

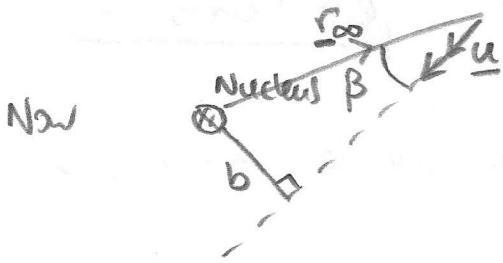
Now since the force on the particle is always in the \hat{r} direction, there can be no net torque. \therefore angular momentum is constant.

$$\underline{L} = m \underline{r} \times \underline{v}$$

In polar coordinates $\underline{r} \times \underline{v} = r \hat{r} \times (\dot{r} \hat{r} + r \dot{\theta} \hat{\theta})$
 $= r^2 \dot{\theta} \hat{r} \times \hat{\theta}$

$$\therefore L = |\underline{L}| = \boxed{m r^2 \dot{\theta}} \quad (4)$$

So $r^2 \dot{\theta}$ is a constant.



is the initial condition, in relation to the closest approach distance b (which is the 'no deflection' scenario)

From the definition of the cross product $|\underline{r} \times \underline{u}| = r u \sin \beta$
 $= ub$ $\therefore m r^2 \dot{\theta} = m u b$

$$\boxed{r^2 \dot{\theta} = u b} \quad (5)$$

$$\therefore \text{in } (3) \quad u^2 = \dot{r}^2 + r^2 \left(\frac{ub}{r^2} \right)^2 + \frac{Ze^2}{4\pi\epsilon_0 M r} \quad (6)$$

Now let $\boxed{z = \frac{1}{r}}$

$$\frac{dz}{dt} = -\frac{1}{r^2} \frac{dr}{dt} \quad \therefore \dot{r} = -r^2 \frac{dz}{dt}$$

Now $\frac{dz}{dt} = \frac{dz}{d\theta} \frac{d\theta}{dt} = \dot{\theta} \frac{dz}{d\theta} \quad \therefore \dot{r} = -r^2 \dot{\theta} \frac{dz}{d\theta}$

Since $r^2 \dot{\theta} = ub \quad \Rightarrow \quad \boxed{\dot{r} = -ub \frac{dz}{d\theta}}$

Hence (6) becomes: $u^2 = u^2 b^2 \left(\frac{dz}{d\theta} \right)^2 + u^2 b^2 z^2 + \frac{Ze^2}{4\pi\epsilon_0 M} z$

Since u is a constant, $\frac{d}{d\theta}(u^2) = 0$

$$\therefore uz^2 \left(2 \frac{dz}{d\theta} \right) \frac{d^2z}{d\theta^2} + 2u^2bz^2 \frac{dz}{d\theta} + \frac{ze^2}{4\pi\epsilon_0 m} \frac{dz}{d\theta} = 0$$

Since $\frac{dz}{d\theta} \neq 0 \quad \forall \theta$:

$$\Rightarrow \boxed{\frac{d^2z}{d\theta^2} + z + \frac{ze^2}{4\pi\epsilon_0 mu^2b^2} = 0}$$

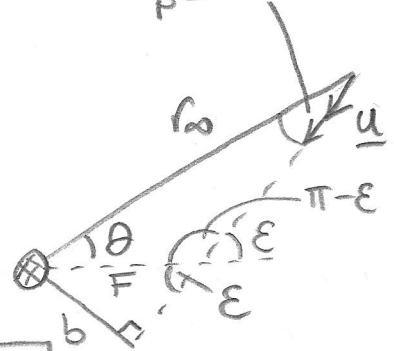
consider solutions of the form $\boxed{z(\theta) = A \cos \theta + B}$

$$\frac{dz}{d\theta} = -A \sin \theta \quad \frac{d^2z}{d\theta^2} = -A \cos \theta = B - z$$

$$\therefore B - z + z + \frac{ze^2}{4\pi\epsilon_0 mu^2b^2} = 0$$

$$\Rightarrow \boxed{B = -\frac{ze^2}{4\pi\epsilon_0 mu^2b^2}}$$

$$\begin{aligned} \beta &= \pi - (\pi - \epsilon + \theta) \\ \beta &= \epsilon - \theta \end{aligned}$$



Now velocity $\underline{v} = \dot{r} \hat{r} + r \dot{\theta} \hat{\theta}$

From above, $\boxed{\dot{r} = -ub \frac{dz}{d\theta}}$ and $\boxed{\dot{\theta} = \frac{ub}{r^2}}$

$$\therefore \underline{v} = -ub \frac{dz}{d\theta} \hat{r} + \frac{ub}{r} \hat{\theta}$$

$$\therefore \underline{v} = A ub \sin \theta \hat{r} + ub z \hat{\theta} \quad \leftarrow \boxed{\frac{1}{r} = z}$$

when $r \rightarrow \infty, z \rightarrow 0, \underline{v} \rightarrow \underline{u}, \theta \rightarrow \epsilon$

$$\therefore u = A ub \sin \epsilon \Rightarrow \boxed{A = \frac{1}{b \sin \epsilon}}$$

$$\text{So } 0 = \frac{1}{b \sin \epsilon} \cos \epsilon - \frac{ze^2}{4\pi\epsilon_0 mu^2b^2} \quad \begin{aligned} [z = A \cos \theta + B] \\ [z \rightarrow 0, \theta \rightarrow \epsilon] \end{aligned}$$

$$\Rightarrow \frac{ze^2}{4\pi\epsilon_0 mu^2 b^2} = \frac{1}{b} \cot \epsilon$$

$$\Rightarrow b = \frac{ze^2}{4\pi\epsilon_0 mu^2} \tan \epsilon \Rightarrow \tan \epsilon = \frac{4\pi\epsilon_0 mu^2 b}{ze^2}$$

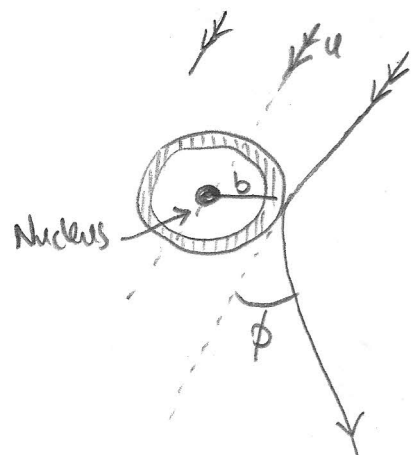
Now from geometry $\phi + 2\epsilon = \pi \quad \therefore \epsilon = \frac{\pi - \phi}{2}$

$$\therefore \tan \epsilon = \frac{\tan\left(\frac{\pi}{2}\right) - \tan\frac{\phi}{2}}{1 + \tan\frac{\pi}{2} \tan\frac{\phi}{2}} = \frac{1 - \tan\frac{\phi}{2}}{\frac{1}{\tan\frac{\pi}{2}} + \tan\frac{\phi}{2}}$$

Now $\tan\frac{\pi}{2} \rightarrow \infty$ so $\tan \epsilon = \cot \frac{\phi}{2}$

$$\boxed{\cot \frac{\phi}{2} = \frac{4\pi\epsilon_0 mu^2 b}{ze^2}}$$

Now consider a beam of particles incident upon a nucleus



The area of the beam cross section which corresponds to scattering within angular range ϕ to $\phi + d\phi$ is $\boxed{2\pi b db}$, since $b = b(\phi)$

$$\text{Now } \frac{d}{d\phi} \cot \frac{\phi}{2} = \frac{d}{d\phi} \frac{1}{\tan \frac{\phi}{2}} = -\frac{1}{\tan^2 \frac{\phi}{2}} \sec^2 \frac{\phi}{2} \times \frac{1}{2}$$

$$= -\frac{1}{\frac{\sin^2 \frac{\phi}{2}}{\cos^2 \frac{\phi}{2}}} \times \frac{1}{\cos^2 \frac{\phi}{2}} \times \frac{1}{2}$$

$$= \boxed{-\frac{1}{2} \frac{1}{\sin^2 \frac{\phi}{2}}}$$

$$-\frac{1}{2} \frac{1}{\sin^2 \frac{\phi}{2}} = \frac{4\pi\epsilon_0 \mu v^2}{ze^2} \frac{db}{d\phi}$$

$$\Rightarrow db = \frac{1 - ze^2}{2 \cdot 4\pi\epsilon_0 \mu v^2} \frac{d\phi}{\sin^2 \frac{\phi}{2}}$$

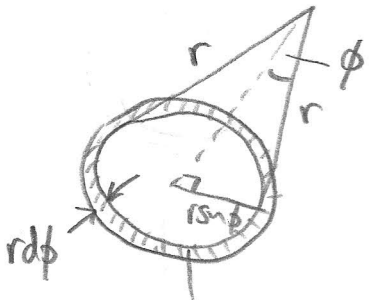
$$\therefore 2\pi b db = 2\pi \left(\frac{ze^2}{4\pi\epsilon_0 \mu v^2} \right) \frac{\cos \frac{\phi}{2}}{\sin^3 \frac{\phi}{2}} \frac{1}{2} \left(\frac{-ze^2}{4\pi\epsilon_0 \mu v^2} \right) \frac{d\phi}{\sin^2 \frac{\phi}{2}}$$

$$\left[b = \frac{ze^2 \cos \frac{\phi}{2}}{4\pi\epsilon_0 \mu v^2 \sin^2 \frac{\phi}{2}} \right] = -\pi \left(\frac{ze^2}{4\pi\epsilon_0 \mu v^2} \right)^2 \frac{\cos \frac{\phi}{2}}{\sin^3 \frac{\phi}{2}} d\phi$$

Now α particles scattered by angle ϕ corresponds to a Solid angle

$$d\Omega = 2\pi \sin \phi d\phi$$

$$= \boxed{4\pi \sin^2 \frac{\phi}{2} \cos^2 \frac{\phi}{2} d\phi}$$



Area of sphere of radius r corresponding to scattered α particles between ϕ and $\phi + d\phi$

$$\text{is } 2\pi r \sin \phi \cdot r d\phi$$

$$\boxed{dA = 2\pi r^2 \sin \phi d\phi}$$

So Solid angle $d\Omega = \frac{dA}{r^2}$

$$= \boxed{2\pi \sin \phi d\phi}$$

Define differential cross section $\frac{d\sigma}{d\Omega}$ as the

area of the incident beam which corresponds to a unit Solid angle of the scattered α particles

$$\text{let } d\sigma = |2\pi b db|$$

$$\therefore \frac{d\sigma}{d\Omega} = \frac{\pi \left(\frac{ze^2}{4\pi\epsilon_0 \mu v^2} \right)^2 \frac{\cos \frac{\phi}{2} d\phi}{\sin^3 \frac{\phi}{2}}}{4\pi \sin^2 \frac{\phi}{2} \cos^2 \frac{\phi}{2} d\phi}$$

$$\boxed{\frac{d\sigma}{d\Omega} = \left(\frac{ze^2}{8\pi\epsilon_0 \mu v^2} \right)^2 \frac{1}{\sin^4 \frac{\phi}{2}}}$$

$$\text{i.e. } \propto \csc^4 \frac{\phi}{2}$$

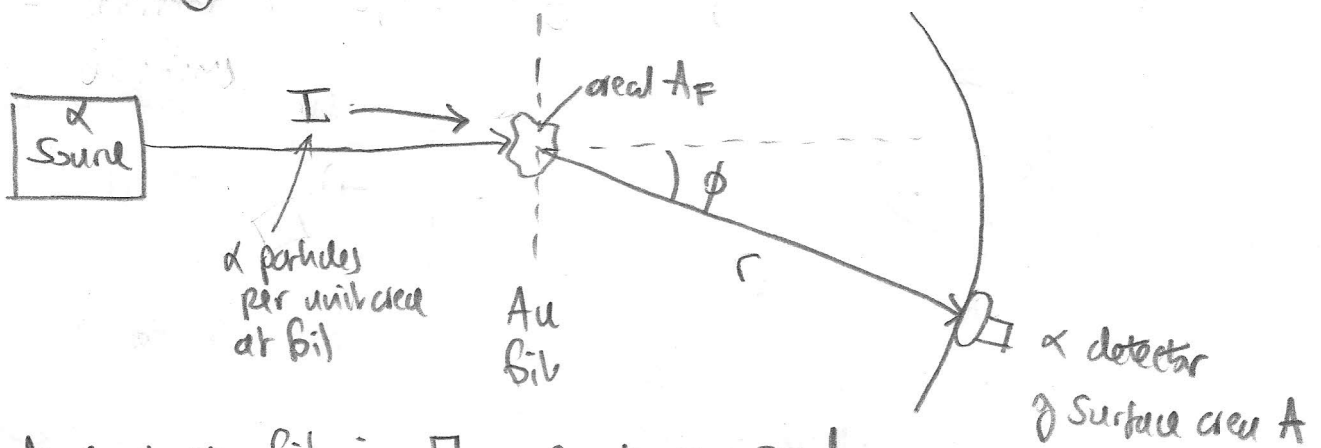
Note if I is the incident flux of α particles per unit area

$|2\pi b db| I$ would be the flux of α particles in scattering direction ϕ (well $\phi \rightarrow \phi + d\phi$)

$$\therefore \frac{d\sigma}{d\Omega} = \frac{|2\pi b db|}{d\Omega} = \frac{|2\pi b db| I}{d\Omega I}$$

= $\frac{\text{flux of } \alpha \text{ particles in scattering direction } \phi \text{ (per unit solid angle)}}{\text{incident flux per unit area}}$

This ratio could be measured experimentally, as in the classic by Rutherford, Geiger and Marsden in 1909-11.



Particle count at foil is Γ_0 counts per second
 $= I A_F$ (assume uniformly illuminated)

This is measured

Particle count at detector is $\Gamma = \frac{A}{r^2} \times \frac{d\sigma}{d\Omega} I \times N_{Au}$

Labels for the equation above:
 $\frac{A}{r^2}$: Solid angle of detector
 $\frac{d\sigma}{d\Omega}$: flux per unit solid angle
 N_{Au} : # atoms in foil within beam

$N_{Au} = \# \text{ atoms in foil of area}$

A_F and thickness $t_F = n A_F t_F$ where $n = \text{number of atoms / unit volume}$

So
$$\Gamma = \frac{A}{r^2} \times \left(\frac{ze^2}{4\pi\epsilon_0\mu v^2} \right)^2 \frac{1}{\sin^4 \frac{\phi}{2}} \frac{\Gamma_0}{A} \times n A t_F$$

$$\Rightarrow \frac{\Gamma}{\Gamma_0} = \frac{A t_F n}{r^2} \left(\frac{ze^2}{4\pi\epsilon_0\mu v^2} \right)^2 \frac{1}{\sin^4 \frac{\phi}{2}}$$

For Au: $z = 79$
 $n \approx 5.9 \times 10^{28} \text{ m}^{-3}$
 $A \sim 15 \text{ mm}^2 = 1.5 \times 10^{-5} \text{ m}^2$
 $t_F \sim 1.5 \mu\text{m} = 1.5 \times 10^{-6} \text{ m}$

Note $E_k = \frac{1}{2} \mu v^2$
 ↑ kinetic energy

So
$$\frac{\Gamma}{\Gamma_0} = \frac{A t_F n}{r^2} \left(\frac{ze^2}{4\pi\epsilon_0 E} \right)^2 \frac{1}{\sin^4 \frac{\phi}{2}}$$

$$\therefore \log_{10} \frac{\Gamma}{\Gamma_0} = \log_{10} \left(\frac{A t_F n}{r^2} \left(\frac{ze^2}{4\pi\epsilon_0 E} \right)^2 \right) + 4 \log_{10} \left(\frac{1}{\sin \frac{\phi}{2}} \right)$$

$\Gamma_0 \approx 1000$ counts/min
 ↑ Typical for 'School Sage' source.

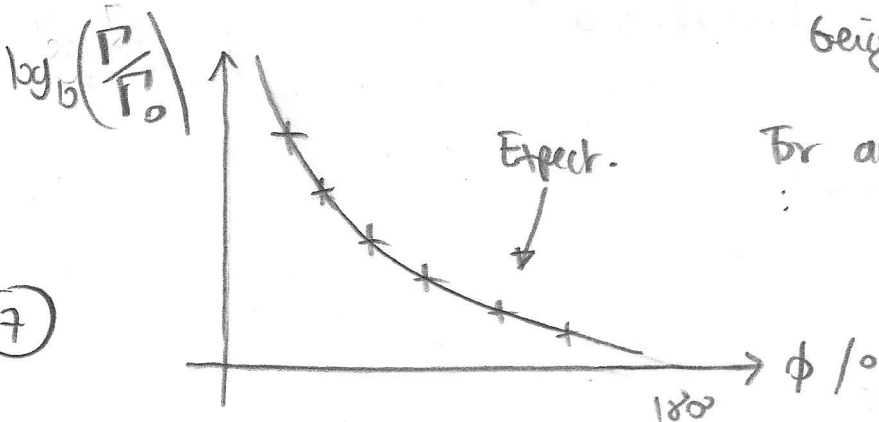
$E_d \sim 3 \text{ MeV}$

$e \sim 1.602 \times 10^{-19} \text{ C}$

$r \sim 0.1 \text{ m}$

$\epsilon_0 = 8.85 \times 10^{-12} \text{ Fm}^{-1}$

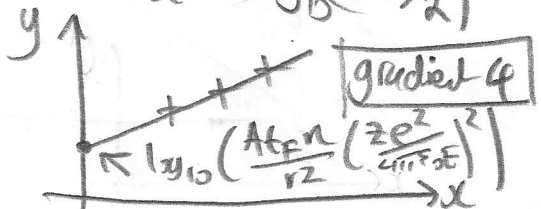
Guess at dimensions of Geiger & Marsden's apparatus



For analysis

$y = \log_{10} \left(\frac{\Gamma}{\Gamma_0} \right)$

$x = \log_{10} \left(\frac{1}{\sin \frac{\phi}{2}} \right) \leftarrow \text{+ve}$

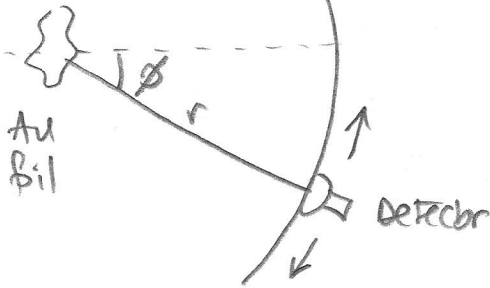


$[\sin \frac{\phi}{2} < 1 \text{ so } \log_{10} \sin \frac{\phi}{2} < 0] \leftarrow \text{Hence use } \frac{1}{\sin \frac{\phi}{2}} \text{ so } \log_{10} \text{ will be +ve.}$

Rutherford Scattering Summary

Experiment:

α Source



$$\frac{\Gamma}{\Gamma_0} = \frac{At_F n}{r^2} \left(\frac{ze^2}{4\pi\epsilon_0 E} \right)^2 \frac{1}{\sin^4 \frac{\phi}{2}}$$

Γ particle count rate at detector

Γ " " " at foil

t_F foil thickness

n # atoms per m^3 in foil

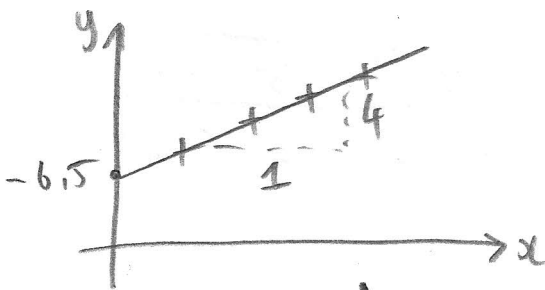
E α particle KE

z atomic number (79 for gold)

e charge on electron

$$\underbrace{\log_{10} \frac{\Gamma}{\Gamma_0}}_y = \log_{10} \left[\frac{At_F n}{r^2} \left(\frac{ze^2}{4\pi\epsilon_0 E} \right)^2 \right] + 4 \log_{10} \left(\frac{1}{\sin \frac{\phi}{2}} \right)_x$$

ie predict $y = 4x + C$



For sensible numbers:

$$\frac{At_F n}{r^2} \left(\frac{ze^2}{4\pi\epsilon_0 E} \right)^2$$

$$\approx 1.9101 \times 10^{-4}$$

