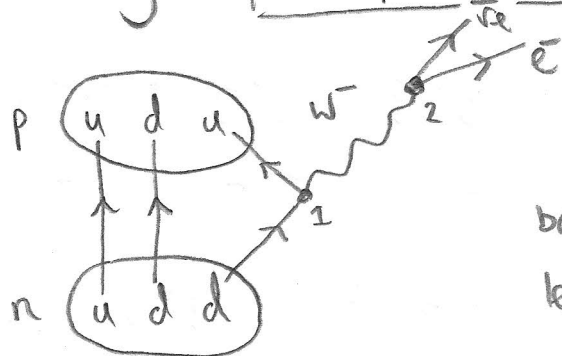


# HIGH ENERGY & PARTICLE PHYSICS

(i)  $\beta^-$  decay:

$$n \rightarrow p + e^- + \bar{\nu}_e$$



charge:  $-\frac{1}{3}e \rightarrow \frac{2}{3}e - 1e \quad \checkmark$   
 baryon #:  $\frac{1}{3} \rightarrow \frac{1}{3} + 0 \quad \checkmark$   
 lepton #:  $0 \rightarrow 0 + 0 \quad \checkmark$

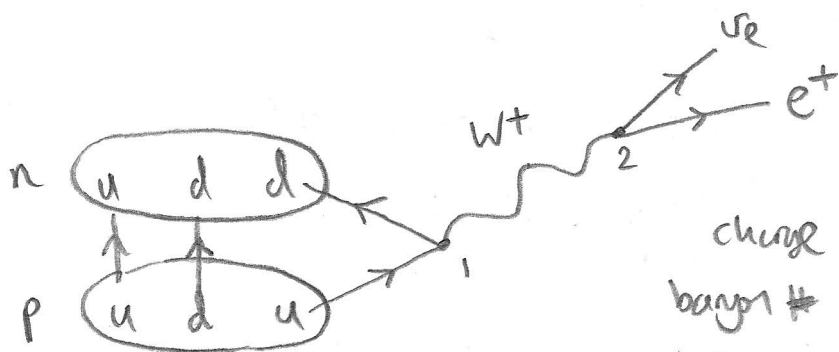
vertex 2:  $W^- \rightarrow e^- + \bar{\nu}_e$

charge:  $-e \rightarrow -e + 0 \quad \checkmark$   
 baryon #:  $0 \rightarrow 0 + 0 \quad \checkmark$   
 lepton #:  $0 \rightarrow 1 - 1 \quad \checkmark$

Note  $\bar{\nu}_e$  would have a lepton # of +1 so it must be an anti neutrino.

(ii)  $\beta^+$  decay:

$$p \rightarrow n + e^+ + \nu_e$$



charge:  $\frac{2}{3}e \rightarrow -\frac{1}{3}e + 1e \quad \checkmark$   
 baryon #:  $\frac{1}{3} \rightarrow \frac{1}{3} + 0 \quad \checkmark$   
 lepton #:  $0 \rightarrow 0 + 0 \quad \checkmark$

vertex 2:  $W^+ \rightarrow \nu_e + e^+$

charge:  $1e \rightarrow 0 + 1e \quad \checkmark$   
 baryon #:  $0 \rightarrow 0 + 0 \quad \checkmark$   
 lepton #:  $0 \rightarrow +1 - 1 \quad \checkmark$

$e^+$  is the antiparticle of  $e^-$ , so has lepton # of -1.

(iii)  $\pi^- + p \rightarrow n + \pi^- + \pi^+$

Quark content:  $\bar{u}d + udu \rightarrow udd + \bar{u}d + u\bar{d}$

$$\pi^- + p \rightarrow n + \pi^- + \pi^+$$

$$\bar{u}d + udu \rightarrow udd + \bar{u}d + u\bar{d}$$

charge:  
/e

$$-\frac{2}{3} - \frac{1}{3} + \frac{2}{3} - \frac{1}{3} + \frac{2}{3} \rightarrow \frac{2}{3} - \frac{1}{3} - \frac{1}{3} - \frac{2}{3} - \frac{1}{3} + \frac{2}{3} + \frac{1}{3}$$

$$-1 + 1 \rightarrow 0 - 1 + 1 \quad \checkmark$$

baryon #:

$$-\frac{1}{3} + \frac{1}{3} + \frac{1}{3} + \frac{1}{3} + \frac{1}{3} \rightarrow \frac{1}{3} + \frac{1}{3} + \frac{1}{3} + -\frac{1}{3} + \frac{1}{3} + \frac{1}{3} - \frac{1}{3}$$

$$0 + 1 \rightarrow 1 + 0 + 0 \quad \checkmark$$

(iv)  $k^+ \rightarrow \pi^+ + \pi^0$  kaon decay to pions  
 $u\bar{s} \rightarrow u\bar{d} + d\bar{d}$  Quark content of these Mesons.

charge:  
/e

$$\frac{2}{3} + \frac{1}{3} \rightarrow \frac{2}{3} + \frac{1}{3} + -\frac{1}{3} + \frac{1}{3}$$

$$1 \rightarrow 1 + 0 \quad \checkmark$$

(v)  $M_{k^\pm} c^2 + B_{k^\pm} = M_u c^2 + M_{\bar{s}} c^2$

$$B_{k^\pm} = 2.4 + 95 - 0.5261 \times 938.27 \quad (\text{MeV})$$

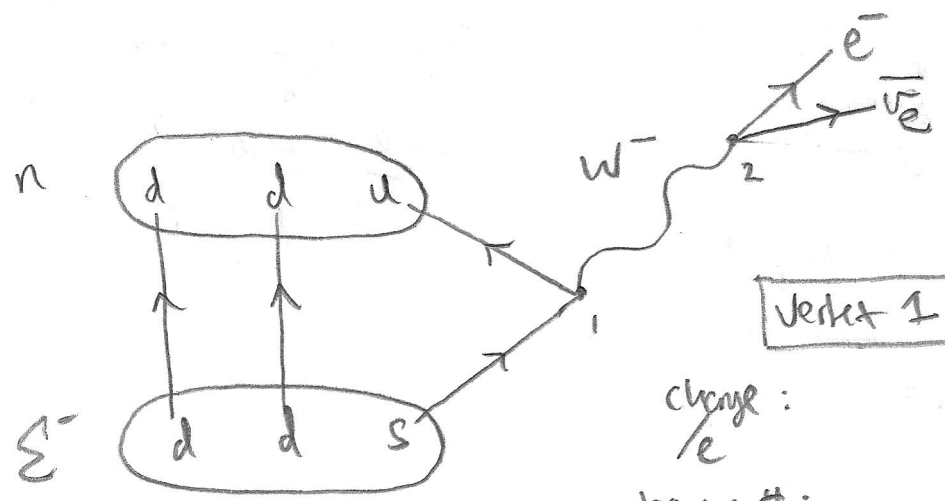
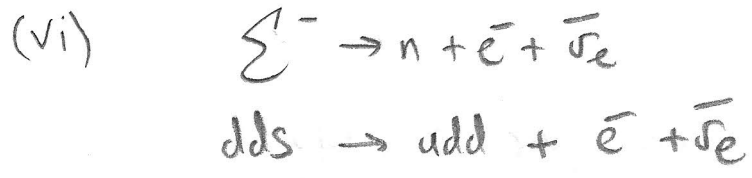
$$B_{k^\pm} = -396.2 \text{ MeV}$$

what does a -ve binding energy mean? Is it like a -ve gravitational potential energy? Interestingly, it appears most of the mass-energy of the kaon is not the mass-energy of the quarks - it is an energy held within the fields associated with the strong force.

$$B_{\pi^\pm} = m_u c^2 + m_{\bar{d}} c^2 - M_{\pi^\pm} c^2$$

$$= 2.4 + 4.8 - 0.14875 \times 938.27 \quad \text{MeV}$$

$$= -132.4 \text{ MeV}$$



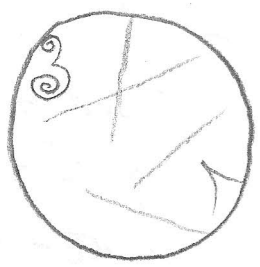
Vertex 1:  $S \rightarrow W^- + u$

change:  $-\frac{1}{3} \rightarrow -1 + \frac{2}{3} \checkmark$   
 $e$   
 baryon #:  $\frac{1}{3} \rightarrow 0 + \frac{1}{3} \checkmark$   
 lepton #:  $0 \rightarrow 0 + 0 \checkmark$

Vertex 2:  $W^- \rightarrow e^- + \bar{\nu}_e$

change:  $-1 \rightarrow -1 + 0 \checkmark$   
 $e$   
 baryon #:  $0 \rightarrow 0 + 0 \checkmark$   
 lepton #:  $0 \rightarrow 1 - 1 \checkmark$

viii) **CLOUD CHAMBER**



- \* Sealed environment containing a supersaturated vapour of water or alcohol.
- \* A high energy charged particle will leave a trail of ionized vapour molecules. These will act as 'condensation centres' about which droplets form. ie resulting in a 'cloud track' that can be photographed.

**BUBBLE CHAMBER**

\* High energy particle leaves a trail of bubbles, that can be photographed. Bubbles form on ionized molecules.

A bubble chamber is superheated, ie a liquid held above its boiling point under pressure. (whereas a cloud chamber is supersaturated). Reduce the pressure carefully, and bubbles form first on ionized molecules.

$$(VIII) \quad \gamma \rightarrow \mu^- + \mu^+$$

If the kinetic energy of muon pair is zero  
i.e. all the  $\gamma$  energy  $\rightarrow$  mass-energy of the muons

$$E_\gamma = 2m_\mu c^2$$

$$E_\gamma = \frac{hc}{\lambda} \quad \text{for a}$$

single photon.

$$\text{So} \quad \frac{hc}{\lambda} = 2m_\mu c^2$$

$$\Rightarrow \quad \boxed{\lambda = \frac{hc}{2m_\mu c^2}}$$

(Note if  $E > 2m_\mu c^2$  then clearly  $\lambda$  is smaller)

So minimum  $\lambda$  is  $\frac{hc}{2m_\mu c^2}$ .

$$\lambda = \frac{6.63 \times 10^{-34} + 2.998 \times 10^8}{2 \times 105.66 \times 10^6 + 1.602 \times 10^{-19}}$$

$$\boxed{\lambda = 5.87 \times 10^{-15} \text{ m}}$$

(ix) Total energy of  $e^- + e^+$  is  $\left(1 - \frac{u^2}{c^2}\right)^{-\frac{1}{2}} + 2m_e c^2$

$$= \left(1 - 0.8^2\right)^{-\frac{1}{2}} + 2 \times 0.51100 \text{ MeV}$$

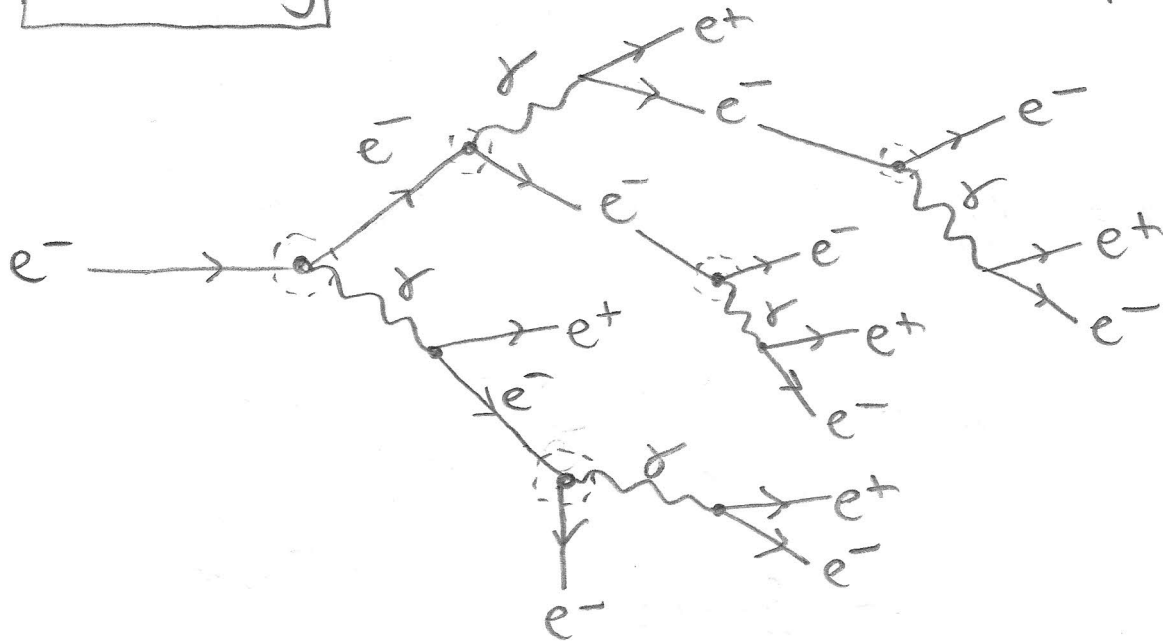
$$= \boxed{1.70 \text{ MeV}}$$

Total energy  
of  $e^- + e^+$   
is  $(KE + MC^2) \times 2$

(X)

# Bremsstrahlung

(radiation from rapidly accelerated particles)



$e^-$  could be accelerated, say by interaction with the dense outer shell electrons of Pb (lead) atoms. This acceleration causes an emission of a  $\gamma$  ray.

The  $\gamma$  ray then decays to form an  $e^+, e^-$  pair.

If the  $e^-$  has sufficient energy, it may produce a  $\gamma$  ray to and create more pairs etc.

$e^+$  could in principle do the same — but perhaps attraction to  $e^-$  (and its annihilation) in Pb might be more likely than repulsion from a Pb nucleus?

Either way, lots of extra  $e^-$  are produced, which enables the same  $e^-$  to be detected.

Note  $(KE)$  of  $e^-$  must be  $> 2m_e c^2$  for at least one pair to be produced. i.e.  $> 2 \times 0.51105 \text{ MeV}$   
 A 100 MeV  $e^-$  could make up to 97 pairs.

(5)

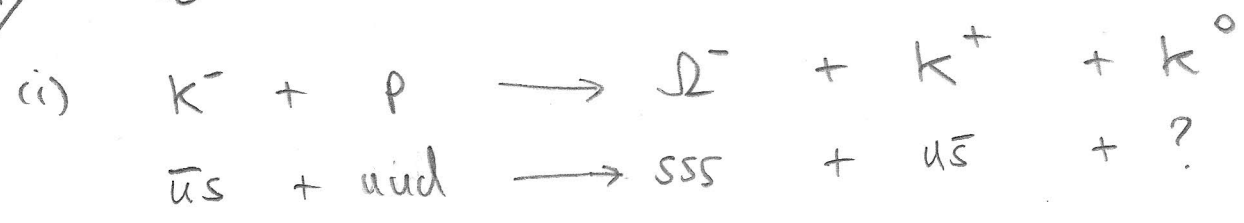
(xi) Koide mass ratio  $\frac{M_u + M_d + M_s}{(\sqrt{M_u} + \sqrt{M_d} + \sqrt{M_s})^2}$

=  $k/4$  <sup>2/3</sup> what is integer  $k$ ?

$\Rightarrow k = \frac{2.4 + 4.8 + 95}{(\sqrt{2.4} + \sqrt{4.8} + \sqrt{95})^2} \approx 9$   
 $= \boxed{5.106}$  so  $\boxed{k \approx 5}$

so  $\frac{M_u + M_d + M_s}{(\sqrt{M_u} + \sqrt{M_d} + \sqrt{M_s})^2} \approx \frac{5}{4} = \boxed{\frac{5}{6} \times \frac{2}{3}}$

2/  $\Omega^- = sss$



Charge:  $-\frac{2}{3} + -\frac{1}{3} + \frac{2}{3} + \frac{2}{3} - \frac{1}{3} \rightarrow +3 + -\frac{1}{3} + \frac{2}{3} + \frac{1}{3} + ?$   
 $-1 + 1 \rightarrow -1 + 1 + \textcircled{+0}$

so charge of  $K^0$  must be 0 (obvious from it's superscript!)

Baryon #:  $-\frac{1}{3} + \frac{1}{3} + 1 \rightarrow 1 + \frac{1}{3} - \frac{1}{3} + \textcircled{0}$

so baryon # of  $K^0$  must be 0

strangeness:  $-1 + 0 \rightarrow -3 + 1 + ?$   
 $? \text{ must be } +1 \text{ if strangeness of } K^0.$

So in summary:  $K^0$  must have an  $\bar{s}$  quark for its strangeness to be +1. (And no other strange quarks)

\* To have a baryon # of 0 it must be a meson i.e. a quark, antiquark pair

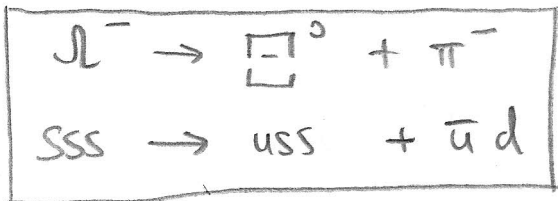
\* charge of  $\bar{s}$  is  $+\frac{1}{3}e$ , so to have a net zero charge, remaining quark must have a charge of  $-\frac{1}{3}e$ .

options are:  $[d, b]$  (since  $s$  not allowed).

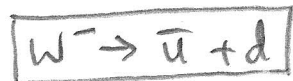
$b$  is too massive to be created in the interaction

so  $\therefore [d]$  so  $K^0 = \bar{s}d$

(ii)



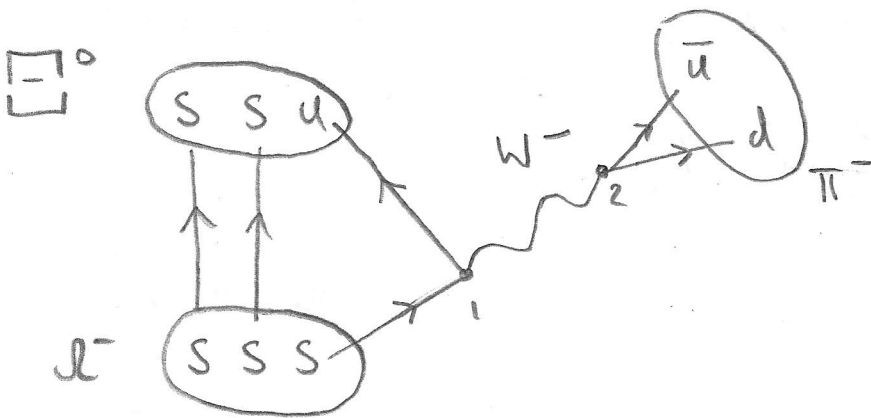
Vertex 2:



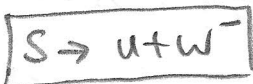
charge:  $-1 \rightarrow -\frac{2}{3} + -\frac{1}{3} \checkmark$

Baryon #:  $0 \rightarrow -\frac{1}{3} + \frac{1}{3} \checkmark$

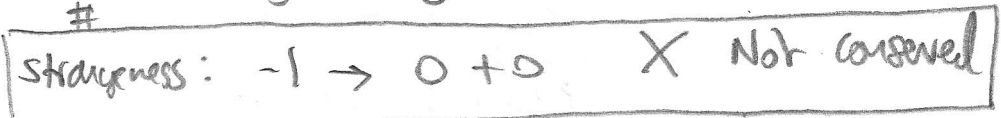
strangeness:  $0 \rightarrow 0 + 0 \checkmark$



Vertex 1: charge:  $-\frac{1}{3} \rightarrow \frac{2}{3} - 1 \checkmark$



Baryon #:  $\frac{1}{3} \rightarrow \frac{1}{3} + 0 \checkmark$

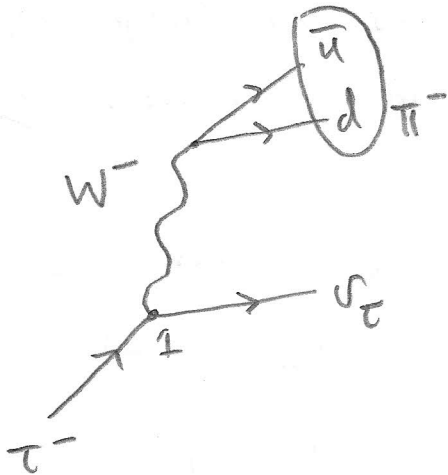
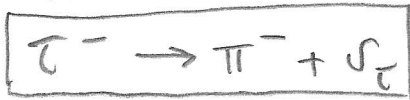


Strangeness not conserved in weak interactions.

(7)

3/

(i)



(ii) lepton number must be conserved at vertex 1



$$\text{lepton \#}: 1 \rightarrow 0 + 1 \quad \checkmark$$

If the  $\nu_{\tau}$  was not produced, lepton # would not be conserved since  $W^-$  is a boson, and  $\therefore$  has a lepton # of 0.

(iii)

$$E = \left(1 - \frac{u^2}{c^2}\right)^{-\frac{1}{2}} mc^2$$

total energy of  $\tau^-$

$\therefore$  kinetic energy is  $E - mc^2$

$$E_k = \left[ \left(1 - 0.2^2\right)^{-\frac{1}{2}} - 1 \right] \times 1776.86 \text{ MeV}$$

$$= \boxed{36.64 \text{ MeV}}$$

(iv)

in  $\tau^- \rightarrow \pi^- + \nu_{\tau}$ , the pion has a speed of  $0.1c$

so

$$E_{\nu_{\tau}} = E_{\tau^-} - E_{\pi^-}$$

$$\left( \text{using } E = \gamma mc^2 \right) \gamma = \left(1 - \frac{u^2}{c^2}\right)^{-\frac{1}{2}}$$

$$= \left(1 - 0.2^2\right)^{-\frac{1}{2}} \times 1776.86 - \left(1 - 0.01\right)^{-\frac{1}{2}} \times 0.14875 \times 938.27$$

$$= 1813.50 - 989.37 \text{ MeV}$$

$$= \boxed{824.13 \text{ MeV}}$$

(MeV)



(v) Relativistic momentum:

$$p = \left(1 - \frac{u^2}{c^2}\right)^{-\frac{1}{2}} m u$$

$$E = \left(1 - \frac{u^2}{c^2}\right)^{-\frac{1}{2}} m c^2 \quad \text{Total energy}$$

$$\text{So } E^2 - p^2 c^2 = \left(1 - \frac{u^2}{c^2}\right)^{-1} m^2 c^4 - \left(1 - \frac{u^2}{c^2}\right)^{-1} m^2 u^2 c^2$$

$$= \left(1 - \frac{u^2}{c^2}\right)^{-1} m^2 c^4 \left(1 - \frac{u^2}{c^2}\right)$$

$$= m^2 c^4$$

$$\text{So } \boxed{E^2 - p^2 c^2 = m^2 c^4}$$

Energy - momentum invariant

$$\text{if } m = 0 \Rightarrow E = p c \quad \text{So } \boxed{p = \frac{E}{c}}$$

$$\therefore p_{\nu_e} = \boxed{24.13 \text{ MeV}/c}$$

$$= \frac{24.13 \times 10^6 + 1.602 \times 10^{-19}}{2.998 \times 10^8}$$

$$= \boxed{4.40 \times 10^{-19} \text{ kgms}^{-1}}$$

$$\text{(vi) a) } \frac{m_e + m_\mu + m_\tau}{(\sqrt{m_e} + \sqrt{m_\mu} + \sqrt{m_\tau})^2} = \frac{0.511 + 105.166 + 1776.86}{(\sqrt{0.511} + \sqrt{105.166} + \sqrt{1776.86})^2}$$

$$= 0.667 \approx \boxed{\frac{2}{3}}$$

Kaide formula for leptons.

⑨

$$b) \frac{M_c + M_b + M_t}{(\sqrt{M_c} + \sqrt{M_b} + \sqrt{M_t})^2} = \frac{1775 + 4180 + 172440}{(\sqrt{1775} + \sqrt{4180} + \sqrt{172440})^2}$$

$$= 0.669 \approx \boxed{\frac{2}{3}}$$

krude formula for heavy quarks.

Are these formula serendipity? or do they hint at a deeper unifying mathematics of the standard model.....

4/

### Compton Scattering



BEFORE



AFTER

(i) Energy-momentum invariant (for photon of mass  $m=0$ )

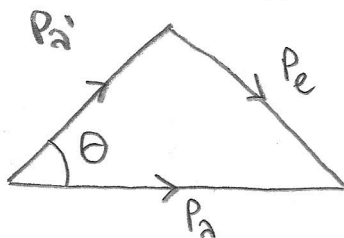
$$E_a^2 - p_a^2 c^2 = 0$$

$$\Rightarrow p_a = E_a / c$$

$$\Rightarrow p_a = \frac{hc}{\lambda c}$$

$$\Rightarrow \boxed{p_a = \frac{h}{\lambda}} \quad \text{de-Broglie}$$

(ii) Momentum conservation  $\underline{p}_a = \underline{p}_{a'} + \underline{p}_e$



By cosine rule:

$$\boxed{p_e^2 = p_a^2 + p_{a'}^2 - 2 p_a p_{a'} \cos \theta}$$

(iii) Conservation of energy

$$m_e c^2 + E_a = E_{a'} + E_e$$

$$\therefore E_e = m_e c^2 + E_a - E_{a'}$$

$$E_a = p_a c \quad E_{a'} = p_{a'} c$$

and

$$E_e^2 - p_e^2 c^2 = m_e^2 c^4$$

} Energy-momentum invariants

$$\text{So } p_e^2 = \frac{E_e^2}{c^2} - m_e^2 c^2$$

$$p_e^2 = \frac{(m_e c^2 + (p_a - p_{a'})c)^2}{c^2} - m_e^2 c^2$$

$$p_e^2 = \frac{m_e^2 c^4 + (p_a - p_{a'})^2 c^2 + 2m_e c^2 (p_a - p_{a'})c - m_e^2 c^4}{c^2}$$

$$p_e^2 = m_e^2 c^2 + (p_a - p_{a'})^2 + 2m_e c (p_a - p_{a'}) - m_e^2 c^2$$

$$\therefore p_e^2 = (p_a - p_{a'})^2 + 2m_e c (p_a - p_{a'})$$

(iv) Equating with  $p_e^2 = p_a^2 + p_{a'}^2 - 2p_a p_{a'} \cos \theta$ :

$$p_a^2 + p_{a'}^2 - 2p_a p_{a'} \cos \theta = p_a^2 + p_{a'}^2 - 2p_a p_{a'} + 2m_e c (p_a - p_{a'})$$

$$2p_a p_{a'} (1 - \cos \theta) = 2m_e c (p_a - p_{a'})$$

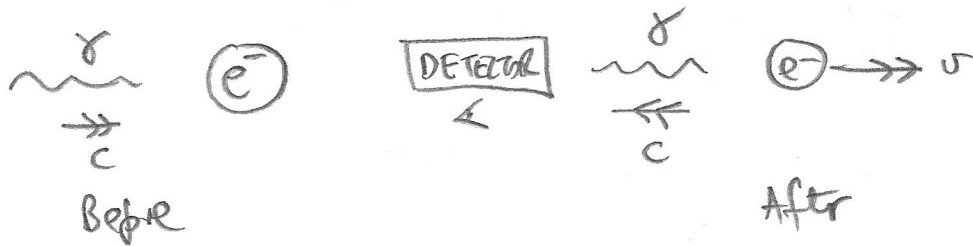
$$\text{Now use } p_a = \frac{h}{\lambda} \quad \text{and} \quad p_{a'} = \frac{h}{\lambda'}$$

$$\Rightarrow \frac{h^2}{\lambda \lambda'} (1 - \cos \theta) = m_e c h \left( \frac{1}{\lambda} - \frac{1}{\lambda'} \right)$$

$$\therefore \frac{h}{m_e c} (1 - \cos \theta) = \lambda' - \lambda$$

$$v) \quad \Delta\lambda = \lambda' - \lambda = \frac{h}{m_e c} (1 - \cos\theta)$$

Maximum  $\Delta\lambda$  is when  $\cos\theta = -1$  i.e.  $\theta = 180^\circ$   
 i.e. direct recoil of the X-ray



$$\therefore \quad \Delta\lambda_{\max} = \frac{2h}{m_e c}$$

$$\text{So } \frac{\Delta\lambda_{\max}}{\lambda} = \frac{2h}{m_e c \lambda} \quad E_\gamma = \frac{hc}{\lambda} \quad \text{Planck}$$

$$\text{So: } \frac{1}{\lambda} = \frac{E_\gamma}{hc}$$

$$\therefore \quad \frac{\Delta\lambda_{\max}}{\lambda} = \frac{2h}{m_e c} \frac{E_\gamma}{hc}$$

$$\therefore \quad \frac{\Delta\lambda_{\max}}{\lambda} = \frac{2E_\gamma}{m_e c^2}$$

$$[m_e = 0.51100 \text{ MeV}/c^2]$$

$$\text{So } \quad E_\gamma = \frac{1}{2} m_e c^2 \frac{\Delta\lambda_{\max}}{\lambda}$$

$$\text{If } \frac{\Delta\lambda_{\max}}{\lambda} = 0.1 \quad \Rightarrow \quad E_\gamma = \frac{1}{2} \times 511 \text{ keV} \times 0.1$$

$$E_\gamma = 25.6 \text{ keV}$$

(vi) To conserve momentum:

$$\parallel x: \quad P_a = P_{a'} \cos \theta + P_e \cos \phi$$

$$\parallel y: \quad 0 = P_{a'} \sin \theta - P_e \sin \phi$$

$$\text{so } P_e \sin \phi = P_{a'} \sin \theta$$

$$P_e \cos \phi = P_a - P_{a'} \cos \theta$$

$$\therefore \tan \phi = \frac{P_{a'} \sin \theta}{P_a - P_{a'} \cos \theta} = \frac{\sin \theta}{\frac{P_a}{P_{a'}} - \cos \theta}$$

$$\text{Now } \frac{P_a}{P_{a'}} = \frac{h}{a} \frac{\lambda'}{h} = \frac{\lambda'}{\lambda}$$

$$\lambda' - \lambda = \frac{h}{m_e c} (1 - \cos \theta)$$

$$\text{so } \frac{\lambda'}{\lambda} - 1 = \frac{h}{m_e c \lambda} (1 - \cos \theta)$$

$$\text{so } \frac{P_a}{P_{a'}} = \frac{\lambda'}{\lambda} = 1 + \frac{h}{m_e c \lambda} (1 - \cos \theta)$$

$$\therefore \tan \phi = \frac{\sin \theta}{1 + \frac{h}{m_e c \lambda} (1 - \cos \theta) - \cos \theta}$$

$$\text{Now } P_e \sin \phi = P_{a'} \sin \theta \quad \therefore P_e = \frac{h}{\lambda'} \frac{\sin \theta}{\sin \phi}$$

$$\lambda' = \lambda + \frac{h}{m_e c} (1 - \cos \theta)$$

$$\therefore p_e = \frac{h \sin \theta}{\sin \phi} \left( \lambda + \frac{h}{m_e c} (1 - \cos \theta) \right)^{-1}$$

Now  $p_e = \left( 1 - \frac{v^2}{c^2} \right)^{-\frac{1}{2}} m_e v$

$$p_e^{-2} = \left( 1 - \frac{v^2}{c^2} \right) \frac{1}{m_e^2 v^2}$$

$$\frac{1}{p_e^2} = \frac{1}{m_e^2 v^2} - \frac{1}{m_e^2 c^2}$$

$$\therefore v^2 = \left( \frac{m_e^2}{p_e^2} + \frac{1}{c^2} \right)^{-1}$$

$$v = \left( \frac{m_e^2}{p_e^2} + \frac{1}{c^2} \right)^{-\frac{1}{2}}$$

So  $v = \left( \frac{m_e^2}{p_e^2} + \frac{1}{c^2} \right)^{-\frac{1}{2}} \leftarrow \text{or } v = c \left( \frac{m_e c}{p_e} \right)^2 + 1)^{-\frac{1}{2}}$

where  $p_e = \frac{h \sin \theta}{\sin \phi} \left( \lambda + \frac{h}{m_e c} (1 - \cos \theta) \right)^{-1}$

and  $\tan \phi = \frac{\sin \theta}{1 + \frac{h}{m_e c \lambda} (1 - \cos \theta) - \cos \theta}$

Perhaps it is preferable to determine  $P_e$  in terms of  $\lambda, \theta$  directly (rather than via  $\phi$ ).

From (iii)  $P_e^2 = (P_a - P_{a'})^2 + 2m_e c (P_a - P_{a'})$  (\*)

$P_a = \frac{h}{\lambda}$  and  $P_{a'} = \frac{h}{\lambda'}$   $\lambda' = \lambda + \frac{h}{m_e c} (1 - \cos\theta)$

So  $P_e = \sqrt{(P_a - P_{a'})^2 + 2m_e c (P_a - P_{a'})}$

Now when  $\cos\theta = -1$  ( $\theta = 180^\circ$ ),  $\lambda'$  is maximized

so since  $P_{a'} = \frac{h}{\lambda'}$  this means  $P_{a'}$  is minimized

$\therefore P_a - P_{a'}$  is maximized, which means  $P_e$  is maximized. (From \*)

$\lambda'_{max} = \lambda + \frac{2h}{m_e c}$

So  $P_a - P_{a'} = \frac{h}{\lambda} - \frac{h}{\lambda + \frac{2h}{m_e c}} = \frac{h}{\lambda} \left( 1 - \frac{1}{1 + \frac{2h}{\lambda m_e c}} \right)$

So  $P_e$  max is  $\sqrt{\frac{h^2}{\lambda^2} \left( 1 - \left( 1 + \frac{2h}{\lambda m_e c} \right)^{-1} \right)^2 + \frac{2m_e c h}{\lambda} \left( 1 - \left( 1 + \frac{2h}{\lambda m_e c} \right)^{-1} \right)}$

Although a better way of finding it is to go back to conservation of energy



$$\lambda' = \lambda + \frac{2h}{m_e c}$$

$$\left( \text{From } \Delta \lambda = \frac{h}{m_e c} (1 - \cos \theta) \right)$$

and  $\theta = +180^\circ$

Conservation of energy

$$\frac{hc}{\lambda} + m_e c^2 = \frac{hc}{\lambda'} + \left(1 - \frac{v^2}{c^2}\right)^{-\frac{1}{2}} m_e c^2$$

$$\therefore \left( \frac{\frac{hc}{\lambda} - \frac{hc}{\lambda'} + m_e c^2}{m_e c^2} \right)^{-2} = 1 - \frac{v^2}{c^2}$$

$$\therefore v = c \sqrt{1 - \left( \frac{m_e c^2}{m_e c^2 + \frac{hc}{\lambda} - \frac{hc}{\lambda'}} \right)^2}$$