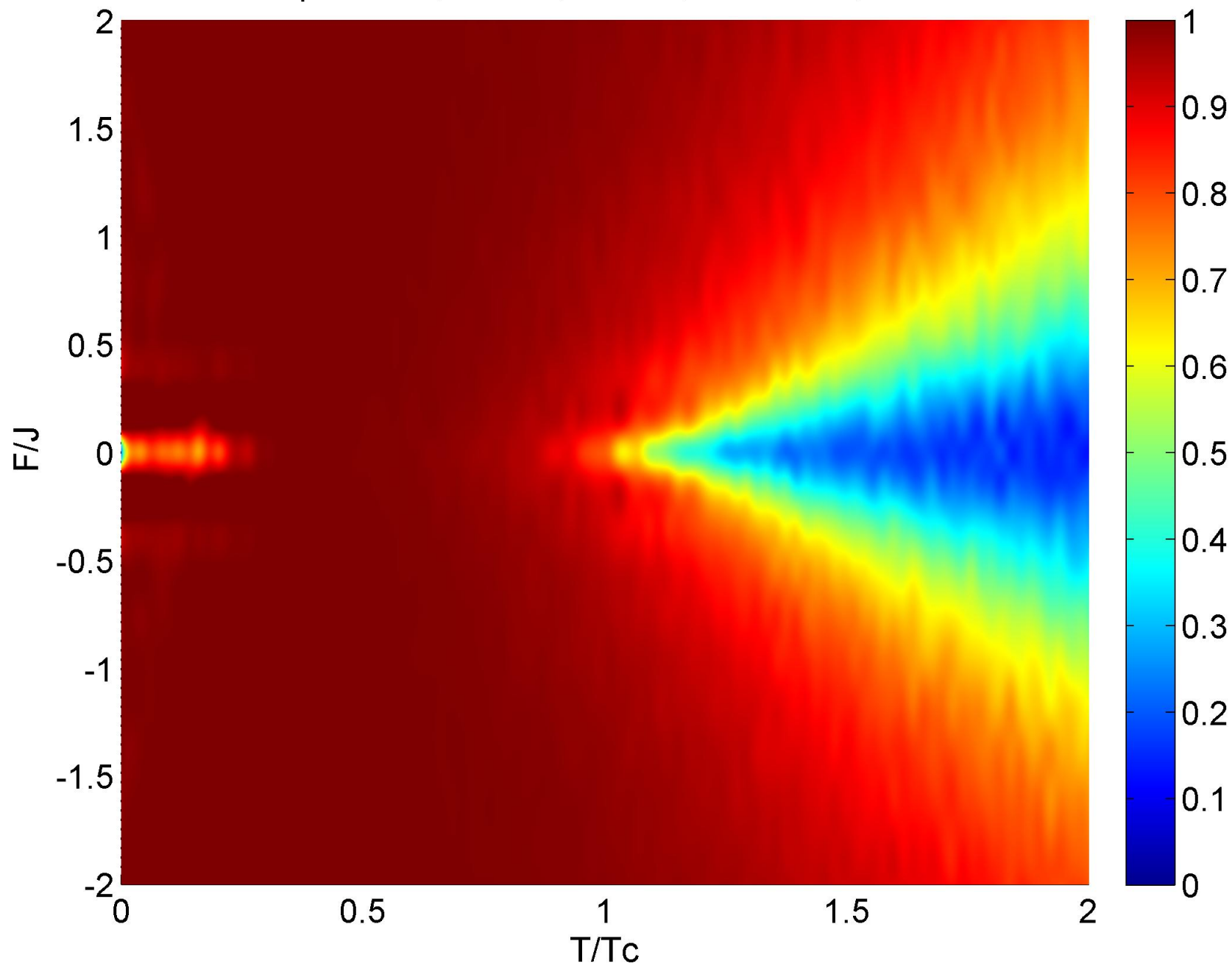


Spin energy/J. $J=0.04\text{eV}$. $T_c=1043\text{ K}$

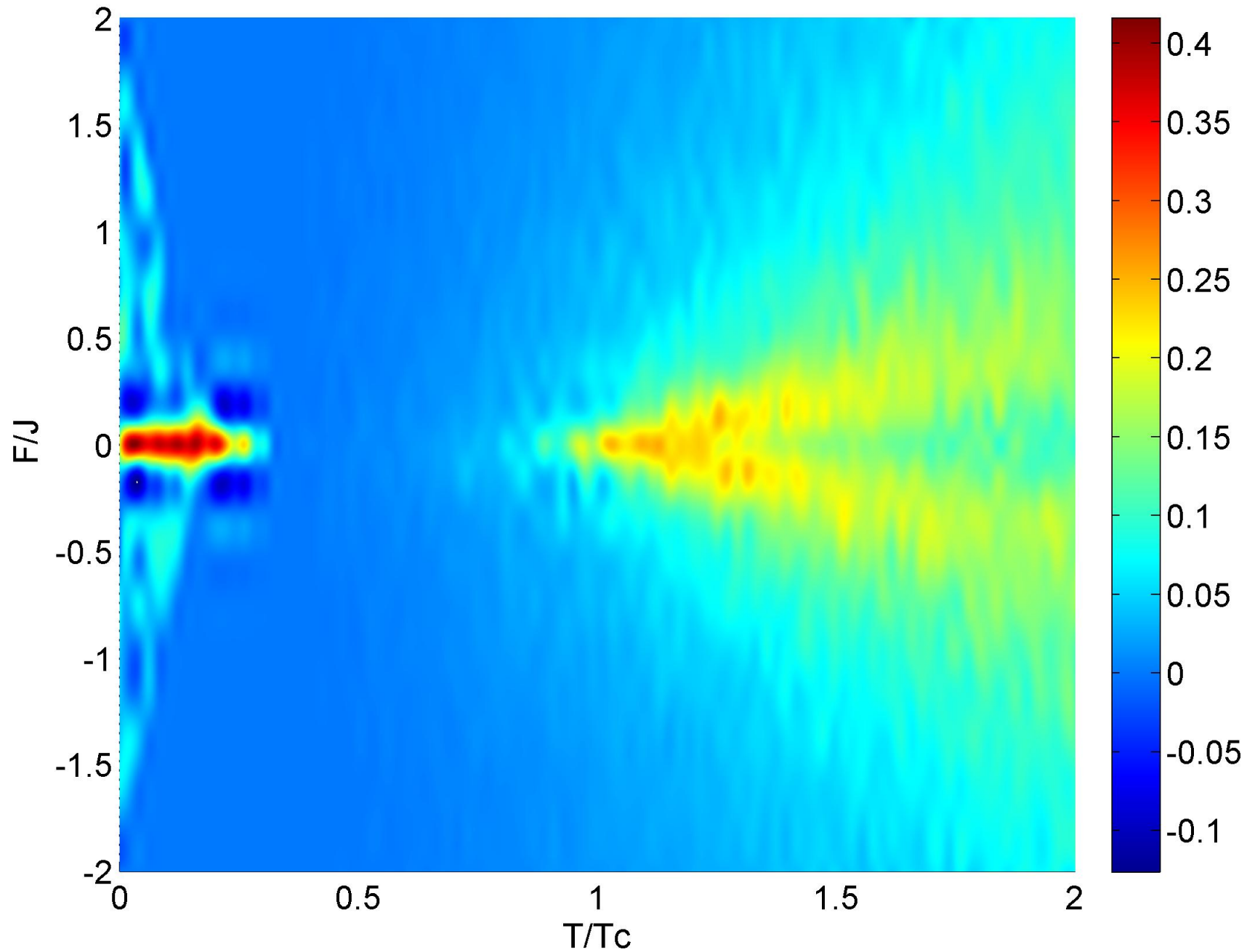


Heat capacity. $J=0.04\text{eV}$. $T_c=1043\text{ K}$ Heat capacity. $J=0.04\text{eV}$. $T_c=1043\text{ K}$ Heat capacity. $J=0.04\text{eV}$. $T_c=1043\text{ K}$ Heat capacity. $J=0.04\text{eV}$. $T_c=1043\text{ K}$ 

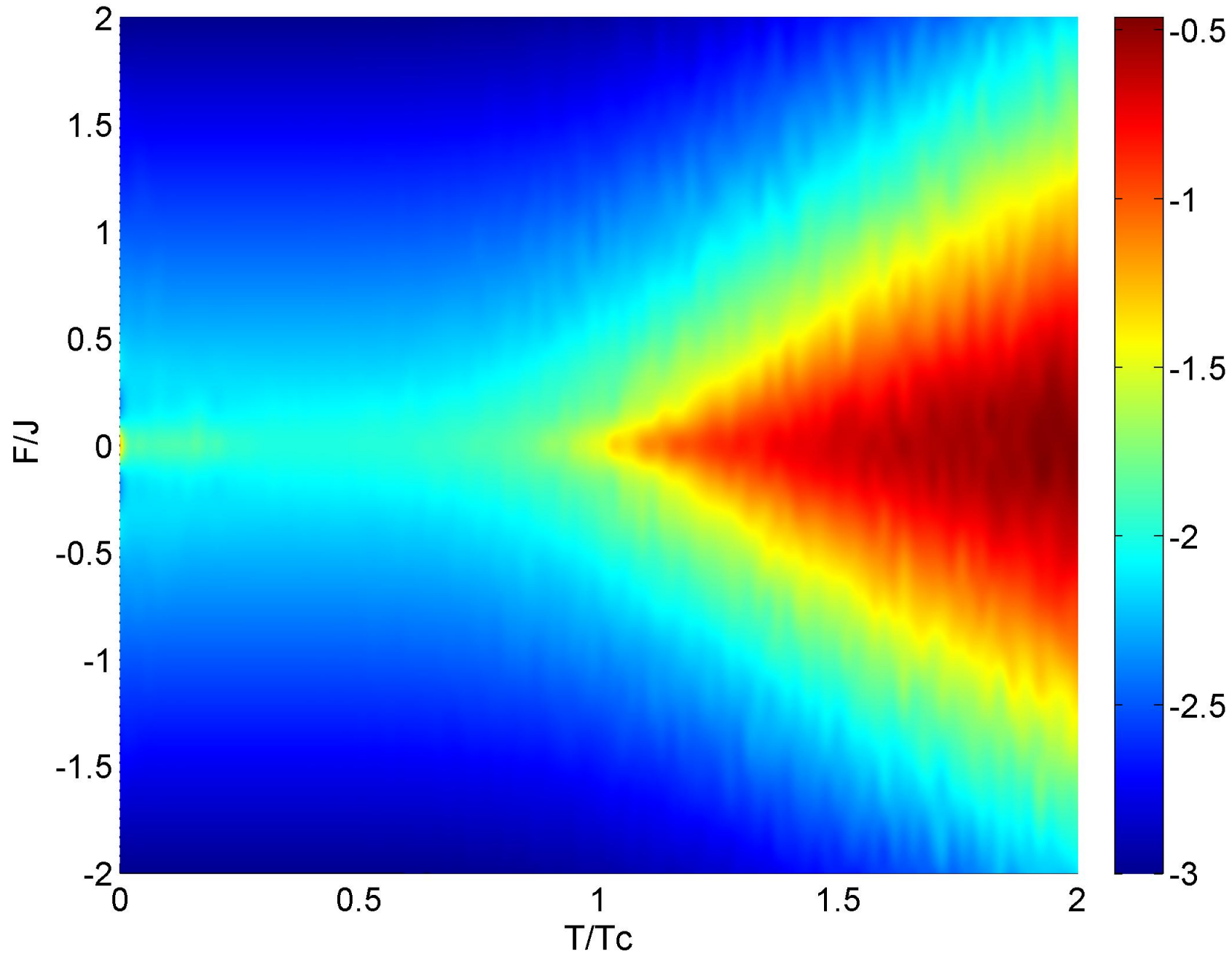
Mean abs spin $N=10$, $R=100$, $I=2000$, $T_c=1043\text{K}$, $J=0.039644\text{eV}$



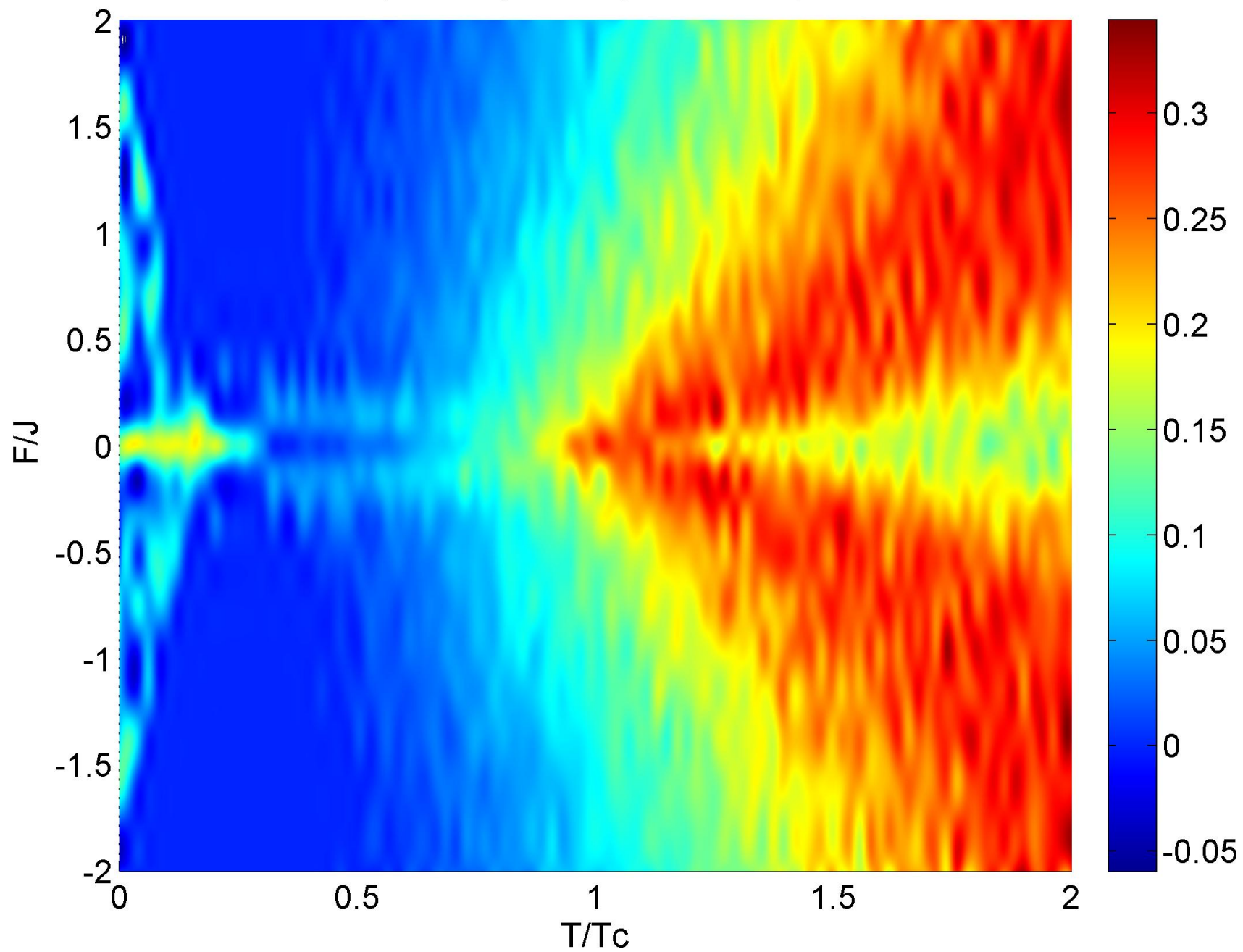
Mean abs spin sd N=10, R=100, I=2000, Tc=1043K, J=0.039644eV



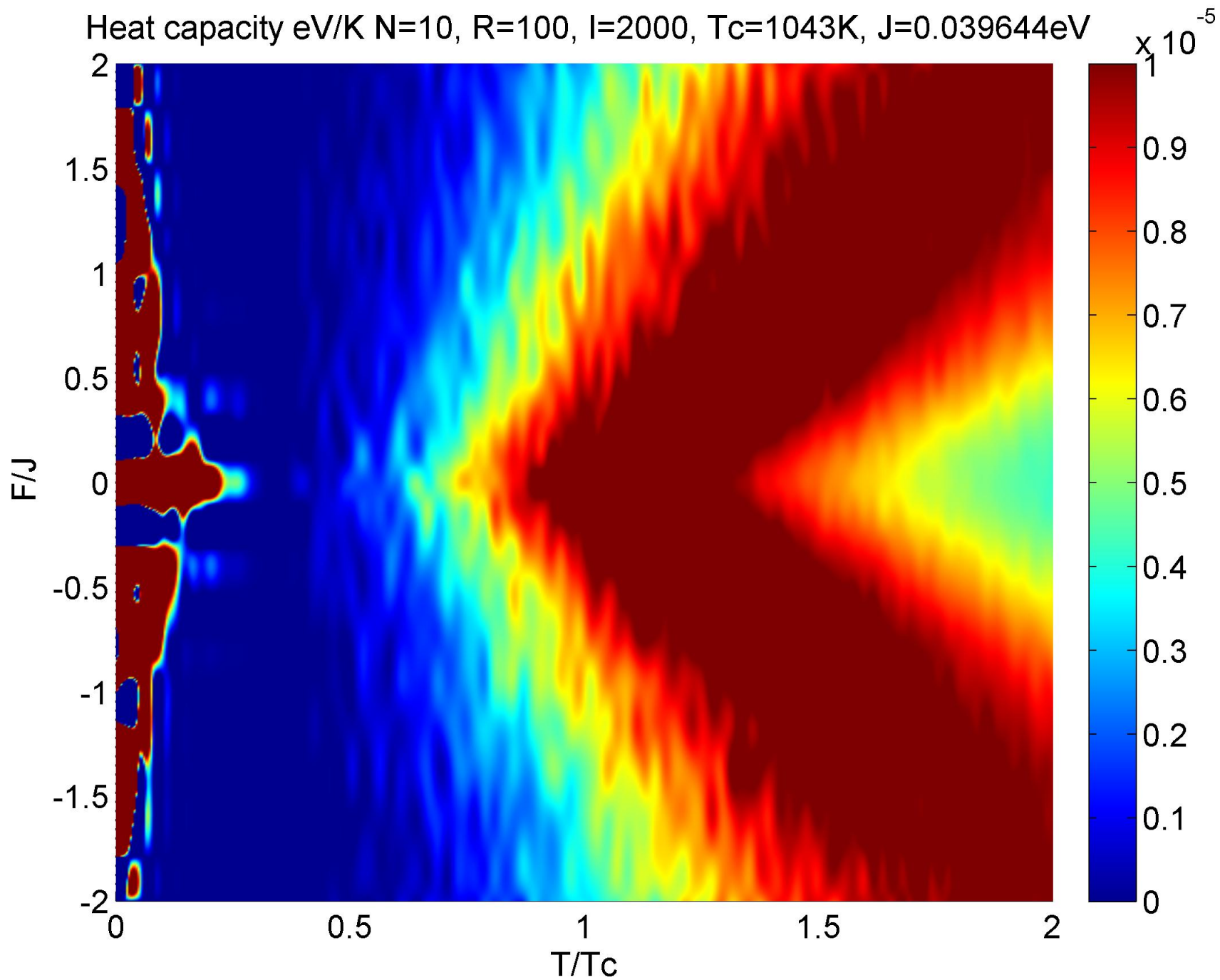
E_{mean}(eV)/J(eV) N=10, R=100, I=2000, T_c=1043K, J=0.039644eV



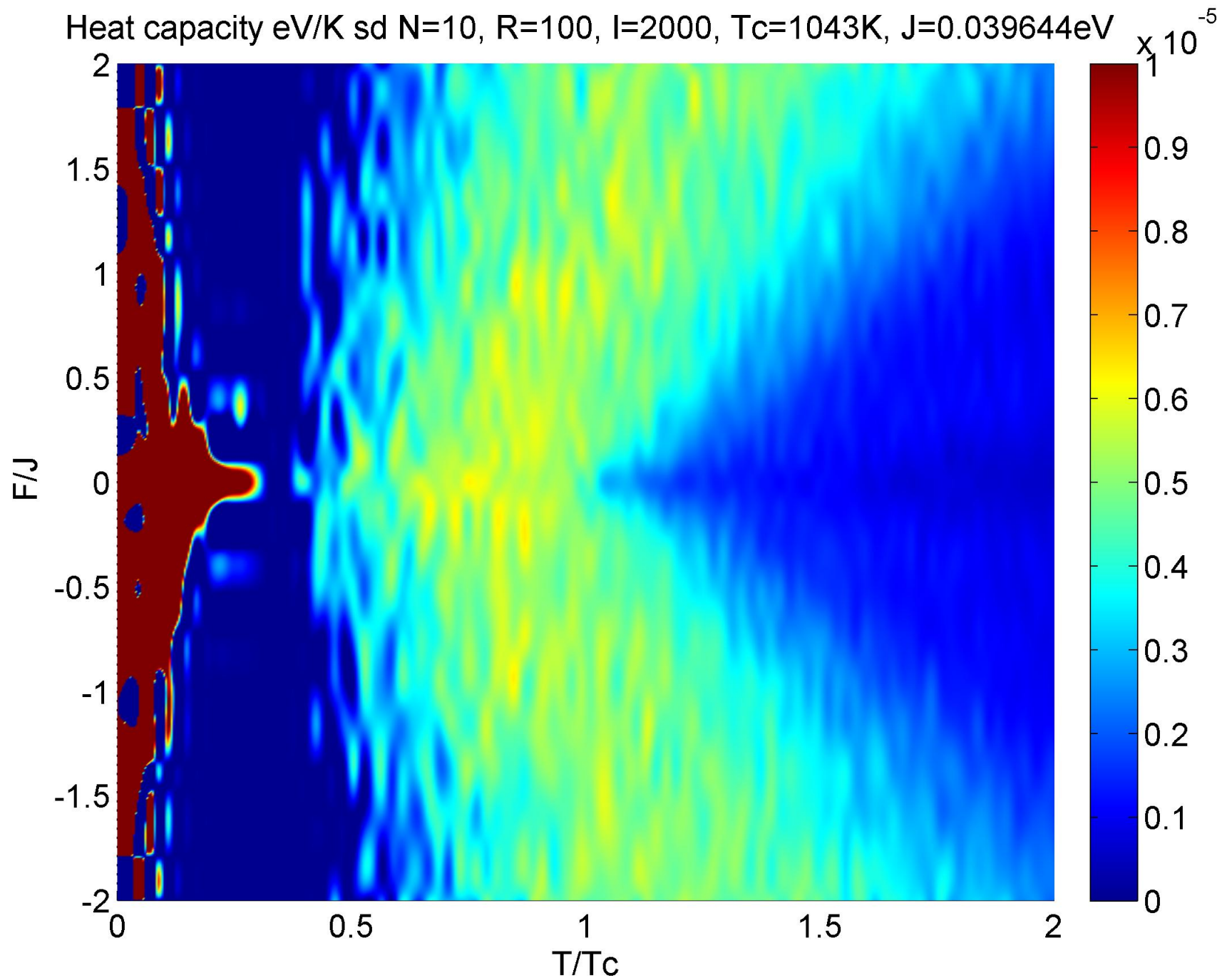
Em/J sd N=10, R=100, l=2000, T_c=1043K, J=0.039644eV



Heat capacity eV/K $N=10$, $R=100$, $I=2000$, $T_c=1043\text{K}$, $J=0.039644\text{eV}$



Heat capacity eV/K sd N=10, R=100, I=2000, T_c=1043K, J=0.039644eV



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